

CONCENTRATION COMPACTNESS OF A LANDAU-LIFSCHITZ TYPE ENERGY WITH RADIAL STRUCTURE

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ABSTRACT. This paper is concerned with the asymptotic behavior of a p -Landau-Lifschitz type functional with radial structure as parameter goes to zero. We study the concentration compactness properties in the case of $p \geq n$ where n is the dimension. By analyzing the functional globally, we show that the singularity of p -Landau-Lifschitz energy concentrates on the origin.

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1. Introduction

Let $G \subset \mathbf{R}^n$ ($n \geq 2$) be a bounded and simply connected domain with smooth boundary ∂G . Denote $S^{n-1} = \{x \in \mathbf{R}^{n+1} : x_1^2 + x_2^2 + \cdots + x_n^2 = 1, x_{n+1} = 0\}$ and $S^n = \{x \in \mathbf{R}^{n+1} : |x| = 1\}$. We sometimes write the vector-valued function $u = (u_1, u_2, \dots, u_n, u_{n+1}) = (u', u_{n+1})$. Let $g = (g', 0)$ be a smooth map from ∂G into S^{n-1} satisfying $\deg(g', \partial G) = d > 0$. Consider the energy functional

$$E_\varepsilon(u, G) = \frac{1}{p} \int_G |\nabla u|^p dx + \frac{1}{2\varepsilon^p} \int_G u_{n+1}^2 dx,$$

where the small parameter $\varepsilon > 0$ and $p \geq n$. By the direct method in the calculus of variations, the minimizer u_ε exists in the space $W' = \{u \in W^{1,p}(G, S^n) :$

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$u|_{\partial G} = g\}$. Clearly, the minimizer is a weak solution of the system

$$-\operatorname{div}(|\nabla u|^{p-2}\nabla u) = u|\nabla u|^p + \frac{1}{\varepsilon^p} [u(u_{n+1})^2 - u_{n+1}e_{n+1}] \quad \text{in } G, \quad (1.1)$$

where $e_{n+1} = (0, 0, \dots, 0, 1)$. Maybe (1.1) has no classical solution since it is degenerate. Therefore, following Uhlenbeck's idea, we consider the regularized functional

$$E_\varepsilon^\tau(u, G) = \frac{1}{p} \int_G (|\nabla u|^2 + \tau)^{p/2} dx + \frac{1}{2\varepsilon^p} \int_G u_{n+1}^2 dx.$$

Similar to the derivation of [8: (4.2) and (4.8)], when $p \geq n$, the regularized minimizer u_ε^τ of $E_\varepsilon^\tau(u, G)$ satisfies

$$\lim_{\tau \rightarrow 0} u_\varepsilon^\tau = u_\varepsilon, \quad \text{in } W^{1,p}(G), \quad (1.2)$$

$$\lim_{\tau \rightarrow 0} u_\varepsilon^\tau = u_\varepsilon, \quad \text{in } C_{loc}^{1,\beta}(G \setminus \{a_j\}_{j=1}^N), \quad \text{for some } \beta \in (0, 1/2). \quad (1.3)$$

where u_ε is a minimizer of $E_\varepsilon(u, G)$, and $a_j (j = 1, 2, \dots, N)$ are singularities of a p -harmonic map with boundary value g' . We call u_ε a *regularized minimizer* of $E_\varepsilon(u, G)$.

In the 2-dimensional case, the functional $E_\varepsilon(u, G)$ was introduced in the study of some simplified model of high-energy physics, which controls the statics of planner ferromagnets and antiferromagnets (cf. [6], [11] and [12]). The asymptotic behavior of minimizers of $E_\varepsilon(u, G)$ has been considered in [4]. One of the main results is the following proposition.

PROPOSITION 1.1. ([4: Theorem 4.2]) *Assume $p = n = 2$, u_ε is a minimizer of $E_\varepsilon(u, G)$, and $a_1, a_2, \dots, a_d$ are singularities of a harmonic map with boundary value g' . Then as $\varepsilon \rightarrow 0$, there exists a subsequence ε_k such that*

$$\lim_{\varepsilon_k \rightarrow 0} \frac{u_{\varepsilon_k}^2}{2\varepsilon_k^2} = \frac{\pi}{2} \sum_{j=1}^d \delta_{a_j}, \quad \text{in the weak } * \text{ topology of } C(\overline{G}), \quad (R1)$$

$$\lim_{\varepsilon_k \rightarrow 0} \frac{|\nabla u_{\varepsilon_k}|^2}{|\log \varepsilon_k|} = 2\pi \sum_{j=1}^d \delta_{a_j}, \quad \text{in the weak } * \text{ topology of } C(\overline{G}), \quad (R2)$$

where δ_{a_j} is the Dirac mass at a_j .

