Research Article

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One-dimensional constrained inversion study of TEM and application in coal goafs' detection

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Abstract: Transient electromagnetic method (TEM) is often used in urban underground space exploration and field geological resource detection. Inversion is the most important step in data interpretation. Because of the volume effect of the TEM, the inversion results are usually multi-solvable. To reduce the multi-solvability of inversion, the constrained inversion of TEM has been studied using the least squares method. The inversion trials were performed using two three-layer theoretical geological models and one four-layer theoretical geological model. The results show that onedimensional least squares constrained inversion is faster and more effective than unconstrained inversion. The induced electromotive force attenuation curves of the inversion model indicate that the same attenuation curve may be used for different geological conditions. Therefore, constrained inversion using known geological information can more accurately reflect the underground geological information.

Keywords: transient electromagnetic method, constrained inversion, the least squares method, 1D TEM forward simulation, induced emf

1 Introduction

Geophysics methods have been widely used to detect the mineral resources [1–5], geological structures [6,7], and

urban underground space construction [8–11]. As the safety of coal mines is increasingly valued by people, it is necessary to detect the water-rich areas around the working face before coal mining, such as goafs and water-conducting structures. As one of the geophysical exploration methods, the transient electromagnetic method (TEM), on the strength of its high efficiency, is widely used in advanced detection of tunnels, accurate positioning of underground pipelines, and detection of holes underground.

At present, the commonly used processing methods are calculating apparent resistivity [12–15], drawing multi-track curves of induced electromotive force (emf) [16] and one-dimensional (1D) inversion [17–19]. The volume effect of TEM makes it difficult for the apparent resistivity and multi-track curves to identify the edge and depth of the target. The general inversion methods without geological information are marked by slow speed, low accuracy, and multi-solvability.

Before proceeding with exploration, we are usually equipped with some underground information, such as the electrical characteristics of the pipeline and the depth of the coal seam. The known information can be used in the transient electromagnetic data processing. Therefore, to improve the inversion accuracy, reduce the multisolvability, and increase the speed, we add known geological information to the initial model for constrained inversion. Compared with other inversion methods, the least squares inversion is much simpler with a high inversion speed, thus adopted for the study of the TEM 1D constraint inversion.

2 Basic principle of TEM and 1D forward simulation

2.1 Basic principle of TEM

TEM involves applying a current to power supply electrodes or transmitting loops, which causes it to form an electromagnetic field in the surrounding space, and then

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suddenly cutting off the current. The instantaneous disappearance of the electromagnetic field will generate the induced electromagnetic field in the surrounding conductor [20]. We obtain geological information by analyzing the inductive electromagnetic fields.

2.2 1D forward algorithm

The basis of inversion is forward simulation. The ground surface electromagnetic response of the horizontal stratum which is generated by step current of loop source can be obtained from the Maxwell equations. The time domain electromagnetic response of rectangular loop excitation source is shown below [21]:

$$\frac{\partial B}{\partial t} = R_1 + R_2,\tag{1}$$

where

$$R_{1} = -\frac{\ln 2}{t} \frac{I}{4\pi} \sum_{j=1}^{J} K_{j} \overline{z_{0}} \int_{-W}^{W} \int_{0}^{\infty} \lambda (1 + r_{\text{TE}})$$

$$\times \left[\frac{L - y}{\rho_{L}} J_{1}(\lambda \rho_{L}) + \frac{L + y}{\rho_{-L}} J_{1}(\lambda \rho_{-L}) \right] d\lambda dx',$$

$$\ln 2 I \int_{0}^{L} \int_{0}^{\infty} \int_{0}^{L} dx'$$

$$R_{2} = -\frac{\ln 2}{t} \frac{I}{4\pi} \sum_{j=1}^{J} K_{j} \overline{z_{0}} \int_{-L}^{L} \int_{0}^{\infty} \lambda (1 + r_{\text{TE}})$$

