$CaTMg_2$ and $CaTCd_2$ (T = Rh, Pd, Pt) with YPd_2Si -type Structure

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The intermetallic calcium compounds $CaTMg_2$ and $CaTCd_2$ (T = Rh, Pd, Pt) were obtained by high-frequency melting of the elements in sealed niobium ampoules or through reactions in muffle furnaces. The polycrystalline samples were characterized by powder X-ray diffraction. They crystallize with a site occupancy variant of YPd₂Si, a ternary ordered version of Fe₃C. The structures of $CaPdMg_2$ and $CaPdCd_2$ were refined from single-crystal diffractometer data: Pnma, a = 792.2(2), b = 803.4(2), c = 572.0(1) pm, wR2 = 0.0663, 1621 F^2 values, 24 variables for $Ca_{0.94}PdMg_{2.06}$ and a = 794.6(2), b = 809.5(3), c = 554.7(2) pm, wR2 = 0.0301, $819 F^2$ values, 23 variables for CaPdCd₂. A small range of homogeneity was observed for $Ca_{1-x}PdMg_{2+x}$. The magnesium and cadmium atoms build up three-dimensional tetrahedral substructures (306–327 pm Mg–Mg and 307-317 pm Cd-Cd) that resemble hexagonal diamond, lonsdaleite. Together with the palladium atoms one obtains three-dimensional, covalently bonded [PdMg2] and [PdCd2] networks which leave cages for the calcium atoms. The latter are bonded to these networks via shorter Ca-Pd contacts $(298-319 \text{ pm in } \text{Ca}_{0.94}\text{PdMg}_{2.06} \text{ and } 295-312 \text{ pm in } \text{Ca}\text{PdCd}_2)$. The course of the interatomic distances is in line with calculated overlap populations. The CaPdMg₂, SrPdMg₂ and CaRhIn₂ structures are all derived from a CaIn₂-related subcell by an ordered filling of transition metal atoms into trigonal prisms. This leads to different herringbone patterns for the networks of puckered and elongated hexagons of magnesium and indium atoms.

Key words: Magnesium, Cadmium, Cementite Phases, Electronic Structure

Introduction

The structure of the S-phase precipitate MgCuAl₂ [1, 2] is the prototype for a series of aluminides, gallides, indides, and thallides ATX_2 where A can be a rare earth or alkaline earth element, and T is an electron-rich transition metal. The X atoms build up lonsdaleite-related networks, and the ATX_2 structures can formally be described as transition metal-filled versions of binary Zintl phases AX_2 (CaIn₂ type) [3]. Although the MgCuAl₂-type structure is typically formed for triels, recently also some magnesium and cadmium compounds, $RETMg_2$ and $RETCd_2$ have been reported [4–7]. This is an astonishing result, given that they contain rare networks of tetrahedraly coordinated magnesium and cadmium atoms.

Even more interesting is the result that MgCuAl₂-type compounds are formed with two different alkaline

earth elements, as shown by the synthesis of $AETMg_2$ (AE = Ca, Sr; T = Pd, Ag, Pt, Au) and isotypic $SrTCd_2$ compounds [8]. To give an example, the magnesium atoms in $CaAuMg_2$ build up the orthorhombically distorted tetrahedral network, and the calcium atoms fill the typical cation position. Thus one observes clear segregation of magnesium and the heavier alkaline earth elements into different substructures without formation of a solid solution. This is similar to the situation in the calcium-rich compounds Ca_4AgMg and Ca_4AuMg [9].

The MgCuAl₂-type phases so far have been observed for valence electron concentrations in the range of 16-20 [10]; however, also size criteria play an important role. During our systematic phase analytical studies of the AE-T-Mg and AE-T-Cd systems with the heavier alkaline earth elements we obtained the new phases $CaTMg_2$ and $CaTCd_2$ (T = Rh, Pd, Pt). These calcium compounds crystallize with the closely related

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YPd₂Si type [11], a ternary ordered version of the cementite structure, Fe₃C [12]. The synthesis, structure, and bonding peculiarities are reported herein.

