Analysis of the Quadrupole Coupling Constants and Mössbauer Isomeric Shifts in Halogen Compounds Within the Gaussian98 Code

O. Kh. Poleshchuk, J. N. Latosińska, and B. Nogaj Institute of Physics, Adam Mickiewicz University, Umultowska 85, 61-614 Poznań, Poland Reprint requests to Dr. J. N. L.; Fax: +48-61-8257758; E-mail: jolanala@amu.edu.pl

Z. Naturforsch. 55 a, 271–275 (2000); received September 11, 1999

Presented at the XVth International Symposium on Nuclear Quadrupole Interactions, Leipzig, Germany, July 25 - 30, 1999.

A comparison of the experimental and calculated nuclear quadrupole coupling constants of diatomic halogen, interhalogen, trihalide ions and complexes of pyridine is a good test of reliability of semiempirical and non-empirical theories as well as their usefulness in the interpretation of quadrupole coupling data from the qualitative and quantitative point of view. Therefore, a practical way of gaining insight into the bonding properties and electronic structure of such systems consists of combining semiempirical MO calculations with experimental data to derive reliable information. In the present work we report the results of ab initio studies of several iodine containing molecules and ions using the calculated NQCC as a test of the quality of the wavefunctions. To demonstrate the quality of our calculations, we compare the calculated halogen-halogen bond length and QCC with the corresponding experimental values for the compounds studied. With a few exceptions, the overall agreement with experiment is most satisfactory. The results on dipole moments and energy levels are also in a good agreement with the experimental data, however these quantities were measured only in the few cases.

Key words: DFT; QCC; Isomeric Mössbauer Shifts; Halogen Compounds.

Introduction

The nuclear quadrupole coupling constants (QCC) of diatomic halogenides, their ions and interhalogen compounds are well known [1]. A good test of various semiempirical as well as nonempirical methods could be the quality of reproduction of NQR parameters. It should be pointed out that all halogen atoms have quadrupole moments either in the ground or in the excited states (in the case of fluorine atom) and therefore could be very useful for testing different methods of calculations.

Usually such an analysis is performed on the basis of the Townes-Dailey approximation [2], which allows a comparison of quadrupole coupling constant (QCC) obtained from experiment and calculated on the basis of density matrix values.

The QCC of halogen atoms have recently been calculated by ab initio methods within the density functional theory (DFT), in particular the $X\alpha$ method

[3 - 6]. However, the latter method is not really a non-empirical one, because it involves an empirical parameter α which varies in the range from 0,6 up to 1. The second group of methods includes: LAPW [7, 8], LMTO [9, 10], KKR [11] and PAW [12], which are the so-called solid state methods. They allow a more precise description of wave functions close to the nuclei and have recently been applied to the calculations of QCC of halogen atoms. The limitation of these methods is that they demand the not always existing knowledge of the molecular packing in a crystal.

In our earlier papers we calculated the QCC by ab initio methods using HONDO and GAUSSIAN94 packages [13, 14] for heavy nuclei, such as iodine, tin and antimony in some complexes containing these elements.

Direct calculations of the electric field gradient at iodine, tin and antimony atoms have shown that the use of a pseudopotential for these atoms does not lead to QCC values reliable enough. The results obtained for chlorine and nitrogen atoms in the extended basis

0932-0784 / 00 / 0100-0271 \$ 06.00 © Verlag der Zeitschrift für Naturforschung, Tübingen · www.znaturforsch.com

set i. e. 6-311G*, are in a good agreement with experimental values.

The aim of this paper is to check the usefulness of the DFT methods for QCC and isomeric shift calculations on halogen atoms.

Calculation Details

The QCC values of halogen atoms were calculated within the density functional theory and using new hybrid functionals, such as B3P86, B3PW91, MPW1PW91, BH, HLYP, and B3LYP. The main aim was to obtain reliable values of the electronic density distribution on atoms as well as the electric field gradients at the nuclei. The full optimisation of geometry was achieved within the GAUSSIAN98 program [15] run on a PC or Cray (PCSS, Domain).