$$\times \left[\frac{W - x}{\rho_{W}} J_{1}(\lambda \rho_{W}) + \frac{W + x}{\rho_{-W}} J_{1}(\lambda \rho_{-W}) \right] d\lambda dy',$$

where λ is the space wave number; K_j is the G–S algorithm coefficient (can be obtained from equation (2) [22]); W is the half-width and L is the half length of the rectangular loop in meters; J_1 is the first-order Bessel function; $\overline{z_0} = i\omega\mu_0$ and $\mu_0 = 4\pi \times 10^{-7}$ show the vacuum permeability, ω angular frequency; r_{TE} is the reflection coefficient (can be obtained from equation (3)); x' and y' are the coordinates of each point on the rectangular emission loop; I is the current; and t is the time in seconds. Solving these expressions require the Hankel transformation [23,24]. Furthermore, ρ_L , ρ_{-L} , ρ_W , and ρ_{-W} are shown below:

$$\begin{split} & \rho_L = [(x'-x)^2 + (L-y)^2]^{1/2} \\ & \rho_{-L} = [(x'-x)^2 + (-L-y)^2]^{1/2} \\ & \rho_W = [(y'-y)^2 + (W-x)^2]^{1/2} \\ & \rho_{-W} = [(y'-y)^2 + (-W-x)^2]^{1/2}, \end{split}$$

$$K_{j} = (-1)^{j+M} \times \sum_{k=m}^{\min(j,M)} \frac{k^{M}(2k)!}{(M-k)!k!(k-1)!(j-k)!(2k-j)!},$$
(2)

where M = J/2; m is the integer of (j + 1)/2; and J is the number of G–S coefficient.

$$r_{\rm TE} = \frac{Y_0 - \overline{Y}_1}{Y_0 + \overline{Y}_1},\tag{3}$$

where $Y_0 = \frac{u_0}{i\omega\mu_0}$. $\overline{Y_1}$ can be obtained from the following recursive formula:

$$\overline{Y_n} = Y_n \frac{\overline{Y_{n+1}} + Y_n \tanh(u_n h_n)}{Y_n + \overline{Y_{n+1}} \tanh(u_n h_n)}$$

$$\vdots$$

$$\overline{Y_n} = Y_N$$

where $Y_n = \frac{u_n}{\overline{Z_n}}$; $\overline{Z} = i\omega\mu_n$ (a total of N layers in the model and n represents the nth layer); $u_n = (\lambda^2 + k_n^2)^{1/2}$; $k_n^2 = \omega^2\mu_n\varepsilon_n - i\omega\mu_n\sigma_n$; μ_n , ε_n , and σ_n are the magnetic permeability, dielectric constant, and electrical conductivity of the nth layer; and h_n is the thickness of the nth layer.

To verify the accuracy of the above algorithm, comparison with the analytical solution of uniform half space is now performed. The analytical formula of uniform half space is as follows [25]:

$$B_z = \frac{I\mu_0}{2a} \left[\frac{3}{\sqrt{\pi}\theta a} e^{-\theta^2 a^2} + \left[1 - \frac{3}{2\theta^2 a^2} \right] \operatorname{erf}(\theta a) \right], \quad (4)$$

where a is the radius of circular emission loop; $\theta = \sqrt{\frac{\mu_0}{4t\rho}}$; t is the time; $\rho = \frac{1}{\sigma}$ is the resistivity of uniform half space;

and
$$\operatorname{erf}(\theta a) = \frac{2}{\sqrt{\pi}} \int_{0}^{\theta a} e^{-t^2} dt$$
 is the error function.