If the term $\frac{u_{n+1}^2}{\varepsilon^p}$ is replaced by $\frac{(1-|u|^2)^2}{2\varepsilon^p}$ and S^n is replaced by $\mathbf{R}^n$, the functional is the Ginzburg-Landau functional. Many researches studied the case $p = n = 2$ (cf. [2] and the references therein). In particular, in [2: Theorems VII.2 and VII.3] derived the similar results to (R1) and (R2). Paper [10] generalized those results to the case of $p > n = 2$. Other work on $p \neq n$ can be seen in [1] and [16]. Nineteen open problems were introduced in [2]. When $p = n > 2$, the asymptotic behavior of minimizers was in [2: Problem 17], which

was studied in [3], [5] and [14]. Their work shows that the study of minimizers of n -Ginzburg-Landau functional is connected tightly with the corresponding properties of the n -harmonic map. Afterwards, paper [9] obtained (R1) and (R2) when $p = n > 2$.

In this paper, we will study the Landau-Lifschitz type functional $E_\varepsilon(u, G)$ with radial structure. Let $G = B = \{x \in \mathbf{R}^n : |x| < 1\}$. The minimizer u_ε of $E_\varepsilon(u, B)$ in the space

$$W = \left\{ u(x) = \left(\frac{x}{|x|} \sin f(r), \cos f(r) \right) \in W^{1,p}(B, S^n) : f(1) = \frac{\pi}{2} \right\},$$

is named *radial minimizer*, where $r = |x|$. Similar to Remark in [7: p. 68], $f_\varepsilon(r)$ also belongs to $\{f \in C[0, 1] : f(0) = 0\}$.

When $p > n = 2$, Paper [7] studied the asymptotics of the radial minimizer. Those results can be extended to the higher dimensions without any obstacle. First, in [7: Proposition 2.1] implies that there exists $C > 0$ which is independent of ε , such that

$$\varepsilon^{p-n} \int_B |\nabla u_\varepsilon|^p dx \leq C, \quad (1.4)$$

$$\frac{1}{\varepsilon^n} \int_B |u_{\varepsilon n+1}|^2 dx \leq C. \quad (1.5)$$

Next, similar to [7: Proposition 3.2], for any compact subset K of $B \setminus \{0\}$, there exists a constant $C > 0$ which is independent of ε , such that

$$E_\varepsilon(u_\varepsilon, K) \leq C. \quad (1.6)$$

In addition, [7: Theorem 1.1] shows that, for any $\gamma \in (0, 1)$, there exists a constant $h = h(\gamma) > 0$ which is independent of ε , such that

$$|u'_\varepsilon(x)|^2 \geq 1 - \gamma^2, \quad \text{for } x \in B \setminus B(0, h\varepsilon). \quad (1.7)$$

In this paper, we establish the properties (R1) and (R2) for $p \geq n > 2$.

THEOREM 1.2. *Assume $p > n > 2$, u_ε is a regularized radial minimizer. Then as $\varepsilon \rightarrow 0$, there exists a subsequence ε_k such that*

$$\frac{1}{2\varepsilon_k^n} |u_{\varepsilon_k n+1}|^2 \rightarrow L\delta_o, \quad \text{weakly star in } C(\overline{B}). \quad (1.8)$$

$$\varepsilon_k^{p-n} |\nabla u_{\varepsilon_k}|^p \rightarrow \frac{np}{p-n} L\delta_o, \quad \text{weakly star in } C(\overline{B}). \quad (1.9)$$

Here δ_o is the Dirac mass at the origin, and the positive constant L satisfies

$$\begin{aligned} & \frac{(n-1)^{p/2}|\partial B|}{np} \sup_{\gamma \in (0,1)} (1-\gamma^2)^{p/2} h^{n-p}(\gamma) \\ & \leq L \leq (1 - \frac{n}{p}) \min_W E_1(u, B) + \frac{(n-1)^{p/2}}{p^2} |\partial B|, \end{aligned} \quad (1.10)$$

where $h(\gamma)$ is a positive constant in (1.7).