All the functionals used contained local and non-local exchange and correlation terms, and differed in the form of the correlation functional. One of them, BLYP, contains a correlation functional introduced by Lee [16] (which includes both local and non-local terms), while the BP86 functional contains the local and non-local terms proposed by Perdew [17]. The other two approaches are the so-called hybrid methods since they incorporate some degree of exact Hartree-Fock exchange. The Becke hybrid functional (B3PW91) [18] can be expressed in the form

$$E_{XC} = (1 - A) * E_{X}^{Slater} + A * E_{X}^{HF} + B * E_{X}^{Becke} + E_{C}^{LSD} + C * E_{C}^{PW91},$$

where $E_{\rm X}^{\rm HF}$ is the Hartree-Fock exchange, $E_{\rm X}^{\rm Slater}$ the local exchange functional taken after Slater, $E_{\rm X}^{\rm Becke}$ the gradient correction exchange taken after Becke [18], $E_{\rm C}^{\rm LSD}$ the local spin density (LCD) local correlation functional [19], and $E_{\rm C}^{\rm PW91}$ the 1991 gradient functional of Perdew and Wang [19]. The coefficients A,B, and C were determined using accurate experimental data [20] (56 atomisation energies, 42 ionisation potentials and 8 proton affinities). The hybrid B3LYP functional, proposed by Stephens et al. [21], replaces $\{E_{\rm C}^{\rm LSD}+C*E_{\rm C}^{\rm PW91}\}$ by $C*E_{\rm C}^{\rm LYP}+(1-C)E_{\rm C}^{\rm VWN},$ where $E_{\rm C}^{\rm LYP}$ is the correlation functional of Lee et al. [16], and $E_{\rm C}^{\rm VWN}$ is the total correlation functional of Vosko et al. [22]. Note that the $(1-C)E_{\rm C}^{\rm VWN}$ term is added because $E_{\rm C}^{\rm LYP}$ contains both local and non-local terms.

Table 1. The experimental and calculated QCC (MHz) for halogen atoms with the application of the different functional in 3-21G* basis set.

Mol.	Nuclei	HF	MP2	B3 P86	B3P W91	MPW1 PW91	BH and HLYP	B3 LYP	Exp.
$\overline{\text{Cl}_2}$	35Cl	98.0	94.6	97.8	97.9	97.8	98.5	98.0	108.5
Br_2^2	⁷⁹ Br	746	704	749	749	747	752	754	810
I	127 _I	2273	2122	2294	2294	2290	2298	2304	2453
I ₂ BrCl	⁷⁹ Br	819	756	820	819	818	824	827	877
	35 CI	88.9	87.5	88.8	88.9	88.8	89.4	89.3	103.6
ClF	35 CI	122.6	112.5	117.4	117.4	117.6	119.7	117.3	145.9
IBr	^{127}I	2568	2352	2550	2549	2547	2568	2554	2731
	⁷⁹ Br	651	628	667	667	665	665	672	697
IF	^{127}I	3180	2771	3070	3067	3071	3112	3048	3438
IC1	^{127}I	2786	2494	2754	2750	2750	2777	2761	2930
	35Cl	76	77	78	78	78	78	78	82.5
HI	127 _I	1739	1621	1720	1720	1710	1714	1822	1823

Table 2. A comparison of the QCC values (MHz) for halogen atoms calculated with the use of two different basis sets.

Compound	Nuclei	B3LYP/6-31G*	B3LYP/3-21G*	Exp.
Cl ₂	³⁵ Cl	104.6	98.0	108.5
Br ₂	⁷⁹ Br	748	754	810
BrČl	⁷⁹ Br	807	827	877
	35Cl	96.2	89.3	103.6
ClF	35C1	132.8	117.3	145.9
Cl ₃ -	35Cl	108.1	106.3	118
3		51.3	51.2	58
BrCl ₂ -	⁷⁹ Br	855	881	_
2	35Cl	45.8	44.8	48
Br ₃ -	⁷⁹ Br	788	803	814
3		363	372	371
Br ₂ -	⁷⁹ Br	348	352	382
Cl_2^2	35Cl	47.9	46.4	54.6

The calculations were carried out with the use of 3-21G* (for Cl, Br, F and I) atoms and 6-31G* basis sets (for Cl and Br), and the QCC values were obtained from the principal components of the electric field gradient tensor along the principal axes by a program written by the authors.

The materials studied were simple compounds, i.e.: halogens, interhalogen ions, and also some iodides of elements from the groups III, IV, and V of the periodic table. The relativistic effects were not taken into account.

Results

The QCC values calculated for halogen atoms with the use of various functionals, are listed in Table 1. The use of the hybrid B3LYP functional leads to the best agreement with the experimental data. The differences between the experimental and calculated QCC

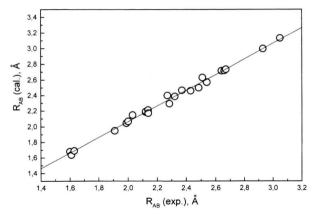


Fig. 1. The correlation between the calculated and experimental bond lengths for the compounds containing halogen atoms.

values for interhalogen molecules and ions do not depend on the halogen type.

