

Research Article

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Ground state solutions of nonlinear Schrödinger equations involving the fractional p -Laplacian and potential wells

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Abstract: The purpose of this paper is to investigate the ground state solutions for the following nonlinear Schrödinger equations involving the fractional p -Laplacian

$$(-\Delta)_p^s u(x) + \lambda V(x)u(x)^{p-1} = u(x)^{q-1}, \quad u(x) \geq 0, \quad x \in \mathbb{R}^N,$$

where $\lambda > 0$ is a parameter, $1 < p < q < \frac{Np}{N-sp}$, $N \geq 2$, and $V(x)$ is a real continuous function on \mathbb{R}^N . For λ large enough, the existence of ground state solutions are obtained, and they localize near the potential well $\text{int}(V^{-1}(0))$.

Keywords: nonlinear Schrödinger equation, ground state solution, fractional p -Laplacian, variational methods

MSC 2020: 35J60, 35B33

1 Introduction and main results

In this paper, we consider the following nonlinear Schrödinger equations involving the fractional p -Laplacian

$$\begin{cases} (-\Delta)_p^s u(x) + \lambda V(x)u(x)^{p-1} = u(x)^{q-1}, & x \in \mathbb{R}^N, \\ u(x) \geq 0, & u(x) \in W^{s,p}(\mathbb{R}^N), \end{cases} \quad (1.1)$$

where $\lambda > 0$ is a parameter, $1 < p < q < \frac{Np}{N-sp}$, $N \geq 2$, and $V(x)$ is a real continuous function on \mathbb{R}^N .

We are interested in the existence of ground state solutions for λ big enough, and their asymptotical behavior as $\lambda \rightarrow \infty$. As far as we know, these kinds of problems were first put forward in [1] by Bartsch and Wang, where they studied the Schrödinger equations. Under suitable conditions imposed on the potential, the loss of compactness caused by the whole space \mathbb{R}^N can be recovered when parameter λ is big enough. Then, many authors began studying the problems with potential well. A lot of results have been obtained.

Bartsch and Parnet [2] also considered the nonlinear Schrödinger equation:

$$\begin{cases} -\Delta u + (a_0(x) + \lambda a(x))u = f(x, u), & x \in \mathbb{R}^N, \\ u(x) \rightarrow 0, & |x| \rightarrow \infty, \end{cases}$$

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where $a_0(x) + \lambda a(x)$ is indefinite. By using a local linking theorem and the critical groups theory, they obtained the existence of solutions and their asymptotical behavior as $\lambda \rightarrow \infty$.

Xu and Chen [3] studied the following Kirchhoff problem:

$$\begin{cases} -\left(a + b \int_{\mathbb{R}^N} |\nabla u|^2 dx\right) \Delta u + \lambda V(x)u = f(x, u), & x \in \mathbb{R}^N, \\ u(x) \in H^1(\mathbb{R}^N), \end{cases}$$

where $f(x, u)$ can be sublinear or superlinear. By using the genus theory, they obtained infinitely many negative solutions.

Aleves et al. [4] dealt with the following Choquard equation:

$$\begin{cases} -\Delta u + (\lambda a(x) + 1)u = \left(\frac{1}{|x|^\mu} * |u|^p\right) |u|^{p-2}u, & x \in \mathbb{R}^N, \\ u(x) \in H^1(\mathbb{R}^N), \end{cases}$$

where $\mu \in (0, 3)$, $p \in (2, 6 - \mu)$, and the potential well $\Omega = \bigcup_{j=1}^k \Omega_j$. They proved the existence of a solution, which is nonzero on any subset Ω_j . Furthermore, its asymptotical behavior was investigated.

Zhao et al. [5] studied the Schrödinger-Poisson system allowing the potential $V(x)$ changes sign

$$\begin{cases} -\Delta u + \lambda V(x)u + K(x)\phi u = |u|^{p-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = K(x)u^2 & \text{in } \mathbb{R}^3, \end{cases}$$

where $p \in (3, 6)$ and $V \in C(\mathbb{R}^3, \mathbb{R})$ are bounded from below. Using the variational method, they obtained the existence and asymptotic behavior of nontrivial solutions.

For the critical problems, Clapp and Ding [6] have studied the nonlinear Schrödinger equation:

$$-\Delta u + \lambda V(x)u = \mu u + u^{2^*-1}, \quad x \in \mathbb{R}^N$$

for $N \geq 4$, $\lambda, \mu > 0$. By using variational methods, the authors established existence and multiplicity of positive solutions, which localize near the potential well for λ large and μ small.

Later, the corresponding results obtained in [6] were generalized to the fractional Schrödinger equations by Niu and Tang [7], where they have studied

$$\begin{cases} (-\Delta)^s u + (\lambda V(x) - \mu)u = |u|^{2_s^*-2}u, & x \in \mathbb{R}^N, \\ u \geq 0, & u \in H^s(\mathbb{R}^N). \end{cases}$$

Under the linear perturbation, [6] and [7] obtained the existence of solutions and their asymptotic behavior. For the nonlinear perturbation, Alves and Barros [8] considered

$$-\Delta u + \lambda V(x)u = \mu u^{p-1} + u^{2^*-1}, \quad x \in \mathbb{R}^N.$$

By employing the Ljusternik-Schnirelmann category, for λ big enough and μ small enough, the aforementioned problem has at least $\text{cat}(\Omega)$ positive solutions.

For more results about these kinds of problems and fractional Schrödinger equations, see, for example, [9–24] and references therein. Motivated by the aforementioned results, we consider equation (1.1). The potential function $V(x)$ satisfies

- (V₁) $V(x) \in C(\mathbb{R}^N, \mathbb{R})$ such that $V(x) \geq 0$, $\Omega := \text{int}V^{-1}(0)$ is a nonempty open set of class $C^{0,1}$ with bounded boundary and $V^{-1}(0) = \bar{\Omega}$;
 (V₂) There exists $M_0 > 0$ such that

$$\mu(\{x \in \mathbb{R}^N : V(x) \leq M_0\}) < \infty,$$

where μ denotes the Lebesgue measure on \mathbb{R}^N .