THEOREM 1.3. Assume $p = n > 2$, u_ε is a regularized radial minimizer. Then as $\varepsilon \rightarrow 0$,

$$\frac{|\nabla u_\varepsilon|^n}{|\log \varepsilon|} \rightarrow (n-1)^{\frac{n}{2}} |\partial B| \delta_o, \quad \text{weakly star in } C(\overline{B}). \quad (1.11)$$

$$\frac{|u_{\varepsilon n+1}|^2}{2\varepsilon^n} \rightarrow \frac{(n-1)^{\frac{n}{2}} |\partial B|}{n^2} \delta_o, \quad \text{weakly star in } C(\overline{B}). \quad (1.12)$$

Remark 1. We conjecture that L in the limits of (1.8) and (1.9) should be independent of the subsequence. If it is true, then (1.8) and (1.9) hold not only for ε_k , but also for the whole u_ε . However, this conjecture seems difficult to prove. How to deal with the absence of the conformal transformation of $\int |\nabla u|^p dx$ when $p \neq n$ is also an open problem.

Remark 2. Theorem 1.2 shows that the energy $\varepsilon^{p-n} E_\varepsilon(u_\varepsilon, B)$ concentrates on the origin. In addition, if $p = n$ in (1.10), then $L = \frac{(n-1)^{n/2}}{n^2} |\partial B|$, (1.8) becomes (1.12). Moreover, if $p = n = 2$ in (1.10), then $L = \frac{\pi}{2}$, and (1.8), (1.11) and (1.12) are identical with (R1) and (R2).

Remark 3. When $p < n$, the energy $E_\varepsilon(u_\varepsilon, B) \leq E_\varepsilon(u_*, B) = C$, where $u_* = (\frac{x}{|x|}, 0)$. The concentration compactness of this p-Landau-Lifschitz energy like (R1) and (R2) does not happen.

2. Proof of Theorem 1.2

In this section, we assume $p > n > 2$, and u_ε is the regularized radial minimizer.

Proof of (1.8) and (1.9). In view of (1.4) and (1.5), there exist two Radon measures ω_1 and ω_2 , such that as $\varepsilon \rightarrow 0$,

$$\varepsilon_k^{p-n} |\nabla u_{\varepsilon_k}|^p \rightarrow \omega_1, \quad \text{weakly star in } C(\overline{B}), \quad (2.1)$$

$$\frac{1}{2\varepsilon_k^n} u_{\varepsilon_k n+1}^2 \rightarrow \omega_2, \quad \text{weakly star in } C(\overline{B}), \quad (2.2)$$

for some subsequence ε_k of ε . Sometimes we also denote u_{ε_k} by u_ε if it is not confusing. Furthermore, (1.6) implies that as $\varepsilon \rightarrow 0$,

$$\begin{aligned} \varepsilon^{p-n} \int_K |\nabla u_\varepsilon|^p dx &\rightarrow 0, \\ \frac{1}{2\varepsilon^n} \int_K u_{\varepsilon_{n+1}}^2 dx &\rightarrow 0, \end{aligned}$$

where K is an arbitrary compact subset of $B \setminus \{0\}$. These results lead to $\text{supp}(\omega_i) = \{0\}$ for $i = 1, 2$. Then we can find constants L and L' such that

$$\omega_1 = L'\delta_o, \quad \omega_2 = L\delta_o. \quad (2.3)$$

Next, we shall point out the relation between L and L' .