Experimental

Synthesis

Starting materials for the syntheses of the $CaTMg_2$ and $CaTCd_2$ (T=Rh, Pd, Pt) samples were granules of distilled calcium (Johnson Matthey, > 99.5%), palladium tear drops, rhodium and platinum sponge (Degussa-Hüls, > 99.9% or Allgemeine Gold- und Silberscheideanstalt, Pforzheim, > 99.9%), a magnesium rod (Johnson Matthey, Ø 16 mm, > 99.5%; the surface of the rod was first cut on a turning lathe in order to remove surface impurities), and a cadmium rod (Johnson Matthey, > 99.9%).

For the preparation of $CaTMg_2$ (T = Rh, Pd, Pt) pieces of calcium, the noble metal sponge or pieces and turnings of the magnesium rod were weighed in the ideal 1:1:2 atomic ratios and arc-welded [13] in small niobium tubes under an

argon pressure of *ca.* 700 mbar. The argon was purified over titanium sponge (900 K), silica gel, and molecular sieves.

The niobium ampoules were sealed in a quartz tube for protection against oxidation and then annealed within 2 h to 1273 K in a muffle furnace. After keeping the temperature for 2 h it was lowered to 873 K within 5 h, kept at 873 K for another 30 h, followed by radiative heat loss. The samples were separated mechanically from the ampoules. No reactions with the ampoule material were evident. The polycrystalline samples are sensitive to moisture and were kept in Schlenk tubes.

The $CaTCd_2$ samples (T=Pd, Pt) were synthesized by a two-step procedure. First calcium, the noble metal and cadmium were weighed in 1:1:2.1 atomic ratios and arcwelded in small niobium tubes. The latter were placed in a water-cooled sample chamber of a high-frequency furnace (Hüttinger Elektronik, Freiburg, type TIG 1.5/300) [14] under flowing argon and annealed at 1300 K for about ten minutes, followed by cooling to 800 K within 30 min. Finally the samples were annealed for two hours at that tempera-

Compound	Structure	a (pm)	b (pm)	c (pm)	$V (\text{nm}^3)$	Reference
	type					
Ca ₃ Pd	Fe ₃ C	769.9	993.7	669.1	0.5119	[44]
CaRhMg ₂	YPd_2Si	787.0(1)	783.8(1)	571.0(1)	0.3522	this work
CaPdMg ₂	YPd_2Si	792.2(2)	803.4(2)	572.0(1)	0.3641	this work
CaPtMg ₂	YPd_2Si	788.5(2)	792.7(2)	566.5(1)	0.3541	this work
CaPdCd ₂	YPd ₂ Si	794.6(2)	809.5(3)	554.7(2)	0.3568	this work
CaPtCd ₂	YPd_2Si	796.6(3)	809.9(2)	552.6(2)	0.3565	this work

Table 1. Lattice parameters (Guinier powder data) of Ca_3Pd and compounds $CaTMg_2$ and $CaTCd_2$ (T=Rh, Pd, Pt). Standard deviations are given in parentheses.

Empirical formula	Ca _{0.94(1)} PdMg _{2.06(1)}	CaPdCd ₂
Formula weight, g mol ⁻¹	194.15	371.28
	194.13	3/1.26
Unit cell dimension (Guinier powder data)	702.2(2)	704 ((2)
a, pm	792.2(2)	794.6(2)
b, pm	803.4(2)	809.5(3)
c, pm	572.0(1)	554.7(2)
Cell volume V , nm ³	0.3641	0.3568
Calculated density, g cm ⁻³	3.54	6.91
Crystal size, μ m ³	$20 \times 80 \times 80$	$50 \times 80 \times 110$
Transmission ratio (min / max)	0.349 / 0.799	0.332 / 0.580
Absorption coefficient, mm ⁻¹	6.5	17.9
F(000), e	358	648
ω range; increment, deg	-60-190; 0.3	0-180; 1
Integration parameters A / B / EMS	6.7 / -6.2 / 0.012	11.6 / 1.7 / 0.011
θ range for data collection, deg	4-46	4-35
Range in hkl	$\pm 15, \pm 16, \pm 11$	$\pm 12, \pm 12, \pm 8$
Total no. of reflections	35970	5338
Independent reflections $/ R_{int}$	1621 / 0.0581	819 / 0.0229
Reflections with $I > 2\sigma(I)/R_{\sigma}$	1191 / 0.0240	758 / 0.0131
Data / ref. parameters	1621 / 24	819 / 23
Goodness-of-fit on F^2	0.997	1.082
$R1 / wR2$ for $I > 2\sigma(I)$	0.0270 / 0.0653	0.0146 / 0.0295
R1 / wR2 for all data	0.0348 / 0.0663	0.0179 / 0.0301
Extinction coefficient	0.038(2)	0.0195(4)
Largest diff. peak / hole, e Å ⁻³	1.25 / -1.30	0.78 / -0.93