We first introduce some notations. For $s \in (0, 1)$, $p \in [1, +\infty)$, define

$$W^{s,p}(\mathbb{R}^N) := \left\{ u \in L^p(\mathbb{R}^N) : \frac{|u(x) - u(y)|}{|x - y|^{\frac{N}{p} + s}} \in L^p(\mathbb{R}^N \times \mathbb{R}^N) \right\},$$

endowed with the norm

$$\|u\|_{W^{s,p}(\mathbb{R}^N)} := \left(\int_{\mathbb{R}^N} |u|^p dx + \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+sp}} dx dy \right)^{\frac{1}{p}},$$

where the term

$$[u]_{W^{s,p}(\mathbb{R}^N)} := \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy \right)^{\frac{1}{p}}$$

is the so-called Gagliardo (semi)norm of u . Moreover, we define

$$W^{s,p}(\Omega) := \left\{ u \in L^p(\Omega) : \frac{|u(x) - u(y)|}{|x - y|^{\frac{N}{p} + s}} \in L^p(\Omega \times \Omega) \right\},$$

endowed with the norm

$$\|u\|_{W^{s,p}(\Omega)} := \left(\int_{\Omega} |u|^p dx + \int_{\Omega} \int_{\Omega} \frac{|u(x) - u(y)|^p}{|x - y|^{N+sp}} dx dy \right)^{\frac{1}{p}}.$$

Let

$$E_\lambda := \left\{ u \in W^{s,p}(\mathbb{R}^N) : \int_{\mathbb{R}^N} \lambda V(x) |u(x)|^p dx < \infty \right\},$$

with the norm

$$\|u\|_\lambda := \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} \lambda V(x) |u(x)|^p dx \right)^{\frac{1}{p}}.$$

The energy functional associated with (1.1) is

$$J_\lambda(u) = \frac{1}{p} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \frac{\lambda}{p} \int_{\mathbb{R}^N} V(x) |u(x)|^p dx - \frac{1}{q} \int_{\mathbb{R}^N} u^+(x)^q dx \quad \text{for } u \in E_\lambda, \quad (1.2)$$

where $u^+ = \max\{u, 0\}$. Then, we can define the Nehari manifold

$$\mathcal{M}_\lambda := \left\{ u \in E_\lambda \setminus \{0\} : \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \lambda \int_{\mathbb{R}^N} V(x) |u(x)|^p dx = \int_{\mathbb{R}^N} u^+(x)^q dx \right\}$$

and

$$c_\lambda := \inf\{J_\lambda(u) : u \in \mathcal{M}_\lambda\}.$$

Consider the following “limit” problem of (1.1)

$$\begin{cases} (-\Delta)_p^s u(x) = u(x)^{q-1}, & x \in \Omega, \\ u(x) \geq 0, & x \in \Omega, \\ u(x) = 0, & x \in \mathbb{R}^N \setminus \Omega, \end{cases} \quad (1.3)$$

Define a subspace E_0 of $W^{s,p}(\mathbb{R}^N)$ as follows:

$$\begin{aligned} E_0 &:= \{u \in W^{s,p}(\mathbb{R}^N) : u(x) = 0 \text{ in } \mathbb{R}^N \setminus \Omega\} \\ \text{tr}_\Omega E_0 &= \{u|_\Omega : u \in E_0\}. \end{aligned} \quad (1.4)$$

The energy functional associated with (1.3) can be defined by

$$\Phi(u) = \frac{1}{p} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy - \frac{1}{q} \int_{\Omega} u^+(x)^q dx \quad \text{for } u \in E_0.$$

Then, the associated Nehari manifold is

$$\mathcal{N} := \left\{ u \in E_0 \setminus \{0\} : \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy = \int_{\Omega} u^+(x)^q dx \right\}$$

and

$$c(\Omega) := \inf\{\Phi(u) : u \in \mathcal{N}\}.$$

Definition 1.1. A function $u_\lambda(x)$ is a ground state solution of (1.1) if c_λ is achieved by $u_\lambda \in \mathcal{M}_\lambda$, which is a critical point of J_λ . Similarly, a function $u(x)$ is a ground state solution of (1.3) if $c(\Omega)$ is achieved by $u \in \mathcal{N}$, which is a critical point of Φ .

Definition 1.2. Let X be a Banach space, $\varphi \in C^1(X, \mathbb{R})$. The function φ satisfies the $(PS)_c$ condition if any sequence $\{u_n\} \subseteq X$, such that

$$\varphi(u_n) \rightarrow c, \quad \varphi'(u_n) \rightarrow 0 \quad (1.5)$$

has a convergent subsequence. The sequence $\{u_n\}$ that satisfies (1.5) is called to be a $(PS)_c$ sequence of φ .

Our main results read as follows:

Theorem 1.3. Suppose (V_1) and (V_2) hold, then for λ large, (1.1) has a ground state solution $u_\lambda(x)$. Furthermore, any sequence $\lambda_n \rightarrow \infty$, $\{u_{\lambda_n}(x)\}$ has a subsequence such that u_{λ_n} converges in $W^{s,p}(\mathbb{R}^N)$ along the subsequence to a ground state solution u of (1.3).

Theorem 1.4. Suppose (V_1) and (V_2) hold. Let $u_n, n \in \mathbb{N}$ be a sequence of solutions of (1.1) with λ being replaced by λ_n ($\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$) such that $\limsup_{n \rightarrow \infty} J_\lambda(u_n) < \infty$. Then, $u_n(x)$ converges strongly along a subsequence in $W^{s,p}(\mathbb{R}^N)$ to a solution u of (1.3).

The following paper is organized as follows: In Section 2, we will give some preliminary results. Section 3 is devoted to the “limit” problem, and Section 4 contains the proofs of the main results. C denotes various generic positive constants, and $o(1)$ will be used to represent quantities that tend to 0 as λ (or n) $\rightarrow \infty$.

2 Preliminary results

Lemma 2.1. Let $\lambda_0 > 0$ be a fixed constant. Then, for $\lambda \geq \lambda_0 > 0$, $V(x)$ satisfying (V_1) and (V_2) , E_λ is continuously embedded in $W^{s,p}(\mathbb{R}^N)$ uniformly in λ .