It is not difficult to see that the radial minimizer u_ε^τ solves the system

$$\begin{aligned} -\text{div}[v^{(p-2)/2} \nabla u] &= u |\nabla u|^2 v^{(p-2)/2} \\ &+ \frac{1}{\varepsilon^p} (u u_{n+1}^2 - u_{n+1} e_{n+1}) \quad \text{in } B, \end{aligned} \quad (2.4)$$

where $v = |\nabla u|^2 + \tau$. Multiplying (2.4) by $x \cdot \nabla u$ and integrating by parts, we can obtain the Pohozaev type identity

$$\begin{aligned} & - \int_{\partial B_R(0)} |x| v^{\frac{p-2}{2}} |\partial_\nu u|^2 ds + \int_{B_R(0)} v^{\frac{p-2}{2}} |\nabla u|^2 dx \\ & - \frac{n}{p} \int_{B_R(0)} v^{\frac{p}{2}} dx + \frac{1}{p} \int_{\partial B_R(0)} |x| v^{\frac{p}{2}} ds \\ & = - \frac{1}{2\varepsilon^p} \int_{\partial B_R(0)} |x| u_{n+1}^2 ds + \frac{n}{2\varepsilon^p} \int_{B_R(0)} u_{n+1}^2 dx \end{aligned} \quad (2.5)$$

for any $R \in (0, 1]$. Hereafter, we denote f_ε by f . By (1.6) and the mean value theorem, there exists $\sigma \in (1/4, 1/2)$ such that

$$r^{n-1} [(f_r)^2 + r^{-2} (n-1) \sin^2 f]^{p/2} \Big|_{r=\sigma} + \frac{r^{n-1}}{\varepsilon^p} \cos^2 f \Big|_{r=\sigma} \leq C. \quad (2.6)$$

Then, we take $R = \sigma$ in (2.5) and multiply it by ε^{p-n} . Letting $\tau \rightarrow 0$ and using (1.2) and (1.3), we have

$$\begin{aligned} & -\varepsilon^{p-n} \sigma^n [(f_r)^2 + \frac{n-1}{r^2} \sin^2 f]^{p/2} \Big|_{\sigma} \\ & + \left(1 - \frac{n}{p}\right) \varepsilon^{p-n} \int_0^{\sigma} [(f_r)^2 + \frac{n-1}{r^2} \sin^2 f]^{p/2} r^{n-1} dr \\ & + \frac{\sigma^n}{p} \varepsilon^{p-n} [(f_r)^2 + \frac{n-1}{r^2} \sin^2 f]^{p/2} \Big|_{\sigma} \\ & = -\frac{\sigma^n}{2\varepsilon^n} \cos^2 f \Big|_{\sigma} + \frac{n}{2\varepsilon^n} \int_0^{\sigma} \cos^2 f r^{n-1} dr. \end{aligned}$$

Using (2.6), we get

$$\left(1 - \frac{n}{p}\right) \varepsilon^{p-n} \int_{B_{\sigma}(0)} |\nabla u_{\varepsilon}|^p dx - \frac{n}{2\varepsilon^n} \int_{B_{\sigma}(0)} u_{\varepsilon n+1}^2 dx \rightarrow 0 \quad (2.7)$$

as $\varepsilon \rightarrow 0$. Combining this result with (2.1)–(2.3), we obtain

$$L' = \frac{np}{p-n} L.$$

Thus, (1.8) and (1.9) are proved. $\square$

Proof of (1.10).

Step 1. Upper bound.

Similar to the proof of [7: Proposition 2.1], it is easy to derive

$$\varepsilon^{p-n} E_{\varepsilon}(u_{\varepsilon}, B) \leq \frac{(n-1)^{p/2}}{p(p-n)} |\partial B| + \min_W E_1(u, B) + C\varepsilon^{p-n}. \quad (2.8)$$

Here $C > 0$ is independent of ε . On the other hand, (1.8) and (1.9) lead to

$$\lim_{\varepsilon \rightarrow 0} \left[\frac{\varepsilon^{p-n}}{p} |\nabla u_{\varepsilon}|^p + \frac{1}{2\varepsilon^n} u_{\varepsilon n+1}^2 \right] = \frac{p}{p-n} L \delta_o, \quad \text{weakly star in } C(\bar{B}). \quad (2.9)$$

This result, together with (2.8), implies the upper bound of L in (1.10).

Step 2. Lower bound.