Table 2. Crystal data and structure refinement for $Ca_{0.94(1)}PdMg_{2.06(1)}$ and $CaPdCd_2$, space group Pnma, Z=4.

ture, followed by quenching. The temperature was controlled through a Sensor Therm Methis MS09 pyrometer with an accuracy of $\pm 30\,\mathrm{K}$. These niobium ampoules were sealed in quartz tubes for protection against oxidation. Finally the ampoules were annealed in a muffle furnace at 823 K for one week. The CaTCd2 samples are only weakly moisture sensi-

EDX data

Semiquantitative EDX analyses of the single crystals studied on the diffractometers were carried out in variable pressure mode with a Zeiss EVO® MA10 scanning electron microscope with wollastonite, palladium, magnesium oxide, and cadmium as standards. The experimentally observed average composition of 26 ± 1 at.-% Ca : 25 ± 1 at.-% Pd : 49 ± 1 at.-% Cd was close to the ideal one. For the magnesium-containing crystal both the refined (23.5 at.-% Ca: 25 at.-% Pd: 51.5 at.-% Mg) and the EDX analyses $(21 \pm 2 \text{ at.-}\% \text{ Ca} : 22 \pm 2 \text{ at.-}\% \text{ Pd} : 57 \pm 2 \text{ at.-}\% \text{ Mg}) \text{ showed}$ a slightly higher magnesium content, indicating a small range of homogeneity. The standard uncertainties account for the irregular shape of the crystal surfaces. No impurity elements were detected.

X-Ray diffraction

Basic characterization of the CaTMg2 and CaTCd2 (T = Rh, Pd, Pt) samples was carried out by powder X-ray diffraction using the Guinier technique: imaging plate detector, Fujifilm BAS-1800, $CuK_{\alpha 1}$ radiation, and α -quartz (a = 491.30, c = 540.46 pm) as an internal standard. The orthorhombic lattice parameters (Table 1) were deduced from least-squares refinements of the Guinier powder data. Correct indexing was ensured by comparison of the experimental patterns with calculated ones [15].

Crystal fragments of CaPdMg2 and CaPdCd2 were selected from the crushed samples annealed in the muffle furnaces. The crystals were first characterized on a Buerger camera (using white Mo radiation) to check their quality. Intensity data of a suitable CaPdCd2 crystal were collected on a Stoe IPDS-II image plate system (graphitemonochromatized Mo K_{α} radiation; $\lambda = 71.073 \text{ pm}$) in oscillation mode. The CaPdMg2 data set was measured on a Stoe Stadi Vari diffractometer equipped with a Mo micro focus source and a Pilatus detection system and scaled subsequently according to the Gaussian-shaped profile of the X-ray source. Numerical absorption corrections were applied to both data sets. Details about the data collections and the crystallographic parameters are summarized in Table 2.

Structure refinements

The diffractometer data sets of CaPdMg2 and CaPdCd2 showed primitive orthorhombic lattices, and the systematic extinctions were compatible with space group Pnma. The starting atomic parameters were deduced from Direct Methods with SHELXS-97 [16, 17], and the two structures were refined with anisotropic displacement parameters for all atoms with SHELXL-97 (full-matrix least-squares on F_0^2 [18, 19]. The refinements readily revealed isotypism with YPd₂Si [11], and in the following cycles the setting of the silicide was used. The occupancy parameters were refined in separate series of least-squares cycles for both data sets. The sites of CaPdCd₂ were fully occupied within two standard deviations, and the ideal composition was assumed again in the following cycles. The magnesium-based crystal showed too low scattering power for the 4c calcium site, indicating a small degree of mixed occupancy with magnesium. This mixed occupancy was refined as a least-squares variable in the following cycles, leading to the composition $Ca_{0.94(1)}PdMg_{2.06(1)}$ for the investigated crystal. The final difference Fourier syntheses revealed no residual peaks. The refined atomic positions, displacement parameters, and interatomic distances are given in Tables 3-5.