Proof. By the definition of $W^{s,p}(\mathbb{R}^N)$ and E_λ , we only need to prove the following inequality:

$$\int_{\mathbb{R}^N} |u(x)|^p dx \leq C \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} \lambda V(x) |u(x)|^p dx \right). \quad (2.1)$$

Define

$$D := \{x \in \mathbb{R}^N : V(x) \leq M_0\}$$

and

$$D^{\delta_0} := \{x \in \mathbb{R}^N : \text{dist}(x, D) \leq \delta_0\}.$$

Take $\zeta \in C^\infty(\mathbb{R}^N, \mathbb{R})$, $0 \leq \zeta \leq 1$, satisfying

$$\zeta(x) = \begin{cases} 1, & x \in D, \\ 0, & x \notin D^{\delta_0}, \end{cases} \quad |\nabla \zeta| \leq C/\delta_0. \quad (2.2)$$

Then, for any function $u \in E_\lambda$, we can obtain

$$\int_{\mathbb{R}^N} (1 - \zeta^p) |u(x)|^p dx = \int_{\mathbb{R}^N \setminus D} (1 - \zeta^p) |u(x)|^p dx + \int_D (1 - \zeta^p) |u(x)|^p dx \leq \frac{1}{\lambda_0 M_0} \lambda \int_{\mathbb{R}^N} V(x) |u(x)|^p dx \quad (2.3)$$

and

$$\begin{aligned} \int_{\mathbb{R}^N} \zeta^p |u(x)|^p dx &= \int_{D^{\delta_0}} \zeta^p |u(x)|^p dx \\ &\leq \mu(D^{\delta_0})^{1 - \frac{p}{p_s^*}} \left(\int_{D^{\delta_0}} |u(x)|^{p_s^*} dx \right)^{\frac{p}{p_s^*}} \\ &\leq C \mu(D^{\delta_0})^{1 - \frac{p}{p_s^*}} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy, \end{aligned} \quad (2.4)$$

where we have used (V_2) and the Sobolev trace inequality

$$\left(\int_{\mathbb{R}^N} |u(x)|^{p_s^*} dx \right)^{1/p_s^*} \leq C \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy \right)^{1/p},$$

for $u \in W^{s,p}(\mathbb{R}^N)$ and $C = C(N, p, s) > 0$. Thus, (2.1) follows from (2.3) and (2.4). \square

Lemma 2.2. *There exists $\sigma > 0$ independent of λ , such that $\|u\|_\lambda \geq \sigma$ for all $u \in \mathcal{M}_\lambda$.*

Proof. From Lemma 2.1, for any $u \in \mathcal{M}_\lambda$,

$$\begin{aligned} 0 &= \langle J'_\lambda(u), u \rangle = \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \lambda \int_{\mathbb{R}^N} V(x) |u(x)|^p dx - \int_{\mathbb{R}^N} u^+(x)^q dx \\ &\geq \|u\|_\lambda^p - C \|u\|_{W^{s,p}(\mathbb{R}^N)}^q \\ &\geq \|u\|_\lambda^p - C \|u\|_\lambda^q, \end{aligned}$$

where $C > 0$ is independent of $\lambda \geq 0$. The aforementioned inequality implies that $\|u\|_\lambda^{q-p} \geq \frac{1}{C}$. Choosing $\sigma = \left(\frac{1}{C}\right)^{\frac{1}{q-p}}$, we obtain $\|u\|_\lambda \geq \sigma$. \square

Lemma 2.3. *Let λ_0 be a fixed positive constant, there exists $c_0 > 0$ independent of $\lambda \geq \lambda_0 > 0$, such that if $\{u_n\}$ is a $(PS)_c$ sequence of J_λ , then either $c \geq c_0$ or $c = 0$. Moreover,*

$$\limsup_{n \rightarrow \infty} \|u_n\|_\lambda^p \leq \frac{pq}{q-p} c. \quad (2.5)$$

Proof. From the definition of $(PS)_c$ sequence,

$$\begin{aligned} c + \|u_n\|_\lambda \cdot o(1) &= J_\lambda(u_n) - \frac{1}{q} \langle J'_\lambda(u_n), u_n \rangle \\ &= \left(\frac{1}{p} - \frac{1}{q} \right) \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy + \lambda \int_{\mathbb{R}^N} V(x) |u_n(x)|^p dx \right) \\ &= \frac{q-p}{pq} \|u_n\|_\lambda^p. \end{aligned}$$

Then, (2.5) holds. On the other side, there is a constant $C > 0$ independent of $\lambda \geq \lambda_0 > 0$, such that

$$\langle J'_\lambda(u), u \rangle = \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy + \lambda \int_{\mathbb{R}^N} V(x) |u(x)|^p dx - \int_{\mathbb{R}^N} u^+(x)^q dx \geq \|u\|_\lambda^p - C \|u\|_\lambda^q.$$

Thus, there exists $\sigma_1 > 0$ independent of λ , such that

$$\frac{1}{4} \|u\|_\lambda^p \leq \langle J'_\lambda(u), u \rangle \text{ for } \|u\|_\lambda < \sigma_1. \quad (2.6)$$

If $c < \frac{\sigma_1^p(q-p)}{pq}$, then

$$\limsup_{n \rightarrow \infty} \|u_n\|_\lambda^p \leq \frac{cpq}{q-p} < \sigma_1^p.$$

Hence, $\|u_n\|_\lambda < \sigma_1$ for n large. It follows from (2.6) that

$$\frac{1}{4} \|u_n\|_\lambda^p \leq \langle J'_\lambda(u_n), u_n \rangle = o(1) \|u_n\|_\lambda,$$

which implies $\|u_n\|_\lambda \rightarrow 0$ as $n \rightarrow \infty$. Therefore, $J_\lambda(u_n) \rightarrow 0$, that is, $c = 0$. Thus, $c_0 = \frac{\sigma_1^p(q-p)}{qp}$ is as required. \square

Lemma 2.4. *There exists $\delta_0 > 0$, such that any $(PS)_c$ sequence $\{u_n\}$ of J_λ with $\lambda \geq 0$ and $c > 0$ satisfies*

$$\liminf_{n \rightarrow \infty} \|u_n^+\|_{L^q(\mathbb{R}^N)}^q \geq \delta_0 c. \quad (2.7)$$