The use of the extended basis set 6-31G* does not improve the accuracy of QCC values (Table 2). This conclusion is very important from the point of view of further applications.

A comparison of the geometrical parameters calculated by B3LYP/3-21G* for molecules and ions containing a halogen atom shows that the bond lengths have been overestimated with respect to experimental values. The correlation between the calculated and experimental bond lengths is as follows (Figure 1):

$$R_{X-Y}^{\text{cal.}} = 0,017 + 1,020 R_{X-Y}^{\text{exp.}}$$

 $(r = 0,998; s = 0,03; n = 31),$ (1)

where $R_{\rm X-X}$ is the halogen-halogen or carbon-halogen bond length.

The factor 1,020, close to unity, confirms high quality of the results. Contrary to the $X\alpha$ method [4], the electric dipole moments (μ, D) calculated for halogen compounds are close to the experimental ones [23]:

$$\mu^{\text{cal.}} = -0, 10 + 1, 22\mu^{\text{exp.}}$$

$$(r = 0, 998; s = 0, 2; n = 12).$$
(2)

The use of the d-orbitals for heavy atoms of halogens leads to significantly better results. It is interesting to note that in the observed correlation (2) we included not only interhalogens, but also such molecules, as KI, CHI₃, CH₂I₂, CH₃I, and ICl^oPy complex, i.e. compounds of different symmetries.

Table 3. The calculated (using B3LYP/3-21 G^*) and experimental QCC and U_p values for halogen atoms.

Comp.	Nuclei	$e^2 Q q_{zz}^{\text{exp.}}$	$e^2 Q q_{zz}^{\mathrm{cal.}}$	$U_{\mathrm{p}}^{\mathrm{exp.}}$	$U_{p}^{cal.}$
•		[Mz]	[Mz]	P	P
					relicional de
Cl_2	35Cl	108,5	98,9	0,989	0,901
Br_2	⁷⁹ Br	810	750	1,052	0,974
I_2	¹²⁷ I	2453	2304	1,070	1,005
F_2	¹⁹ F*	127,2	98,0	0,978	0,753
BrCl	35Cl	103,6	89,3	0,944	0,814
	⁷⁹ Br	877	827	1,139	1,074
ClF	35Cl	145,9	117,3	1,329	1,069
IBr 127I	2731	2554	1,191	1,114	
	⁷⁹ Br	697	672	0,905	0,873
IF	^{127}I	3438	3048	1,499	1,329
ICl	^{127}I	2930	2761	1,278	1,204
	35Cl	82,5	78,3	0,752	0,714
HI	^{127}I	1823	1822	0,795	0,794
HF	19F*	40	58	0,308	0,446
Cl ₃ -	35Cl	118; 58	106; 51		0,969; 0,466
BrCl ₂ -	35Cl	48	44,8	0,437	0,408
Br_3^{-2}	⁷⁹ Br	814; 371	803; 372		1,043; 0,483
ICl ₂ -	127 I	3050; 40	2797	1,330	1,220
10.2	35Cl	2020, 10	40,6	0,364	0,370
I_3^-	127 I	2478: 1142			1,068; 0,488
*3	35Cl	45,1	47,2	0,411	0,430
Cl ₂ -	35Cl	54,6	46,4	0,498	0,423
Br_2^{-}	⁷⁹ Br	382	352	0,496	0,457
IBr_2^2	127 I	2770	2575	1,208	1,123
1012	⁷⁹ Br	322	343	0,418	0,445
CF_4	19F*	59,7	47,6	0,459	0,366
CI ₄	127 I	2102	2256	0,917	0,984
CHI ₃	127 I	2029	2107	0,885	0,919
	127 I	1920	1980	0,837	0,863
CH_2I_2	127 I	1739	1864	0,758	0,803
CH ₃ I KI	127 I	0	250	0,738	0,109
	127 I	1242	1209	0,542	0,109
BI ₃	19F*				
BF_3	127 _I	30,2	25,3	0,232	0,195
Geľ ₄	127 I 127 I	1500	1491	0,654	0,650
SiI ₄	127 I 127 I	1335	1365	0,582	0,595
SnI ₄	19F*	1364	1381	0,595	0,602
SF ₆		60	46,5	0,461	0,358
SeF ₆	¹⁹ F*	55,7	44,2	0,428	0,340