Atom	Wyckoff site	х	у	Z	$U_{ m eq}$
$\overline{\text{Ca}_{0.94(1)}\text{PdMg}_2}$	2.06(1)				
Caa	4 <i>c</i>	0.03241(7) 0.033	1/4	0.10005(10) 0.102	221(2)
Pd	4c	0.37578(3) 0.377	1/4	0.88544(4) 0.884	215(1)
Mg	8 <i>d</i>	0.16774(10) 0.169	0.05683(9) 0.056	0.60470(14) 0.604	223(1)
CaPdCd ₂					
Ca	4c	0.03320(8)	1/4	0.10208(10)	140(1)
Pd	4c	0.37781(3)	1/4	0.88151(4)	117(1)
Cd	8 <i>d</i>	0.17207(2)	0.05570(2)	0.59466(3)	135(1)

Table 3. Atomic coordinates and anisotropic displacement parameters (pm²) for $\begin{array}{c} Ca_{0.94(1)}PdMg_{2.06(1)} \quad \text{and} \\ CaPdCd_2. \quad \textit{U}_{eq} \quad \text{is} \quad \text{defined} \end{array}$ as one third of the trace of the orthogonalized $U_{\rm ij}$ tensor. The parameters of idealized CaPdMg2 obtained by geometry optimization are listed in italics.

^a The calcium site shows mixed occupancy of 94(1) % Ca + 6(1) % Mg.

Atom	U_{11}	U_{22}	U_{33}	U_{23}	U_{13}	U_{12}
Ca _{0.94(1)} PdN	Ig _{2.06(1)}					
Ca ^a	208(2)	248(3)	207(3)	0	11(2)	0
Pd	200(1)	216(1)	229(1)	0	-10(1)	0
Mg	232(3)	200(3)	236(3)	-12(2)	-11(2)	-2(2)
CaPdCd ₂						
Ca	125(2)	183(3)	112(2)	0	5(2)	0
Pd	116(1)	113(1)	121(1)	0	-6(1)	0
Cd	133(1)	116(1)	156(1)	-27(1)	-20(1)	2(1)

Table 4. Anisotropic displacement parameters (pm^2) for $Ca_{0.94(1)}PdMg_{2.06(1)}$ and $CaPdCd_2$.

Further details of the crystal structure investigation may be obtained from Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (fax: +49-7247-808-666; E-mail: crysdata@fiz-karlsruhe.de, http://www.fiz-karlsruhe.de/request_for_deposited_data.html) on quoting the deposition number CSD-425508 (Ca_{0.94} PdMg_{2.06}) and CSD-425509 (CaPdCd₂).

Computational methods

Two computational methods within the DFT [20, 21] were used in a complementary manner. The Vienna ab initio simulation package (VASP) code [22, 23] allows geometry optimization and total energy calculations. For this we use the projector augmented wave (PAW) method [23, 24], built within the generalized gradient approximation (GGA) scheme following Perdew, Burke and Ernzerhof (PBE) [25]. Also preliminary calculations with local density approximation LDA [26] led to an underestimated volume versus the experiment. The conjugate-gradient algorithm [27] is used in this computational scheme to relax the atoms. The tetrahedron method with Blöchl corrections [24] as well as a Methfessel-Paxton [28] scheme were applied for both geometry relaxation and total energy calculations. Brillouin zone (BZ) integrals were approximated using the special k-point sampling. The optimization of the structural parameters was performed until the forces on the atoms were less than $0.02 \,\text{eV}\,\text{Å}^{-1}$ and all stress components less than 0.003 eV Å⁻³. The calculations are converged at an energy cut-off of 400 eV for the plane-wave basis set with respect to the k-point integration with a starting mesh of $4 \times 4 \times 4$ up to $8 \times 8 \times 8$ for best convergence and relaxation to zero strains. The calculations are scalar-relativistic and assume spin-degenerate total