Proof. From the definition of $(PS)_c$ sequence,

$$c = \lim_{n \rightarrow \infty} \left(J_\lambda(u_n) - \frac{1}{p} \langle J'_\lambda(u_n), u_n \rangle \right) = \left(\frac{1}{p} - \frac{1}{q} \right) \lim_{n \rightarrow \infty} \int_{\mathbb{R}^N} u_n^+(x)^q dx = \frac{(q-p)}{qp} \lim_{n \rightarrow \infty} \|u_n^+(x)\|_{L^q(\mathbb{R}^N)}^q,$$

which implies (2.7) with $\delta_0 \leq \frac{qp}{q-p}$. \square

Lemma 2.5. *Let C_1 be any fixed constant. Then, for any $\varepsilon > 0$, there exists $\Lambda_\varepsilon > 0$ and $R_\varepsilon > 0$, such that if $\{u_n\}$ is a $(PS)_c$ sequence of J_λ with $\lambda \geq \Lambda_\varepsilon$, $c \leq C_1$, then*

$$\limsup_{n \rightarrow \infty} \int_{B_{R_\varepsilon}^c} u_n^+(x)^q dx \leq \varepsilon, \quad (2.8)$$

where $B_{R_\varepsilon}^c = \{x \in \mathbb{R}^N : |x| \geq R_\varepsilon\}$.

Proof. For $R > 0$, let

$$A(R) := \{x \in \mathbb{R}^N : |x| > R, V(x) \geq M_0\}$$

and

$$B(R) := \{x \in \mathbb{R}^N : |x| > R, V(x) < M_0\}.$$

It follows from Lemma 2.3 that

$$\begin{aligned} \int_{A(R)} |u_n(x)|^p dx &\leq \frac{1}{\lambda M_0} \int_{\mathbb{R}^N} \lambda V(x) |u_n(x)|^p dx \\ &\leq \frac{1}{\lambda M_0} \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} \lambda V(x) |u_n(x)|^p dx \right) \\ &\leq \frac{1}{\lambda M_0} \left(\frac{pq}{q-p} C_1 + o(1) \right). \end{aligned} \quad (2.9)$$

From Hölder inequality and (2.5), we can see that, for $1 < r < N/(N - ps)$,

$$\int_{B(R)} |u_n(x)|^p dx \leq \left(\int_{\mathbb{R}^N} |u_n(x)|^{pr} dx \right)^{1/r} \mu(B(R))^{1/r'} \leq C \|u_n\|_\lambda^p \cdot \mu(B(R))^{1/r'} \leq C \frac{pq}{q-p} C_0 \cdot \mu(B(R))^{1/r'}, \quad (2.10)$$

where $C = C(N, r) > 0$ and $1/r + 1/r' = 1$. By interpolation inequality and Sobolev embedding inequality, we can obtain

$$\begin{aligned} \int_{B_R^c} u_n^+(x)^q dx &\leq \left(\int_{B_R^c} |u_n(x)|^p dx \right)^{\frac{q(1-\theta)}{p}} \cdot \left(\int_{B_R^c} |u_n(x)|^{p_s^*} dx \right)^{\frac{q\theta}{p_s^*}} \\ &\leq \left(\int_{B_R^c} |u_n(x)|^p dx \right)^{\frac{q(1-\theta)}{p}} \left(\int_{\mathbb{R}^N} |u_n(x)|^{p_s^*} dx \right)^{\frac{q\theta}{p_s^*}} \\ &\leq C \left(\int_{B_R^c} |u_n(x)|^p dx \right)^{\frac{q(1-\theta)}{p}} \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy \right)^{\frac{q\theta}{p}} \\ &\leq C \left(\int_{A(R)} |u_n(x)|^p dx + \int_{B(R)} |u_n(x)|^p dx \right)^{\frac{q(1-\theta)}{p}} \|u_n\|_\lambda^{q\theta}, \end{aligned}$$

where $\theta = \frac{N}{s} \frac{q-p}{pq}$. Then, the result follows from (2.9), (2.10) and (V₂). \square

Lemma 2.6. (Brézis-Lieb lemma, 1983) Let $\{u_n\} \subset L^p(\mathbb{R}^N)$, $1 \leq p < \infty$. If

- (a) $\{u_n\}$ is bounded in $L^p(\mathbb{R}^N)$,
- (b) $u_n \rightarrow u$ almost everywhere on \mathbb{R}^N , then

$$\lim_{n \rightarrow \infty} (|u_n|_p^p - |u_n - u|_p^p) = |u|_p^p. \quad (2.11)$$

Lemma 2.7. Let $\lambda \geq \lambda_0 > 0$ be fixed and let $\{u_n\}$ be a $(PS)_c$ sequence of J_λ . Then, up to a subsequence, $u_n \rightharpoonup u$ in E_λ with u being a weak solution of (1.1). Moreover, $u_n^1 = u_n - u$ is $(PS)_{c'}$ sequence with $c' = c - J_\lambda(u)$.

Proof. By Lemma 2.3, $\{u_n\}$ is bounded in E_λ . Then, up to a subsequence $u_n \rightharpoonup u$ in E_λ as $n \rightarrow \infty$, and

$$u_n \rightharpoonup u \quad \text{in} \quad W^{s,p}(\mathbb{R}^N), \quad (2.12)$$

$$u_n \rightharpoonup u \quad \text{in} \quad L^q(\mathbb{R}^N), \quad p \leq q < p_s^*, \quad (2.13)$$

$$u_n \rightarrow u \quad \text{in} \quad L_{\text{loc}}^q(\mathbb{R}^N), \quad p \leq q < p_s^*, \quad (2.14)$$

$$u_n \rightarrow u \quad \text{a.e. in } \mathbb{R}^N, \quad (2.15)$$

where $p_s^* = \frac{Np}{N-ps}$ is the fractional critical Sobolev exponent. Hence, for any $\varphi \in E_\lambda$, we have

$$\begin{aligned} \langle J'_\lambda(u_n), \varphi \rangle &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^{p-2} (u_n(x) - u_n(y)) (\varphi(x) - \varphi(y))}{|x - y|^{N+ps}} dx dy \\ &\quad + \lambda \int_{\mathbb{R}^N} V(x) |u_n(x)|^{p-2} u_n(x) \varphi(x) dx - \int_{\mathbb{R}^N} u_n^+(x)^{q-1} \varphi(x) dx \\ &\rightarrow \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y)) (\varphi(x) - \varphi(y))}{|x - y|^{N+ps}} dx dy \\ &\quad + \lambda \int_{\mathbb{R}^N} V(x) u(x)^{p-1} \varphi(x) dx - \int_{\mathbb{R}^N} u^+(x)^{q-1} \varphi(x) dx = \langle J'_\lambda(u), \varphi \rangle. \end{aligned}$$

Therefore,

$$\langle J'_\lambda(u), \varphi \rangle = \lim_{n \rightarrow \infty} \langle J'_\lambda(u_n), \varphi \rangle = 0, \quad (2.16)$$

which implies that u is a critical point of J_λ .