Basically, the calculations by the $X\alpha$ method should provide a correlation between the calculated and experimental QCC. We have analysed results of the calculations performed with the $X\alpha$ method [4 - 6] for halogens, interhalogens and their ions. The correlation is as follows:

$$U_{\rm p}^{\rm cal.} = 0,06 + 0,91 U_{\rm p}^{\rm exp.}$$

 $(r = 0,971; s = 0,08; n = 26),$ (3)

where $U_p = e^2 Q q_{zz}/e^2 Q q_{at.}$ is the number of unbalanced p-electrons of halogen atoms. The nuclear QCC used here were taken for ¹²⁷I: 2293 MHz;

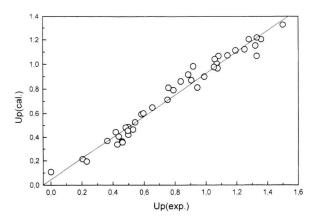


Fig. 2. The correlation between the calculated and experimental $U_{\rm p}$ values for the compounds containing halogen atoms.

⁷⁹Br: 770 MHz; ³⁵Cl: 109,74 MHz, ¹⁹F* = 130 MHz [25, 26].

The relations (1) and (2) indicating a good agreement between calculations and experiment, have encouraged us to perform the calculations of QCC of halogen atoms and isomeric shifts for iodine compounds using the B3LYP/3-21G*. The experimental [1, 3 - 5, 24] and calculated parameters obtained for quite a number of halogen compounds are given in Table 3. On the basis of these results the following correlation (more reliable then (3)) between the experimental and inferred QCC values for all halogen atoms (Figure 2), was found:

$$U_{\rm p}^{\rm cal} = 0,05 + 0,88U_{\rm p}^{\rm exp}$$

 $(r = 0,986; s = 0,05; n = 41).$ (4)

The most important is that correlation (4) was obtained for all nuclei, i. e. on the basis of NQR results (on ³⁵Cl, ⁷⁹Br, ¹²⁷I nuclei) and the perturbed angular distribution method (TDPAD) [27 - 29] (on ¹⁹F* nuclei). It is necessary also to note that such a correlation is valid for all compounds studied, in spite of different environment of the halogen atoms concerned.

Calculations of the isomeric shifts (δ) of iodine compounds seem very important since they can be used to infer about QCC values. The experimental δ values are known for some iodine containing compounds. It is well known that for Mössbauer atoms, such as Sn, Sb, and I, the magnitude of the isomeric shift depends simultaneously on 5s- ($N_{\rm s}$) and 5p-($N_{\rm p}$) orbital populations of these atoms. It was shown (for example [30]), that for SnCl₄ the isomer shift ¹¹⁹Sn depends to a greater degree on the 5s- than

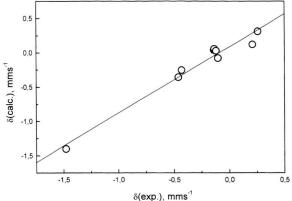


Fig. 3. The correlation between the calculated and experimental isomeric shifts for the compounds containing halogen atoms.

5p-state population. In the case of iodine compounds the following dependence between the isomeric shifts 129 I (rather I_2) and the s- and p-orbital populations is observed (Figure 3):

$$\delta = -10, 9N_s - 1, 0N_p + 23, 3$$

$$(r = 0, 969; s = 0, 10; n = 9)$$
(5)

The experimental [3, 31] and calculated isomeric shifts for the following molecules CI₄, CHI₃, CH₂I₂, CH₃I, SiI₄, SnI₄, KI, ICl and IBr have been compared. The obtained correlation does not seem to be very reliable (Fig. 3), probably due to a low precision of the Mössbauer experiments, but it is also possible to confirm a conclusion about the greater contribution of the 5s-orbital than 5p-orbital of iodine atom to the isomeric shift. The correlation (5) is close enough to the one observed earlier for interhalogen compounds [31].

Conclusions

The DFT in the formulation as in the Gaussian 98 program, in combination with the use of the hybrid functional B3LYP (i.e. taking into account all electrons of atoms) provides very reliable QCC and Mössbauer isomeric shift values. Moreover, the use of the extended basis set 6-311G* does not improve the QCC values significantly.

Acknowledgements

The authors wish to express gratitude to the Russian Federation Education Ministry for support of the research.