Then all-electron calculations with the GGA were carried out for a full description of the electronic structure and the properties of chemical bonding, using the full potential scalar-relativistic augmented spherical wave (ASW) method [29, 30]. In the minimal ASW basis set, we chose the outermost shells to represent the valence states, and the matrix elements were constructed using partial waves up to

 $l_{\rm max}+1=3$ for Pd and $l_{\rm max}+1=2$ for Ca and Mg. Self-consistency was achieved when charge transfers and energy changes between two successive cycles were below 10^{-8} and $10^{-6}\,{\rm eV}$, respectively. BZ integrations were performed using the linear tetrahedron method within the irreducible wedge. Besides the site projected density of states, we have discussed qualitatively the pair interactions based on the overlap population (COOP) [31]. In the plots, positive, negative, and zero COOP indicate bonding, anti-bonding, and non-bonding interactions, respectively.

Table 5. Interatomic distances (pm) in the structures of $Ca_{0.94}PdMg_{2.06}$ and $CaPdCd_2$. All distances of the first coordination spheres are listed. Standard deviations are all smaller or equal than 0.2~pm.

Ca _{0.9}	₄ PdN	Mg _{2.06}		CaPdCd	\mathbf{l}_2	
Ca ^a :	1	Pd	298.4	Ca: 1	Pd	295.3
	1	Pd	304.2	1	Pd	299.9
	1	Pd	319.4	1	Pd	311.9
	2	Mg	338.3	2	Cd	334.0
	2	Mg	340.4	2	Cd	340.7
	2	Mg	342.3	2	Cd	340.8
	2	Mg	344.8	2	Cd	340.8
	2	Mg	348.2	2	Cd	345.0
Pd:	2	Mg	277.6	Pd: 2	Cd	277.1
	2	Mg	278.6	2	Cd	277.1
	2	Mg	278.7	2	Cd	282.1
	1	Ca	298.4	1	Ca	295.3
	1	Ca	304.2	1	Ca	299.9
	1	Ca	319.4	1	Ca	311.9
Mg:	1	Pd	277.6	Cd: 1	Pd	277.1
	1	Pd	278.6	1	Pd	277.1
	1	Pd	278.7	1	Pd	282.1
	1	Mg	305.5	1	Cd	306.5
	1	Mg	310.4	1	Cd	314.6
	2	Mg	327.3	2	Cd	316.9
	1	Ca	338.3	1	Ca	334.0
	1	Ca	340.4	1	Ca	340.7
	1	Ca	342.3	1	Ca	340.8
	1	Ca	344.8	1	Ca	340.8
	1	Ca	348.2	1	Ca	345.0

^a The calcium site shows mixed occupancy of 94(1)% Ca + 6(1)% Mg.

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Discussion

Crystal chemistry

The new calcium compounds CaTMg2 and CaTCd2 (T = Rh, Pd, Pt) crystallize with the orthorhombic YPd₂Si-type structure (space group *Pnma*) [11], a ternary ordered variant of the well known cementite type, Fe₃C [12]. So far, the YPd₂Si type has been observed for REPd2Si and REPt2Si silicides [11, 32-34], REPd₂Ge germanides [35-39], UAu₂Al [40], CePd₂Ga [41], YPd₂Ga [42], and the phosphide MgNi₂P [43]. In these intermetallics, the small silicon, germanium, aluminum, gallium, or phosphorus atoms fill trigonal-prismatic sites. The inverse ordering pattern is realized in the magnesium and cadmium compounds reported herein. The transition metal atoms fill trigonal-prismatic sites formed by the magnesium, respectively cadmium atoms. As an example for the new ternary phases we discuss the CaPdMg₂ structure.