Let $u_n^1 = u_n - u$, we will show that as $n \rightarrow \infty$,

$$J_\lambda(u_n^1) \rightarrow c - J_\lambda(u) \quad (2.17)$$

and

$$J'_\lambda(u_n^1) \rightarrow 0. \quad (2.18)$$

To show (2.17), we observe that

$$\begin{aligned} J_\lambda(u_n^1) &= \frac{1}{p} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n^1(x) - u_n^1(y)|^p}{|x - y|^{N+ps}} dx dy + \frac{\lambda}{p} \int_{\mathbb{R}^N} V(x) |u_n^1(x)|^p dx - \frac{1}{q} \int_{\mathbb{R}^N} u_n^{1+}(x)^q dx \\ &= J_\lambda(u_n) - J_\lambda(u) + \frac{\lambda}{p} \int_{\mathbb{R}^N} V(x) (|u_n^1(x)|^p - |u_n(x)|^p + |u(x)|^p) dx \\ &\quad + \frac{1}{p} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n^1(x) - u_n^1(y)|^p - |u_n(x) - u_n(y)|^p + |u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy \\ &\quad + \frac{1}{q} \int_{\mathbb{R}^N} u_n^+(x)^q dx - \frac{1}{q} \int_{\mathbb{R}^N} u_n^{1+}(x)^q dx - \frac{1}{q} \int_{\mathbb{R}^N} u^+(x)^q dx. \end{aligned} \quad (2.19)$$

From Lemma 2.6, $\int_{\mathbb{R}^N} u_n^+(x)^q dx - \int_{\mathbb{R}^N} u^+(x)^q dx - \int_{\mathbb{R}^N} u_n^{1+}(x)^q dx \rightarrow 0$ as $n \rightarrow \infty$. Conversely, we know that $\|u_n\|_\lambda^p - \|u\|_\lambda^p - \|u_n^1\|_\lambda^p \rightarrow 0$, as $n \rightarrow \infty$. Thus, from (2.19), we indeed have obtained (2.17). Now we come to show (2.18). From (2.16), we have for any $\varphi \in E_\lambda$

$$\langle J'_\lambda(u_n^1), \varphi \rangle = \langle J'_\lambda(u_n), \varphi \rangle - \int_{\mathbb{R}^N} (u_n^+)^{q-1} \varphi(x) dx + \int_{\mathbb{R}^N} (u_n^+)^{q-1} \varphi(x) dx - \int_{\mathbb{R}^N} (u^+)^{q-1} \varphi(x) dx + o(1).$$

Since $J'_\lambda(u_n) \rightarrow 0$ and $u_n \rightarrow u$ in $L^q(\mathbb{R}^N)$, we have

$$\lim_{n \rightarrow \infty} \sup_{\|\varphi\|_\lambda \leq 1} \int_{\mathbb{R}^N} ((u_n^+)^{q-1}(x) \varphi(x) - (u^+)^{q-1} \varphi(x) + (u^+)^{q-1} \varphi(x)) dx = 0.$$

Thus, we have

$$\lim_{n \rightarrow \infty} \langle J'_\lambda(u_n^1), \varphi \rangle = 0 \quad \text{for any } \varphi \in E_\lambda,$$

which implies (2.18), and this completes the proof. \square

Proposition 2.8. Suppose (V_1) and (V_2) hold. Then, for any $C_0 > 0$, there exists $\Lambda_0 > 0$ such that J_λ satisfies the $(PS)_c$ condition for all $\lambda \geq \Lambda_0$ and $c \leq C_0$.

Proof. Choose $0 < \varepsilon < \delta_0 c_0/2$, where c_0 and δ_0 are the constants in Lemmas 2.3 and 2.4, respectively. Let $\Lambda_0 := \Lambda_\varepsilon$, where $\Lambda_\varepsilon > 0$ is from Lemma 2.5.

Assume $\{u_n\}$ is a $(PS)_c$ sequence of J_λ with $\lambda \geq \Lambda_0$ and $c \leq C_0$. By Lemma 2.7, $u_n^1 = u_n - u$ is a $(PS)_{c'}$ sequence of J_λ with $c' = c - J_\lambda(u)$. If $c' > 0$, it follows from Lemma 2.3 that $c' \geq c_0$. From Lemma 2.4, we can obtain

$$\liminf_{n \rightarrow \infty} \|u_n^{1+}(\cdot)\|_{L^q(\mathbb{R}^N)}^q \geq \delta_0 c' \geq \delta_0 c_0.$$

Conversely, Lemma 2.5 implies

$$\limsup_{n \rightarrow \infty} \int_{B_{R_\varepsilon}^c} u_n^{1+}(x)^q \leq \varepsilon < \frac{\delta_0 c_0}{2}.$$

Noting $u_n^1 \rightarrow 0$ in $L_{\text{loc}}^q(\mathbb{R}^N)$, $p \leq q < p_s^*$, a contradiction follows from the aforementioned two inequalities. Therefore, $c' = 0$. Thus, $u_n^1 \rightarrow 0$ in E_λ by Lemma 2.3. \square

Corollary 2.9. For any $q \in (p, p_s^*)$, there exists $\Lambda_0 > 0$, such that c_λ is achieved for all $\lambda \geq \Lambda_0$ at some $u_\lambda \in E_\lambda$, which is a ground state solution of (1.1).