From (1.7), we can deduce that, for any $\sigma > 0$, there exists $C > 0$ independent of ε , such that

$$\begin{aligned} \int_{h\varepsilon}^{\sigma} [(f_r)^2 + r^{-2}(n-1) \sin^2 f]^{p/2} r^{n-1} dr & \geq \int_{h\varepsilon}^{\sigma} r^{n-p-1} (n-1)^{p/2} \sin^p f dr \\ & \geq \frac{(n-1)^{p/2}}{p-n} (1 - \gamma^2)^{p/2} h^{n-p} (\gamma) \varepsilon^{n-p} - C(\sigma). \end{aligned} \quad (2.10)$$

Applying (2.7), we obtain that,

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \varepsilon^{p-n} E_\varepsilon(u_\varepsilon, B_\sigma(0)) &= \lim_{\varepsilon \rightarrow 0} \frac{|\partial B| \varepsilon^{p-n}}{p} \int_0^\sigma [(f_r)^2 + r^{-2}(n-1) \sin^2 f]^{p/2} r^{n-1} dr \\ &\quad + \lim_{\varepsilon \rightarrow 0} \frac{|\partial B|}{2\varepsilon^2} \int_0^\sigma \cos^2 f r^{n-1} dr \\ &= \frac{|\partial B|}{n} \lim_{\varepsilon \rightarrow 0} \varepsilon^{p-n} \int_0^\sigma [(f_r)^2 + r^{-2}(n-1) \sin^2 f]^{p/2} r^{n-1} dr. \end{aligned}$$

Inserting (2.10) into this result, we deduce that for any $\eta \in (0, 1)$,

$$\lim_{\varepsilon \rightarrow 0} \varepsilon^{p-n} E_\varepsilon(u_\varepsilon, B_\sigma(0)) \geq \frac{(n-1)^{p/2} |\partial B|}{n(p-n)} (1 - \gamma^2)^{p/2} h^{n-p}(\gamma).$$

Taking the supremum and writing

$$H := \sup_{\gamma \in (0,1)} (1 - \gamma^2)^{p/2} h^{n-p}(\gamma),$$

we have

$$\lim_{\varepsilon \rightarrow 0} \varepsilon^{p-n} E_\varepsilon(u_\varepsilon, B_\sigma(0)) \geq \frac{(n-1)^{p/2} |\partial B|}{n(p-n)} H.$$

Combining this with (2.9), we obtain the lower bound of L in (1.10). $\square$

3. Proof of Theorem 1.3

In this section, we assume $p = n > 2$, and u_ε is the regularized radial minimizer. In view of [13: Theorem 2.7], (1.7) is still true.

PROPOSITION 3.1. *Assume h is the constant in (1.7). Then for any $\sigma > 0$, we can find $C > 0$ which is independent of ε , such that*

$$\int_{B_{h\varepsilon}(0) \cup (B \setminus B_\sigma(0))} |\nabla u_\varepsilon|^n dx \leq C; \quad (3.1)$$

$$\int_{B_\sigma(0) \setminus B_{h\varepsilon}(0)} |\nabla f|^n dx \leq C; \quad (3.2)$$

$$\frac{1}{2\varepsilon^n} \int_B u_{\varepsilon n+1}^2 dx \leq C. \quad (3.3)$$

Proof. Clearly,

$$|\nabla u_\varepsilon|^n \geq |\nabla f|^n + \sin^n f \left| \nabla \frac{x}{|x|} \right|^n. \quad (3.4)$$

According to [13: Lemmas 2.2 and 3.1], we know

$$E_\varepsilon(u_\varepsilon, B) \leq \frac{1}{n}(n-1)^{n/2}|\partial B||\log \varepsilon| + C; \quad (3.5)$$

$$\int_{B_\sigma(0) \setminus B_{h\varepsilon}(0)} \left| \nabla \frac{x}{|x|} \right|^n dx \geq (n-1)^{n/2}|\partial B||\log \varepsilon| - C. \quad (3.6)$$

Eq. (3.5) subtracts (3.6). Then by applying (3.4) we obtain

$$\begin{aligned} \frac{1}{n} \int_{B_{h\varepsilon}(0) \cup (B \setminus B_\sigma(0))} |\nabla u_\varepsilon|^n dx + \frac{1}{n} \int_{B_\sigma(0) \setminus B_{h\varepsilon}(0)} |\nabla f|^n dx + \frac{1}{2\varepsilon^n} \int_B u_{\varepsilon_{n+1}}^2 dx \\ \leq C + \frac{1}{n} \int_{B_\sigma(0) \setminus B_{h\varepsilon}} (1 - \sin^n f) \left| \nabla \frac{x}{|x|} \right|^n dx. \end{aligned}$$