From equation (4), we can obtain induced emf.

$$\frac{\partial B_z}{\partial t} = \frac{I\rho}{a^3} \left[3 \operatorname{erf}(\theta a) - \frac{2}{\sqrt{\pi}} \theta a \left[3 + 2\theta^2 a^2 \right] e^{-\theta^2 a^2} \right]. \quad (5)$$

The equivalent relationship between the radius a of circular emission loop and the side length 2W and 2L of rectangular emission loop is shown below.

$$a=\sqrt{\frac{4WL}{\pi}}.$$

Set up a uniform half-space model for comparison. The conductivity of uniform half space is $0.01\,\mathrm{S\,m^{-1}}$, the transmitting current is 1 ampere, and the length of transmitting is 100 m. The analytical solution of the induced emf curve is consistent with the 1D forward simulation result (Figure 1).

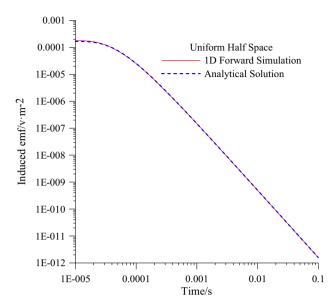


Figure 1: Comparison of analytical solution and 1D forward simulation.

3 Constrained inversion by the least squares method

3.1 Basic principle of constrained inversion

The inversion process of the least squares method iteratively calculates the modification by fitting the observation data, continuously obtains new model parameters, and finally obtains a model in which fit between the response data and the observation data meets the given accuracy. Then the model that fits the given accuracy is the inversion result. The least squares inversion method is the most commonly used inversion method. Its objective function is the sum of the squares of fitting error of the data and the sum of the squares of the model parameters. The iterative formula for the inversion is obtained by minimizing the objective function [26].

First, build an objective fitting function of 1D TEM inversion as shown below:

$$\varphi = \sum_{i=1}^{M} |u_{ai} - u_{ci}^{\star}|^2 = ||u - u^{\star}||_2^2 = (u - u^{\star})^T (u - u^{\star}), (6)$$

where M is the number of the measured data; u_{ai} is the measured induced emf; u_{ci}^{\star} is the theoretically calculated induced emf; u is the vector of measured data; and u^{\star} is the vector that is calculated by the forward model.

The theoretically calculated induced emf is approximated by Taylor's formula, omitting quadratic and above terms. Then the observation data and model approximate linear relationship are shown below:

$$u_{ci}^{\star} \approx u_{ci} + \sum_{i=1}^{N} \left(\frac{\partial u_{ci}}{\partial m_{j}} \Delta m_{j} \right),$$
 (7)

where N is the number of model parameters; u_{ci} is the calculated emf from initial model; m_j is the jth parameter; and Δm_j is the jth model modification. Then, write equation (7) as a matrix and substitute it into equation (6):

$$\varphi = (u - u_0 - J \cdot \Delta m)^T (u - u_0 - J \cdot \Delta m), \tag{8}$$

where J is the Jacobian matrix; Δm is the model modification vector; and u_0 is the emf vector that is calculated from the initial model.

$$J = \begin{bmatrix} \frac{\partial u_{c1}}{\partial m_1} & \dots & \frac{\partial u_{c1}}{\partial m_j} & \dots & \frac{\partial u_{c1}}{\partial m_N} \\ \vdots & \ddots & \vdots & \ddots & \vdots \\ \frac{\partial u_{ci}}{\partial m_1} & \dots & \frac{\partial u_{ci}}{\partial m_j} & \dots & \frac{\partial u_{ci}}{\partial m_N} \\ \vdots & \ddots & \vdots & \ddots & \vdots \\ \frac{\partial u_{cM}}{\partial m_1} & \dots & \frac{\partial u_{cM}}{\partial m_j} & \dots & \frac{\partial u_{cM}}{\partial m_N} \end{bmatrix},$$

$$\Delta m = \begin{bmatrix} \Delta m_1 \\ \vdots \\ \Delta m_j \\ \vdots \\ \Delta m_N \end{bmatrix} \quad u_0 = \begin{bmatrix} u_{c1} \\ \vdots \\ u_{ci} \\ \vdots \\ u_{cM} \end{bmatrix} \quad u = \begin{bmatrix} u_{a1} \\ \vdots \\ u_{ai} \\ \vdots \\ u_{aM} \end{bmatrix}$$