Proof. By Ekeland variational principle, there is a PS sequence $u_n \in E_\lambda$, such that

$$J_\lambda(u_n) \rightarrow c_\lambda \quad \text{and} \quad J'_\lambda(u_n) \rightarrow 0.$$

By Proposition 2.8, there exists some $u_\lambda \in E_\lambda$, such that, up to subsequence, $u_n \rightarrow u_\lambda$ in E_λ as $n \rightarrow \infty$ and λ is sufficiently large. It is not difficult to show that

$$J_\lambda(u_n) \rightarrow J_\lambda(u_\lambda) \quad \text{and} \quad J'_\lambda(u_n) \rightarrow J'_\lambda(u_\lambda).$$

Therefore, we have $J_\lambda(u_\lambda) = c_\lambda$ and $J'_\lambda(u_\lambda) = 0$. This means that u_λ is a ground state solution of (1.1). \square

3 Limit problem

Lemma 3.1. Let $1 < p < q < p_s^* := \frac{pN}{N-ps}$, $N \geq 2$. Then, $tr_\Omega E_0$ is compactly embedded in $L^q(\Omega)$.

Proof. Since $tr_\Omega E_0 \subset W^{s,p}(\Omega)$ and $W^{s,p}(\Omega) \hookrightarrow L^p(\Omega)$ are compact for $p < q < p_s^*$, $N \geq 2$, the result follows. \square

Lemma 3.2. The infimum $c(\Omega)$ is achieved by a function $u \in \mathcal{N}$, which is a ground state solution of (1.3).

Proof. By Ekeland variational principle, there is a PS sequence $u_n \in E_0$, such that

$$\Phi(u_n) \rightarrow c(\Omega) \quad \text{and} \quad \Phi'(u_n) \rightarrow 0.$$

Thus, by Lemma 3.1, we can easily obtain a subsequence of $\{u_n\}$ (still denote it itself), such that $u_n \rightarrow u$ in E_0 . Therefore, u is a ground state solution of (1.3). \square

Remark 3.3. Assume set $\Omega = \text{int} V^{-1}(0)$ has more than one isolated component, for example, $\Omega = \Omega_1 \cup \Omega_2$ with $\Omega_1 \cap \Omega_2 = \emptyset$. Suppose that $u \in \mathcal{N}$ is a nonnegative solution of (1.3) with $u(x) = 0$ in Ω_1 and $u(x) \not\equiv 0$ in Ω_2 . Then, we have $(-\Delta)_p^s u(x) = \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^{p-2} (u(x) - u(y))}{|x - y|^{N+ps}} dy < 0$ in Ω_1 . Conversely, $(-\Delta)_p^s u(x) = u(x)^{q-1} = 0$ for

$x \in \Omega_1$. This contradiction shows that the nonnegative solution $u(x)$ of (1.3) must be $u(x) \not\equiv 0$ in both Ω_1 and Ω_2 . However, the Laplacian case can have a nonnegative solution $u(x)$ satisfying $u(x) = 0$ in Ω_1 and $u(x) \not\equiv 0$ in Ω_2 . The difference between the two phenomena is attributed to the nonlocality of fractional operators and the locality of Laplacian operators.

4 The proof of the main results

Lemma 4.1. $c_\lambda \rightarrow c(\Omega)$ as $\lambda \rightarrow \infty$.

Proof. From the definition of c_λ and $c(\Omega)$, we know that $c_\lambda \leq c(\Omega)$, $\lambda > 0$. Furthermore, c_λ is monotone increasing about the parameter $\lambda > 0$. Then, there exists a constant k , such that

$$\lim_{n \rightarrow \infty} c_{\lambda_n} = k,$$

where $\lambda_n \rightarrow \infty$. It follows from Lemma 2.3 that $k > 0$. By Corollary 2.9, for n large enough, there exists a sequence $u_n \in \mathcal{M}_{\lambda_n}$, such that $J'_{\lambda_n}(u_n) = 0$ and $J_{\lambda_n}(u_n) = c_{\lambda_n}$. If $k < c(\Omega)$, it is easy to see that $\{u_n\}$ is bounded in $W^{s,p}(\mathbb{R}^N)$; thus, we can assume that $u_n \rightharpoonup u$ in E and

$$u_n(x) \rightarrow u(x) \quad \text{in } L^\theta_{\text{loc}}(\mathbb{R}^N) \quad \text{for } p \leq \theta < p_s^*. \quad (4.1)$$

Claim 1: $u|_{\Omega^c} = 0$. In fact, if $u|_{\Omega^c} \neq 0$, then there exists a compact subset $F \subset \Omega^c$ with $\text{dist}(F, \Omega) > 0$, such that $u|_F \neq 0$. It follows from (4.1) that

$$\int_F |u_n(x)|^p dx \rightarrow \int_F |u(x)|^p dx > 0.$$

However, there exists $\varepsilon_0 > 0$, such that $V(x) \geq \varepsilon_0 > 0$, $x \in F$. Thus,

$$J_{\lambda_n}(u_n) \geq \frac{q-p}{pq} \lambda_n \int_F V(x) |u_n(x)|^p dx \geq \frac{q-p}{pq} \lambda_n \varepsilon_0 \int_F |u_n(x)|^p dx \rightarrow \infty \quad \text{as } n \rightarrow \infty,$$

which is a contradiction. Therefore, $u \in E_0$.

Claim 2: $u_n \rightarrow u$ in $L^q(\mathbb{R}^N)$ for $p < q < p_s^*$. Indeed, if not, then by the concentration-compactness lemma from the study by Loins [25], there exist $\delta > 0$, $\rho > 0$ and $x_n \in \mathbb{R}^N$ with $|x_n| \rightarrow \infty$, such that

$$\liminf_{n \rightarrow \infty} \int_{B_\rho(x_n)} |u_n(x) - u(x)|^p dx \geq \delta > 0.$$