Using in [13: Lemma 3.2], we can deduce (3.1)–(3.3). Proposition 3.1 is proved. $\square$

Proof of (1.11). By virtue of (3.5), there exists a Radon measure μ_1 such that

$$\lim_{\varepsilon_k \rightarrow 0} \frac{|\nabla u_{\varepsilon_k}|^n}{|\log \varepsilon_k|} = \mu_1 \quad \text{in the weak } * \text{ topology of } C(\overline{B}).$$

Here ε_k is some subsequence of ε .

In view of (3.1), for arbitrary compact subset K of $\overline{B} \setminus \{0\}$, there holds

$$\lim_{\varepsilon \rightarrow 0} \frac{|\nabla u_\varepsilon|^n}{|\log \varepsilon|} = 0 \quad \text{in } L^1(K).$$

Therefore, we can find $m_1 \geq 0$ such that

$$\mu_1 = m_1 \delta_o. \quad (3.7)$$

By virtue of (3.5),

$$m_1 \leq (n-1)^{n/2}|\partial B|. \quad (3.8)$$

On the other hand, (3.6) implies

$$m_1 \geq (n-1)^{n/2}|\partial B|.$$

Combining this with (3.8) yields

$$m_1 = (n-1)^{n/2}|\partial B|.$$

Substituting this result into (3.7), and noting the limit $\mu_1 = m_1 \delta_o$ is unique, we get (1.11). $\square$

Proof of (1.12).

Step 1.

By virtue of (3.3), there exists a Radon measure μ_2 , such that

$$\lim_{\varepsilon_k \rightarrow 0} \frac{u_{\varepsilon_k n+1}^2}{2\varepsilon_k^n} = \mu_2, \quad \text{in the weak } * \text{ topology of } C(\overline{B}). \quad (3.9)$$

Similar to the derivation of [8: (4.2), (4.6), (4.8) and Theorem 1.3 in], we can also obtain (1.2) and (1.3), and as $\varepsilon_k \rightarrow 0$,

$$\sup_K \frac{u_{\varepsilon_k n+1}^2}{2\varepsilon_k^n} \leq C\varepsilon_k^n \rightarrow 0; \quad (3.10)$$

$$\lim_{\varepsilon_k \rightarrow 0} u_{\varepsilon_k} = \left(\frac{x}{|x|}, 0 \right) \quad \text{in } C^{1,\beta}(K) \text{ for some } \beta \in (0, 1/2), \quad (3.11)$$

where K is an arbitrary compact subset of $B \setminus \{0\}$. Eqs. (3.9) and (3.10) lead to

$$\mu_2 = m_2 \delta_o, \quad 0 \leq m_2 \in \mathbf{R}.$$

Step 2.

Let $R = \sigma$ and $p = n$ in (2.5), then

$$\begin{aligned} \frac{1}{n} \int_{\partial B_\sigma(0)} v^{n/2} |x| ds - \int_{\partial B_\sigma(0)} v^{(n-2)/2} |\partial_\nu u|^2 |x| ds - \tau \int_{B_\sigma(0)} v^{(n-2)/2} dx \\ = -\frac{1}{2\varepsilon^n} \int_{\partial B_\sigma(0)} |x| u_{n+1}^2 ds + \frac{n}{2\varepsilon^n} \int_{B_\sigma(0)} u_{n+1}^2 dx. \end{aligned}$$

Here $v = |\nabla u|^2 + \tau$. Letting $\tau \rightarrow 0$ and using (1.2) and (1.3), we have

$$\begin{aligned} \frac{1}{n} \int_{\partial B_\sigma(0)} |\nabla u_\varepsilon|^n |x| ds - \int_{\partial B_\sigma(0)} |\nabla u_\varepsilon|^{n-2} |\partial_\nu u_\varepsilon|^2 |x| ds \\ = -\frac{1}{2\varepsilon^n} \int_{\partial B_\sigma(0)} |x| u_{n+1}^2 ds + \frac{n}{2\varepsilon^n} \int_{B_\sigma(0)} u_{n+1}^2 dx. \end{aligned} \quad (3.12)$$