If the *j*th parameter is known, we can add the known information to the Jacobian matrix. The elements of the *j*th column of the *J* matrix are all 0. Therefore, the *J* matrix can be rewritten as below:

$$J = \begin{bmatrix} \frac{\partial u_{c1}}{\partial m_1} & \cdots & \frac{\partial u_{c1}}{\partial m_{j-1}} & 0 & \frac{\partial u_{c1}}{\partial m_{j+1}} & \cdots & \frac{\partial u_{c1}}{\partial m_N} \\ \vdots & & & \vdots & & \vdots \\ \frac{\partial u_{ci}}{\partial m_1} & \cdots & \frac{\partial u_{ci}}{\partial m_{j-1}} & 0 & \frac{\partial u_{ci}}{\partial m_{j+1}} & \cdots & \frac{\partial u_{ci}}{\partial m_N} \\ \vdots & & & & \vdots & & \ddots & \vdots \\ \frac{\partial u_{cM}}{\partial m_1} & \cdots & \frac{\partial u_{cM}}{\partial m_{j-1}} & 0 & \frac{\partial u_{cM}}{\partial m_{j+1}} & \cdots & \frac{\partial u_{cM}}{\partial m_N} \end{bmatrix}$$

Let $(u-u_0)=\bar{u}$, to minimize the objection (equation (8)), then $\frac{\partial \varphi}{\partial \Delta m}=0$, and

$$\Delta m = (J^T J)^{-1} J^T \bar{u}. \tag{9}$$

The model modification Δm will be obtained by solving the equation (9) by singular value decomposition. Then the new model can be derived with Δm .

Use the relative mean square deviation to measure the accuracy of the fit as below:

RMSD =
$$\sqrt{\frac{1}{M} \sum_{i=1}^{M} \left(\frac{u_{ai} - u_{ci}}{u_{ai}} \right)^2}$$
. (10)

3.2 Inversion and analysis of theoretical data

To verify the feasibility of the TEM least squares constrained inversion, we use two three-layer (H-type and K-type) geoelectric models, a four-layer (HK-type) geoelectric model, and a five-layer model as examples for inversion trial calculation (Figure 2). Parameters of three models are shown in Table 1 (C is the conductivity and T is the thickness). The length and width of rectangular transmitting loop are all 200 m. The transmitting current is 1 ampere.

Figure 3 shows the inversion result models and the theoretical model of H-type. It can be observed from Figure 3(a) that the constrained inversion and unconstrained inversion result models of H-type geoelectric model are consistent with the theoretical model. Errors of the induced emf curves are all within 1%. However, the unconstrained inversion took 2,564 s and 17 iterations, and the constrained inversion took 1,367 s and 13 iterations.

Figure 4 is the inversion result model and the theoretical model of K-type. The K-type model inversion result shows that the result model of constrained inversion is closer to the theoretical model than unconstrained

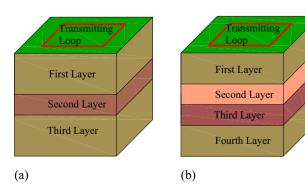
inversion result model, and the constrained inversion took only 950 s and nine iterations, whereas unconstrained inversion took 1,370 s and nine iterations. Moreover, the second layer of unconstrained inversion model of K-type has larger conductivity and is thicker than the theoretical model, even though they have the same induced emf attenuation curves.

Figure 5 shows the inversion result models and the theoretical model of HK-type. It still shows that constrained inversion is better than unconstrained inversion. Although their attenuation curves are the same, unconstrained inversion cannot invert the third layer, but constrained inversion can. Constrained inversion took 4,819 s and 26 iterations, whereas unconstrained inversion took 6,412 s and 26 iterations.