- E. A. C. Lucken, Nuclear quadrupole coupling constants, Acad. Press. London 1969.
- [2] C. H. Townes and B. P. Dailey, J. Chem. Phys. 17, 782 (1949).
- [3] L. A. Eriksson, O. L. Malkina, V. G. Malkin, and D. R. Salahub, Int. J. Quantum Chem. 63, 575 (1997).
- [4] G. A.Bowmaker and P. D. W. Boyd, J. Mol. Struct. 150, 327 (1987).
- [5] G. A.Bowmaker and P. D. W. Boyd, J. Chem. Soc. Faraday Trans. 2. 83, 2211 (1987).
- [6] G. A.Bowmaker, P. D. W. Boyd, and R. J. Sorrenson, J. Chem. Soc., Faraday Trans. 2. 80, 1125 (1984).
- [7] P. Blaha, K. Schwarz, and P. Herzig, Phys. Rew. Lett. 54, 1192 (1985).
- [8] K. Schwarz and P. Blaha, Z. Naturforsch. 47a, 197 (1992).
- [9] H. Petrilli, M. Marszalek, and H. Saitovich, Z. Naturforsch. 51a, 537 (1996).
- [10] A. Svane, N. E. Christensen, C. O. Rodriguez, and M. Methfessel, Phys. Rev. B. 55, 12572 (1997).
- [11] H. Akai, M. Akai, S. Blugel, B. Drittler, H. Ebert, K. Terakura, R. Zeller, and P. H. Dederichs, Prog. Theor. Phys. Suppl. 101, 11 (1990).
- [12] H. M. Petrilli, P. E. Blochl, P. Blaha, and K.S chwarz, Phys. Rev. B. 57, 14690 (1998).
- [13] O. Kh. Poleshchuk, G. N. Dolenko, J. Koput, and J. Latosińska, Russian Coord. Chem. 23, 643 (1997).
- [14] O. Kh. Poleshchuk, J. Koput, and J. Latosińska, Russian Coord. Chem., 24, 72 (1998).
 [15] GAUSSIAN 98W. User's Reference. Editors
- [15] GAUSSIAN 98W. User's Reference. Editors E. Fritsch, M. J. Fritsch, Gaussian Inc., Pittsburgh, 1998.
- [16] C. Lee, W. Yang, and R. G. Parr, Phys. Rev. B. 37, 785 (1988).

- [17] J. P. Perdew, Phys. Rev. B. 33, 8822 (1986).
- [18] A. D. Becke, J. Chem. Phys. 98, 5648 (1993).
- [19] J. P. Perdew and Y. Wang, Phys. Rew. B. 45, 13244 (1992).
- [20] L. A. Curtiss, K.Raghavachari, G. W.Trucks, and J. A. Pople, J. Chem. Phys. 94, 7221 (1991).
- [21] P. J. Stephens, F. J. Devlin, C. F. Chablowski, and M. J. Frisch, J. Phys. Chem. 98, 11623 (1994).
- [22] S. H. Vosko, L. Wilk, and M. Nusair, Can. J. Phys. 58, 1200 (1980).
- [23] O. A. Osipov, V. I. Minkin, and A. D. Garnovskii, Spravochnik po dipolnim momentam, Viishava Shkola, 1971.
- [24] G. K. Semin, T. A. Babushkina, and G. G. Yakobson, Nuclear Quadrupole Resonance in Chemistry, John Wiley & Sons, New York 1975.
- [25] G. A. Jeffrey and T. Sakurai, Progress in Solid State Chemistry 1, 380 (1964).
- [26] J. Olliges, A. Loetz, and J. Voitlander, J. Magn. Res. 69, 302 (1986).
- [27] E. Bertholdt, M. Frank, F. Gubitz, W. Kreische, B. Loesch, C. H. Ott, B. Roeseler, M. Schneider, F. Schwab, K. Stammler, and G. Weeske, J. Mol. Struct. 192, 199 (1989).
- [28] J. A. S. Smith, Z. Naturforsch. 41a, 453 (1986).
- [29] M. Frank, F. Gubitz, W. Kreische, A. Labahn, C. Ott, B. Roseler, F. Schwab, and G. Weeske, Hyperfine Interactions. 34, 259 (1987).
- [30] V. A. Varnek and A. A. Varnek, Zh. Struct. Khim. 38, 1160 (1997).
- [31] Molecular Interactions, Ed. H. Ratajczak and W. J. Orville-Thomas, John Wiley & Sons, New York 1981.