Then, we have

$$\begin{aligned} J_{\lambda_n}(u_n) &= \frac{q-p}{pq} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x-y|^{N+ps}} dx dy + \frac{q-p}{pq} \int_{\mathbb{R}^N} \lambda_n V(x) |u_n(x)|^p dx \\ &\geq \frac{q-p}{pq} \lambda_n \int_{B_\rho(x_n) \cap \{x: V(x) \geq M_0\}} V(x) |u_n(x)|^p dx \\ &= \frac{q-p}{pq} \lambda_n \int_{B_\rho(x_n) \cap \{x: V(x) \geq M_0\}} V(x) |u_n(x) - u(x)|^p dx \\ &\geq \frac{q-p}{pq} \lambda_n \left(M_0 \int_{B_\rho(x_n)} |u_n(x) - u(x)|^p dx - M_0 \int_{B_\rho(x_n) \cap \{x: V(x) \leq M_0\}} |u_n(x)|^p dx \right) \\ &\geq \frac{q-p}{pq} \lambda_n \left(M_0 \int_{B_\rho(x_n)} |u_n(x) - u(x)|^p dx - o(1) \right) \rightarrow \infty \quad \text{as } n \rightarrow \infty, \end{aligned}$$

as a contradiction. So $u_n \rightarrow u$ in $L^p(\mathbb{R}^N)$. Therefore, it is easy to see that $u \geq 0$ is a solution for problem (1.3). Furthermore,

$$k = \lim_{n \rightarrow \infty} c_{\lambda_n} = \lim_{n \rightarrow \infty} J_{\lambda_n}(u_n) = \lim_{n \rightarrow \infty} \left(\frac{1}{p} - \frac{1}{q} \right) \int_{\mathbb{R}^N} u_n^+(x)^q dx = \left(\frac{1}{p} - \frac{1}{q} \right) \int_{\Omega} u^+(x)^q dx,$$

which means $u \in \mathcal{N}$, and then, $k \geq c(\Omega)$, a contradiction. Hence, $\lim_{\lambda \rightarrow \infty} c_\lambda = c(\Omega)$. \square

Proof of Theorem 1.3. By Corollary 2.9, there exists $u_n \in \mathcal{M}_{\lambda_n}$, such that $J_{\lambda_n}(u_n) = c_{\lambda_n}$ ($\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$). It is easy to see that $\{u_n\}$ is bounded in $W^{s,p}(\mathbb{R}^N)$. Then, without loss of generality, $u_n \rightharpoonup u$ in $W^{s,p}(\mathbb{R}^N)$ and $u_n \rightarrow u$ in $L^{\theta}_{\text{loc}}(\mathbb{R}^N)$ for $p < \theta < p^*$.

Now we prove that $u_n \rightarrow u$ strongly in $W^{s,p}(\mathbb{R}^N)$ and u is a ground state solution of (1.3). First, as the proof of Lemma 4.1, $u \geq 0$ is a solution of problem (1.3) and $u_n^+ \rightarrow u^+$ strongly in $L^q(\mathbb{R}^N)$.

Now we claim that

$$\lambda_n \int_{\mathbb{R}^N} V(x) |u_n(x)|^p dx \rightarrow 0$$

and

$$\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy \rightarrow \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy.$$

Indeed, if either

$$\liminf_{n \rightarrow \infty} \lambda_n \int_{\mathbb{R}^N} V(x) |u_n(x)|^p dx > 0$$

or

$$\liminf_{n \rightarrow \infty} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy > \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy.$$

Thus, we have

$$\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy < \int_{\Omega} u^+(x)^q dx.$$

Therefore, there is $\alpha \in (0, 1)$, such that $\alpha u \in \mathcal{N}$ and

$$\begin{aligned} c(\Omega) &\leq \Phi(\alpha u) = \frac{q-p}{pq} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|\alpha u(x) - \alpha u(y)|^p}{|x - y|^{N+ps}} dx dy \\ &< \frac{q-p}{pq} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{N+ps}} dx dy \\ &\leq \lim_{n \rightarrow \infty} \frac{q-p}{pq} \left(\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{N+ps}} dx dy + \int_{\mathbb{R}^N} \lambda_n V(x) |u_n(x)|^p dx \right) \\ &= \lim_{n \rightarrow \infty} J_{\lambda_n}(u_n) = c(\Omega), \end{aligned}$$

which is a contradiction. By now we complete the proof of Theorem 1.3. \square

Proof of Theorem 1.4. Suppose $\{u_n\} \subset W^{s,p}(\mathbb{R}^N)$ is a solution of (1.1) with λ being replaced by λ_n ($\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$). It follows from $\limsup_{n \rightarrow \infty} J_{\lambda_n}(u_n) < \infty$ that such a sequence $\{u_n\}$ is bounded in $W^{s,p}(\mathbb{R}^N)$.

Suppose that $u_n \rightharpoonup u$ in $W^{s,p}(\mathbb{R}^N)$ and $u_n \rightarrow u$ in $L_{\text{loc}}^\theta(\mathbb{R}^N)$ for $p < \theta < p_s^*$. Similar to the proof of Lemma 4.1, $u|_{\Omega^c} = 0$ and $u \in E_0$ is solution of (1.3). Moreover, $u_n \rightarrow u$ in $L^\theta(\mathbb{R}^N)$ for $p < \theta < p_s^*$. Noting $u_n \in \mathcal{M}_{\lambda_n}$ and $u \in \mathcal{N}$, we can obtain

$$\begin{aligned} & \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y) - u(x) + u(y)|^p}{|x - y|^{n+ps}} dx dy + \int_{\mathbb{R}^N} \lambda_n V(x) |u_n(x) - u(x)|^p dx \\ &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u_n(x) - u_n(y)|^p}{|x - y|^{n+ps}} dx dy + \int_{\mathbb{R}^N} \lambda_n V(x) |u_n(x)|^p dx - \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{|u(x) - u(y)|^p}{|x - y|^{n+ps}} dx dy - \int_{\mathbb{R}^N} \lambda_n V(x) |u(x)|^p dx \\ &+ o(1) \\ &= \int_{\mathbb{R}^N} u_n^+(x)^q dx - \int_{\Omega} u^+(x)^q dx + o(1) = o(1). \end{aligned}$$

Thus, $u_n \rightarrow u$ in $W^{s,p}(\mathbb{R}^N)$. This completes the proof of Theorem 1.4. \square