Figure 6 shows the inversion result models and the theoretical model of HKH-type. It still shows that constrained inversion is better than unconstrained inversion. Although their attenuation curves are the same, unconstrained inversion can only invert the first high-conductivity layer, but not the second high-conductivity layer, but after constraining the first high-conductivity layer, the second high-conductivity layer can be well reversed.

4 Instance analysis

In mining areas where the coal seam is shallow, the randomness of small coal kilns gives rise to many unknown goafs. The presence of these goafs poses a great threat to the safety of coal mine production and compromises the stability of ground buildings and ground engineering structures. Usually, goafs and coal seam have the same depth and thickness. Therefore, the thickness of both the cover layer and the target area can be restricted during the inversion.



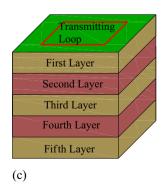


Figure 2: The 1D models of TEM forward simulation. (a) H or K model; (b) HK model; and (c) HKH model.

Table 1: The parameters of geoelectric models

	H-type		K-type		HK-type		HKH-type	
	<i>C</i> (S m ⁻¹)	<i>T</i> (m)	<i>C</i> (S m ⁻¹)	<i>T</i> (m)	<i>C</i> (S m ⁻¹)	<i>T</i> (m)	<i>C</i> (S m ⁻¹)	<i>T</i> (m)
First layer	0.01	100	0.1	100	0.01	100	0.01	50
Second layer	0.1	50	0.01	50	0.1	20	0.1	50
Third layer	0.01		0.1		0.005	30	0.01	50
Fourth layer					0.01		0.1	50
Fifth layer							0.01	

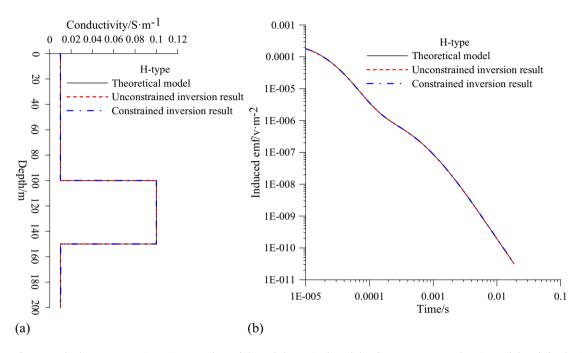


Figure 3: The least squares inversion result models and theoretical models of H-type. (a) Geoelectric model and (b) the induced emf attenuation curves.

We used TEM to detect shallow goafs in a mining area in Inner Mongolia, China. The side length of the transmitting loop is $20 \times 20\,\mathrm{m}$, and the emission current is 18.6 ampere. The instrument is a transient electromagnetic instrument named Tsikl-7. The observation data are normalized by receiving coil area and current. According to known geological data, the average thickness of the coal seam in the locality is 15.6 m, and the average depth is 66.4 m. Therefore, we limited the burial depth and thickness of the target layer during the inversion. However, because it is not known whether water is available, there is no restriction on resistivity.

Figure 7 is the inversion result of one of the observation points. In Figure 7(a), there are two conductive layers in the result, the first being a shallow aquifer at 10–20 m underground, and the second 66.4–80 m underground, which is caused by water accumulation in the

goaf. Because the coal mine goaf is rich in sulfides and metal ions, the conductivity of the water-containing goaf is much greater than that of the shallow aquifer. Figure 7(b) represents the comparison between the attenuation curves of the inversion result data and measured data. Because of various noise in actual field work, the inversion curve cannot fit the measured curve as well as the theoretical model.