Clearly, (3.10) shows that

$$\frac{1}{2\varepsilon_k^n} \int_{\partial B_\sigma(0)} |x| u_{n+1}^2 ds \rightarrow 0 \quad (3.13)$$

when $\varepsilon_k \rightarrow 0$. Next, from (3.11) we can deduce that as $\varepsilon_k \rightarrow 0$,

$$\begin{aligned} f &\rightarrow \frac{\pi}{2}, & \text{for } r = \sigma; \\ |\partial_\nu u|^2 &= |f_r|^2 \rightarrow 0, & \text{for } r = \sigma; \\ \int_{\partial B_\sigma(0)} |\nabla u_{\varepsilon_k}|^{n-2} |x| ds &\leq C. \end{aligned}$$

Thus, as $\varepsilon_k \rightarrow 0$,

$$\int_{\partial B_\sigma(0)} |\nabla u_{\varepsilon_k}|^{n-2} |\partial_\nu u_{\varepsilon_k}|^2 |x| ds \rightarrow 0. \quad (3.14)$$

At last, there holds

$$|\nabla u|^n \geq \sin^n f \left| \nabla \frac{x}{|x|} \right|^n = \left| \nabla \frac{x}{|x|} \right|^n - (1 - \sin^n f) \left| \nabla \frac{x}{|x|} \right|^n.$$

On the other hand, (3.10) implies that as $\varepsilon_k \rightarrow 0$,

$$\int_{\partial B_\sigma(0)} (1 - \sin^n f) \left| \nabla \frac{x}{|x|} \right|^n ds \rightarrow 0.$$

Therefore,

$$\begin{aligned} \lim_{\varepsilon_k \rightarrow 0} \frac{1}{n} \int_{\partial B_\sigma(0)} |\nabla u_{\varepsilon_k}|^n |x| ds &\geq \frac{1}{n} \int_{\partial B_\sigma(0)} \left| \nabla \frac{x}{|x|} \right|^n |x| ds \\ &= (n-1)^{n/2} \frac{|\partial B|}{n}. \end{aligned} \quad (3.15)$$

Letting $\varepsilon_k \rightarrow 0$ in (3.12), and inserting (3.13)–(3.15) into this result, we obtain

$$\lim_{\varepsilon_k \rightarrow 0} \frac{n}{2\varepsilon_k^n} \int_{B_\sigma} u_{\varepsilon_k n+1}^2 dx \geq (n-1)^{n/2} \frac{|\partial B|}{n}.$$

This result, together with (3.9), implies

$$m_2 \geq (n-1)^{n/2} \frac{|\partial B|}{n^2}. \quad (3.16)$$

Step 3.

By Struwe's idea (cf. [15]), we claim that, if $\varepsilon = \varepsilon_k$ is a sequence converging to zero, then there exists an integer $k_0 > 0$ which is independent of ε , such that as $k > k_0$,

$$\frac{n}{2\varepsilon_k^n} \int_B u_{\varepsilon_k n+1}^2 dx \leq (n-1)^{n/2} \frac{|\partial B|}{n}. \quad (3.17)$$

In fact, if we denote $\nu(\varepsilon) = \inf\{E_\varepsilon(u, G) : u \in W\}$, then the map $\varepsilon \rightarrow E_\varepsilon(u, G)$ is not increasing, and

$$-\frac{\partial}{\partial \varepsilon} E_\varepsilon(u, B) = \frac{n}{2\varepsilon^{n+1}} \int_B u_{n+1}^2 dx.$$

Thus, for the minimizer $u = u_\varepsilon$ of $E_\varepsilon(u, G)$,

$$\begin{aligned} \frac{n}{2\varepsilon^{n+1}} \int_B u_{n+1}^2 dx &= \lim_{\delta \rightarrow 0} \frac{E_\varepsilon(u, G) - E_{\varepsilon+\delta}(u, G)}{\delta} \\ &\leq \lim_{\delta \rightarrow 0} \frac{\nu(\varepsilon) - \nu(\varepsilon + \delta)}{\delta} = -\nu'(\varepsilon), \end{aligned}$$