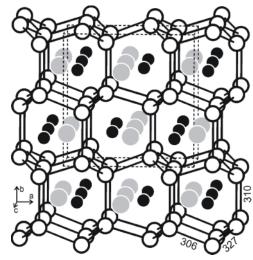
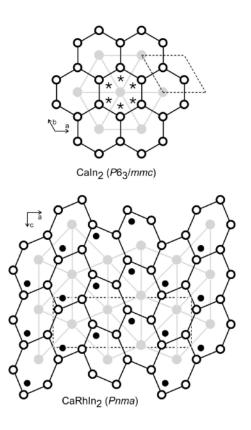


Fig. 1. View of the CaPdMg $_2$ structure approximately along the crystallographic c axis. Calcium, palladium and magnesium atoms are drawn as medium grey, black filled and open circles, respectively. The three-dimensional lonsdaleite-related magnesium network is emphasized. Relevant Mg–Mg distances are indicated.



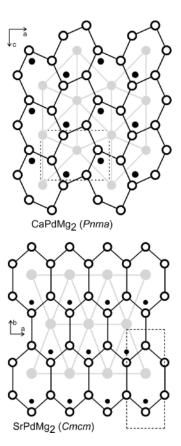


Fig. 2. View upon one puckered indium or magnesium layer and its adjacent calcium (strontium) layer in the structures of CaIn₂, CaRhIn2, CaPdMg2, and SrPdMg₂. Calcium (strontium), rhodium (palladium) indium and (magnesium) atoms are drawn as medium grey, black filled and open circles, respectively. The three ternary structures arise from different insertion of the transition metal atoms into distorted CaIn2-type subcells as emphasized by grey lines (the ideal interstitial sites in CaIn2 are marked with asterisks; within one hexagon of the ternaries only one position is filled).

The single-crystal X-ray diffraction data for CaPdMg2 revealed a small homogeneity range $Ca_{1-x}PdMg_{2+x}$. In this context it is interesting to note that binary Ca₃Pd [44] with cementite structure is known, while related binaries are not known for the corresponding Ca-Rh and Ca-Pt systems. The refined composition of the single crystal was Ca_{0.94}PdMg_{2.06}. One can thus assume an extension towards higher and lower calcium contents as well. Detailed preparative work on these solid solutions is in progress. From Table 1 it becomes clear that substitution of calcium by the smaller magnesium atoms leads to anisotropic changes of the lattice parameters, i. e. the a parameter is increasing, while b and c are decreasing. Towards the magnesium-rich side one can expect a change in structure type since binary Mg₃Pd [45] crystallizes with the hexagonal Cu₃P type.

A view of the CaPdMg₂ structure approximately along the *c* axis is presented in Fig. 1. Similar to the MgCuAl₂-type phases [8, 10], for CaPdMg₂ we also observe formation of a pronounced lonsdaleite-related (hexagonal diamond) magnesium substructure. The Mg–Mg distances range from 306 to 327 pm, in close agreement with the average Mg–Mg distance of 320 pm in *hcp* magnesium [46]. We can thus expect substantial Mg–Mg bonding. In the magnesium substructures of the MgCuAl₂-type phases one observes much longer Mg–Mg distances and larger anisotropy in Mg–Mg bonding [8].

The shortest distances within the CaPdMg2 structure occur between the palladium and magnesium atoms. Within the distorted trigonal PdMg6 prisms the Pd-Mg distances range from 278 to 279 pm, slightly longer than the sum of the covalent radii of 264 pm [47]. Together the palladium and magnesium atoms build up a complex three-dimensional [PdMg₂] network which leaves cavities for the calcium atoms. The latter have coordination number 13 with three palladium and ten magnesium atoms in their coordination shell. The calcium atoms are bonded to the [PdMg₂] network through short Ca-Pd bonds (298 – 319 pm), close in length to the sum of the covalent radii of 302 pm. A similar bonding situation occurs also in CaPdIn₂ (294-316 pm Ca-Pd) [48] and CaPdSn₂ (305 pm) [49].