5 Conclusion

From the present research work, we can arrive at the following conclusions:

 The 1D least squares constrained inversion result model of TEM is better than the unconstrained

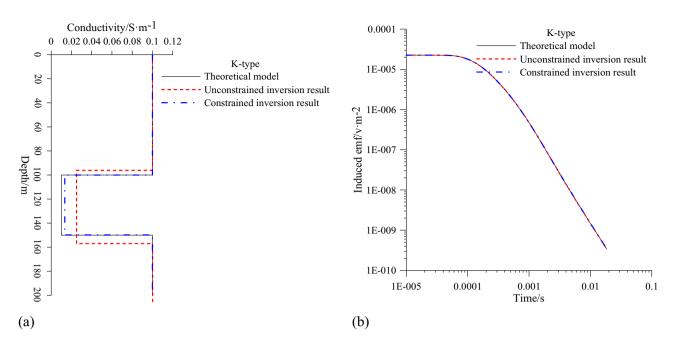


Figure 4: The least squares inversion result models and theoretical models of K-type. (a) Geoelectric model and (b) the induced emf attenuation curves.

inversion on the grounds that it takes less time and iterations, and the inversion results can better reflect the real situation underground.

The induced emf attenuation curves of the constrained inversion result model and unconstrained inversion result model are consistent with the theoretical model's attenuation curve, and the errors are all within 1%, indicating that different inversion models have the same attenuation curves. Therefore, constrained inversion under known geological information can reduce multi-solvability.

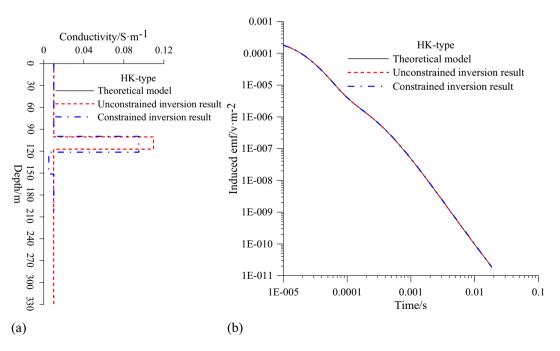


Figure 5: The least squares inversion result models and theoretical models of HK-type. (a) Geoelectric model and (b) the induced emf attenuation curves.

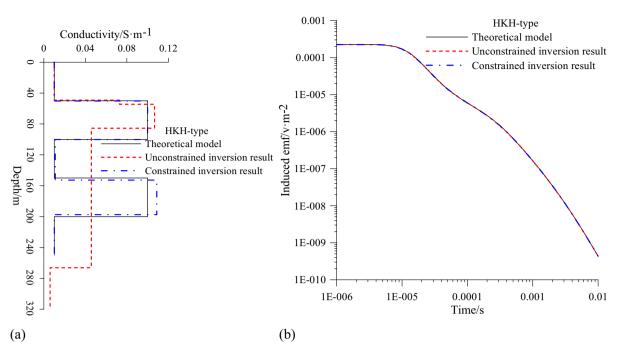


Figure 6: The least squares inversion result models and theoretical models of HKH-type. (a) Geoelectric model and (b) the induced emf attenuation curves.

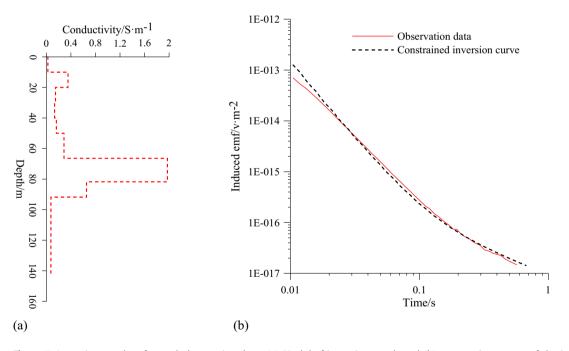


Figure 7: Inversion results of actual observation data. (a) Model of inversion result and (b) attenuation curves of the inversion model and measured data.

3. The actual observation data show that when the depth and thickness of a certain layer are constrained, the resistivity and shallow electrical information of the modified position can still be effectively reversed under the influence of noise. **Acknowledgments:** The authors thank the reviewers and editors for their valuable advice and help. This research was funded by the Nature Science Foundation of China (Grant No. 41674133) and National Science and Technology Support Program of China (Grant No. 2013BAK06B01).

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