since $\nu(\varepsilon + \delta) \leq E_{\varepsilon+\delta}(u, G) \leq E_\varepsilon(u, G) = \nu(\varepsilon)$. Suppose (3.17) is not true, there must be a fixed $\varepsilon_0 > 0$, such that for $0 < \varepsilon < \varepsilon_0$,

$$-\nu'(\varepsilon) > (n-1)^{n/2} |\partial B| (n\varepsilon)^{-1}.$$

Integrating from ε to ε_0 , we have

$$\nu(\varepsilon) = - \int_\varepsilon^{\varepsilon_0} \nu'(\varepsilon) d\varepsilon + \nu(\varepsilon_0) > \frac{1}{n} (n-1)^{n/2} |\partial B| \log \frac{\varepsilon_0}{\varepsilon} + \nu(\varepsilon_0).$$

It contradicts (3.5) as long as ε is sufficiently small. (3.17) is proved.

Combining (3.17) with (3.9), we get

$$m_2 \leq (n-1)^{n/2} \frac{|\partial B|}{n^2}.$$

This result, together with (3.16), means

$$m_2 = (n-1)^{n/2} \frac{|\partial B|}{n^2}.$$

Substituting this into (3.9), and noting the limit $\mu_2 = m_2 \delta_o$ is unique, we complete the proof of (1.12). $\square$

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REFERENCES

- [1] ALMOG, Y.—BERLYAND, L.—GOLOVATY, D.—SHAFRIR, I.: *Global minimizers for a p-Ginzburg-Landau-type energy in $\mathbf{R}^2$* , J. Funct. Anal. **256** (2009), 2268–2290.
- [2] BETHUEL, F.—BREZIS, H.—HELEIN, F.: *Ginzburg-Landau Vortices*, Birkhauser, Boston, 1994.

- [3] HAN, Z.—LI, Y.: *Degenerate elliptic systems and applications to Ginzburg-Landau type equations, Part I*, Calc. Var. Partial Differential Equations **4** (1996), 171–202.
- [4] HANG, F.—LIN, F.: *Static theory for Planar Ferromagnets and Antiferromagnets*, Acta. Math. Sinica **17** (2001), 541–580.
- [5] HONG, M.: *Asymptotic behavior for minimizers of a Ginzburg-Landau type functional in higher dimensions associated with n -harmonic maps*, Adv. Differential Equations **1** (1996), 611–634.
- [6] KOMINEAS, S.—PAPANICOLAOU, N.: *Vortex dynamics in two-dimensional antiferromagnets*, Nonlinearity **11** (1998), 265–290.
- [7] LEI, Y.: *Asymptotic analysis of the radial minimizers of an energy functional*, Methods Appl. Anal. **10** (2003), 67–80.
- [8] LEI, Y.: *Asymptotic properties of minimizers of a p -energy functional*, Nonlinear Anal. **68** (2008), 1421–1431.
- [9] LEI, Y.: *Some results on an n -Ginzburg-Landau type minimizer*, J. Comput. Appl. Math. **217** (2008), 123–136.
- [10] LEI, Y.: *Singularity analysis of a p -Ginzburg-Landau type minimizer*, Bull. Sci. Math. **134** (2010), 97–115.
- [11] MOSER, R.: *Ginzburg-Landau vortices for thin ferromagnetic films*, Appl. Math. Res. Express. AMRX (2003), 1–32.
- [12] PAPANICOLAOU, N.—SPATHIS, P. N.: *Semitopological solutions in planar ferromagnets*, Nonlinearity **12** (1999), 285–302.
- [13] QI, L.—LEI, Y.: *On energy estimates for a Landau-Lifschitz type functional in higher dimensions*, J. Korean Math. Soc. **46** (2009), 1207–1218.
- [14] SANDIER E.: *Ginzburg-Landau minimizers from $\mathbf{R}^{n+1}$ to $\mathbf{R}^n$ and minimal connections*, Indiana Univ. Math. J. **50** (2001), 1807–1844.
- [15] STRUWE, M.: *Une estimation asymptotique pour le modèle Ginzburg-Landau*, C. R. Math. Acad. Sci. Paris **317** (1993), 677–680.
- [16] WANG, C.: *Limits of solutions to the generalized Ginzburg-Landau functional*, Comm. Partial Differential Equations **27** (2002), 877–906.

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