At large, one observes similar structural features for the MgCuAl₂- and YPd₂Si-type *AET*Mg₂ compounds. But what is the difference between the two types? The *CaMg*₂ substructures of both families resemble the hexagonal Zintl phase CaIn₂. This relationship was first reported for the indides $SrTIn_2$ (T = Rh, Pd, Ir, Pt) [3]. In Fig. 2 we present cut-outs of the CaIn₂ [50], CaPdMg₂, CaRhIn₂ [51], and SrPdMg₂ structures [8]. For each compound one layer of puckered hexagons and one Ca, CaPd, CaRh, and SrPd layer are shown. We start the structural comparison with the unfilled hexagonal CaIn₂ type. This structure leaves six possible sites (marked by asterisks) for the incorporation of the transition metal atom. For geometric reasons, only one of these sites is filled per hexagon, and adjacent sites between two hexagons cannot be occupied simultaneously. With the three structures presented in Fig. 2 one knows three different coloring variants. In all of these compounds the puckered hexagons are elongated as a consequence of the Pd-Mg and Rh-In bonding. The elongated hexagons all point to one direction in SrPdMg₂, while CaPdMg₂ and CaRhIn₂ show simple and double herringbone patterns. The different coloring variants allow for the formation of 1-1-2 compounds for diverse element combinations and valence electron concentrations (VEC). The YPd₂Si-type CaTMg₂ compounds reported herein have low VEC of 15 and 16, while the MgCuAl₂-type compounds show a range from 16 to 20. Further phase analytical work on the Ba-T-Mg and Ba-T-Cd systems is under way in order to elucidate potential 1-1-2 compounds with much larger cations.

Geometry optimization results and discussions

In Table 3 the calculated positional parameters for CaPdMg₂ show good agreement with the experimental data but tend to provide a slightly lower volume. We further derive the equilibrium zero pressure parameters from the energy-volume, E(V), equation of state (EOS) with calculations around minima found from geometry optimization. The resulting values are plotted in Fig. 3 using a 3rd order Birch EOS [52] providing E_0 , V_0 , B_0 and B', respectively, as the equilibrium energy, the volume, the bulk modulus and its pressure derivative given in the insert. The fit curves produce a slightly larger volume versus experiment and hence correct the calculated value from geometry optimization because GGA tends to provide larger volumes and a certain underbinding, contrary to the LDA. The bulk modulus which translates the resistance of the compound to isotropic compression has a very low value of $B_0 = 48$ GPa. Compared with the constituent ele-

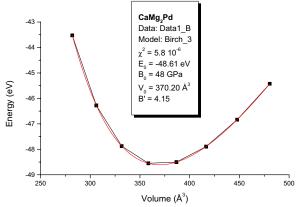


Fig. 3 (color online). CaPdMg₂: Energy-volume curve and fit values from Birch EOS. χ^2 : goodness of fit.

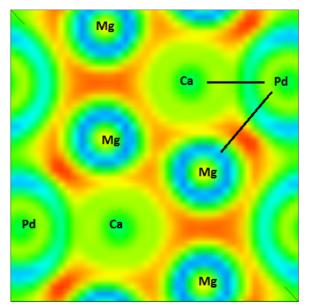


Fig. 4 (color online). CaPdMg₂: Electron localization function ELF projection onto the *ab* plane highlighting the CaPd and Mg-Pd bonds and the localization (red spots) along Mg-Mg. Also notice the dominant green ELF indicating free electron-like behavior (see text for color definitions).

ments in GPa units, $B_0(Pd) = 180$, $B_0(Mg) = 45$ GPa and $B_0(Ca) = 17$, it cannot be the consequence of their averaging (96 GPa) but should rather arise from the nature of its electronic structure.

A first impression is given by the mapping of the electron localization function (ELF) [53]. ELF is a normalized (unit-less) function: 0 < ELF < 1, with ELF = 0 indicating zero localization (blue contours),

ELF = 1 for strong localization (red contours) and ELF = 1/2 value for free electron gas behavior (green contours). The plot shown in Fig. 4 is for a plane cross-

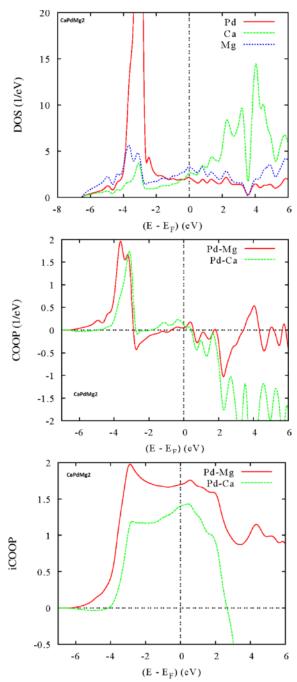


Fig. 5 (color online). Electronic structure and chemical bonding in $CaPdMg_2$.