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References

- [1] T. Bartsch and Z. Wang, *Existence and multiplicity results for some superlinear elliptic problems on \mathbb{R}^N* , Comm. Partial Differential Equations **20** (1995), no. 9–10, 1725–1741, DOI: <https://doi.org/10.1080/03605309508821149>.
- [2] T. Bartsch and M. Parnet, *Nonlinear Schrödinger equations near an infinite well potential*, Calc. Var. Partial Differential Equations **51** (2014), 363–379, DOI: <https://doi.org/10.1007/s00526-013-0678-5>.
- [3] L. Xu and H. Chen, *Nontrivial solutions for Kirchhoff-type problems with a parameter*, J. Math. Anal. Appl. **433** (2016), 455–472, DOI: <https://doi.org/10.1016/j.jmaa.2015.07.035>.
- [4] C. O. Alves, A. B. Nobrega, and M. Yang, *Multi-bump solutions for Choquard equation with deepening potential well*, Calc. Var. Partial Differential Equations **55** (2016), 48, DOI: <https://doi.org/10.1007/s00526-016-0984-9>.
- [5] L. Zhao, H. Liu, and F. Zhao, *Existence and concentration of solutions for the Schrödinger-Poisson equations with steep well potential*, J. Differential Equations **255** (2013), 1–23, DOI: <https://doi.org/10.1016/j.jde.2013.03.005>.
- [6] M. Clapp and Y. Ding, *Positive solutions of a Schrödinger equation with critical nonlinearity*, Z. Angew. Math. Phys. **55** (2004), 592–605, DOI: <https://doi.org/10.1007/s00033-004-1084-9>.
- [7] M. Niu and Z. Tang, *Least energy solutions for nonlinear Schrödinger equation involving the fractional Laplacian and critical growth*, Discrete Contin. Dyn. Syst. **37** (2017), 3963–3987, DOI: <https://doi.org/10.3934/dcds.2017168>.
- [8] C. O. Alves and L. M. Barros, *Existence and multiplicity of solutions for a class of elliptic problem with critical growth*, Monatsh. Math. **187** (2018), 195–215, DOI: <https://doi.org/10.1007/s00605-017-1117-z>.
- [9] D. Applebaum, *Lévy processes – from probability to finance and quantum groups*, Notices Amer. Math. Soc. **51** (2004), 1336–1347.
- [10] A. Bahrouni, *Trudinger-Moser type inequality and existence of solution for perturbed nonlocal elliptic operators with exponential nonlinearity*, Commun. Pure Appl. Anal. **16** (2017), 243–252, DOI: <https://doi.org/10.3934/cpaa.2017011>.
- [11] A. Bahrouni, S. Bahrouni, and M. Xiang, *On a class of nonvariational problems in fractional Orlicz-Sobolev spaces*, Nonlinear Anal. **190** (2020), 111595, DOI: <https://doi.org/10.1016/j.na.2019.111595>.

- [12] J. L. Bona and Y. A. Li, *Decay and analyticity of solitary waves*, J. Math. Pures Appl. **76** (1997), no. 5, 377–430, DOI: [https://doi.org/10.1016/S0021-7824\(97\)89957-6](https://doi.org/10.1016/S0021-7824(97)89957-6).
- [13] A. de Bouard and J. C. Saut, *Symmetries and decay of the generalized Kadomtsev-Petviashvili solitary waves*, SIAM J. Math. Anal. **28** (1997), no. 5, 1064–1085, DOI: <https://doi.org/10.1137/S0036141096297662>.
- [14] X. Cabré and J. Tan, *Positive solutions of nonlinear problems involving the square root of the Laplacian*, Adv. Math. **224** (2010), no. 5, 2052–2093, DOI: <https://doi.org/10.1016/j.aim.2010.01.025>.
- [15] L. Caffarelli and L. Silvestre, *An extension problem related to the fractional Laplacian*, Comm. Partial Differential Equations **32** (2007), no. 8, 1245–1260, DOI: <https://doi.org/10.1080/03605300600987306>.
- [16] M. Cheng, *Bound state for the fractional Schrödinger equation with unbounded potential*, J. Math. Phys. **53** (2012), 043507, DOI: <https://doi.org/10.1063/1.3701574>.
- [17] J. Dávila, M. del Pino, and J. Wei, *Concentrating standing waves for the fractional nonlinear Schrödinger equation*, J. Differential Equations **256** (2014), no. 2, 858–892, DOI: <https://doi.org/10.1016/j.jde.2013.10.006>.
- [18] S. Dipierro, G. Palatucci, and E. Valdinoci, *Existence and symmetry results for a Schrödinger type problem involving the fractional Laplacian*, Matematiche **68** (2013), 201–216, DOI: <https://doi.org/10.4418/2013.68.1.15>.
- [19] P. Felmer, A. Quaas, and J. Tan, *Positive solutions of the nonlinear Schrödinger equation with the fractional Laplacian*, Proc. Roy. Soc. Edinburgh Sect. A **142** (2012), no. 6, 1237–1262, DOI: <https://doi.org/10.1017/S0308210511000746>.
- [20] R. L. Frank and E. Lenzmann, *Uniqueness of nonlinear ground states for fractional Laplacians in \mathbb{R}* , Acta Math. **210** (2013), no. 2, 261–318, DOI: <https://doi.org/10.1007/s11511-013-0095-9>.
- [21] T. Jin, Y. Li, and J. Xiong, *On a fractional Nirenberg problem, part I: blow up analysis and compactness of solutions*, J. Eur. Math. Soc. **16** (2014), no. 6, 1111–1171, DOI: <https://doi.org/10.4171/JEMS/456>.
- [22] M. Maris, *On the existence, regularity and decay of solitary waves to a generalized Benjamin-Ono equation*, Nonlinear Anal. **51** (2002), no. 6, 1073–1085, DOI: [https://doi.org/10.1016/S0362-546X\(01\)00880-X](https://doi.org/10.1016/S0362-546X(01)00880-X).
- [23] M. Niu and Z. Tang, *Least energy solutions of nonlinear Schrödinger equations involving the fractional Laplacian and potential wells*, Sci. China Math. **60** (2017), 261–276, DOI: <https://doi.org/10.1007/s11425-015-0830-3>.
- [24] J. Tan and J. Xiong, *A Harnack inequality for fractional Laplace equations with lower order terms*, Discrete Contin. Dyn. Syst. **31** (2011), no. 3, 975–983, DOI: <https://doi.org/10.3934/dcds.2011.31.975>.
- [25] P. L. Lions, *The concentration-compactness principle in the calculus of variations. The locally compact case, Part I*, Ann. Inst. H. Poincaré Anal. Non Linéaire **1** (1984), no. 2, 109–145, DOI: [https://doi.org/10.1016/S0294-1449\(16\)30428-0](https://doi.org/10.1016/S0294-1449(16)30428-0).