ing the Mg atoms. The departure of electrons from electropositive Ca and Mg is shown by blue areas (zero localization spots), not seen for Ca because, the atoms are not in the same plane. Nevertheless the dominating green color indicates a free electron itinerant (s, p) behavior, despite the presence of low-energy lying Pd d states in the valence band. The highlighted Mg–Pd and Ca–Pd bonds show rather different features with localized ELF (red-like spots) between Mg and Pd and only green contours between Ca and Pd. The red spots in between Mg and neighboring Pd atoms are a relevant feature and indicate bonding properties discussed in the next section.

We further analyze the charge density issued from the self consistent calculations using the AIM (atoms in molecules theory) approach [54] developed by Bader who devised an intuitive way of splitting molecules into atoms as based purely on the electronic charge density. Typically in chemical systems, the charge density reaches a minimum between atoms and this is a natural region to separate them from each other. Such an analysis does not constitute a tool for evaluating absolute ionizations. Bader's analysis is done using a fast algorithm operating on a charge density grid. The program [55] reads in charge densities obtained from high precision VASP calculations and outputs the total charge associated with each atom. Further, in order to include core electrons within the PAW method in VASP, use is made of the LAECHG = .TRUE. parameter in the 'INCAR' control file followed by operating a summation of core and valence charge densities using a simple perl program [55]. The results of computed charge changes (ΔQ) are such that they lead to neutrality when the respective multiplicities are accounted for; the obtained values are:

$$\Delta Q(Mg) = +0.236$$
; $\Delta Q(Ca) = +0.919$ and $\Delta Q(Pd) = -1.391$.

Due to the large electronegativity difference between Pd ($\chi=2.2$) on one hand and Mg ($\chi=1.3$) and Ca ($\chi=1.0$) on the other, charge transfer is observed from Ca and Mg (positive charge) towards Pd (negative charge). However, the more pronounced electropositive character of calcium leads to a more extensive ionization (close to 1) while Mg is much less ion-

ized. This corroborates the ELF plots described above and indicates a rather iono-covalent Mg–Pd compared to a more ionic Ca–Pd bond. The charge distribution and the course of the bond strengths resemble those in CaAuIn₂ [48].

Electronic structure and chemical bonding

The electronic density of states and the chemical bonding are analyzed based on experimental data and assuming a spin-degenerate total spin (NSP) configuration. At self consistent convergence the charge transfer follows the trends observed above.

Figure 5 shows the site projected density of states (PDOS) for CaPdMg₂ accounting for site multiplicities and the total number of FU per cell. Along the x axis the energy is with respect to the Fermi level E_F . The valence band is dominated by low energy filled Pd-d states with contributions from Mg and Ca PDOS underneath, pointing to quantum mixing between them. The twice larger number of Mg vs. Ca leads to higher PDOS, but most of the Ca PDOS are empty and found above E_F in the conduction band. This is in agreement with the larger ionization of Ca as detailed above. The crossing of E_F by low intensity itinerant broad PDOS from all components signals a metal with mainly s character orbitals.

The quantum mixing between Pd and Ca/Mg can be further assessed using the COOP criterion of chemical bonding. The relative strengths of chemical bonding can be discussed within the valence band below $E_{\rm F}$ with one atom of each kind for the sake of comparing magnitudes. The Mg–Pd bond is found with larger intensity than the Ca–Pd bond with main contributions below the Pd-d PDOS. This follows from the larger Ca–Pd vs. Mg–Pd distances (298 vs. 278 pm). Low magnitude antibonding COOP (<0) are observed around -2 eV, but they do not change the trend of stronger Mg–Pd bonding. This is shown with the integrated iCOOP criterion in $3^{\rm rd}$ panel to be interpreted by the principle: the larger the area below the curve, the stronger the bonding.

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