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**NECESSARY AND SUFFICIENT CONDITIONS FOR THE EXISTENCE
OF LINEAR CONNECTION DETERMINED BY THE FIELD
OF A TENSOR OF TYPE (1,1) COVARIANTLY CONSTANT
IN THE THREE DIMENSIONAL SPACE**

1. Introduction

Let M be a three-dimensional differential manifold and ξ a chart belonging to its atlas $\text{atl } M$. By $\mathcal{F}(M)$, resp. $\mathcal{X}(M)$, we denote respectively the ring of smooth functions and the module of smooth vector field on M .

Let T be a vector field of type (1,1) and ∇ a linear connection on M . Let V be a vector field belonging to $\mathcal{X}(M)$. Then $\nabla_V T$ denotes the covariant derivative of the given vector field T along the vector field V .

In the present paper we shall give a necessary and sufficient condition for the existence of a linear connection ∇ determined by the field of a tensor T covariantly constant, i.e. satisfying the condition

$$(1.1) \quad \nabla_V T = 0$$

for every $V \in \mathcal{X}(M)$. We call such a tensor field parallel with respect to the linear connection ∇ .

Condition (1.1) can be written out with respect to an arbitrary chart $\xi \in \text{atl } M$

$$(1.2) \quad \nabla_{\sigma} T_{\lambda}^{\mu} := \partial_{\sigma} T_{\lambda}^{\mu} + \Gamma_{\sigma\alpha}^{\mu} T_{\lambda}^{\alpha} - \Gamma_{\sigma\lambda}^{\alpha} T_{\alpha}^{\mu} = 0$$

$$(\sigma, \alpha, \lambda, \mu = 1, 2, 3),$$

where T_{λ}^{μ} are given coordinates of the tensor field T in the chart ξ , $\partial_{\sigma} T_{\lambda}^{\mu}$ are partial derivatives of the field T and $\Gamma_{\sigma\alpha}^{\beta}$ are unknown coordinates of the linear connection in the given chart. (1.2) can be treated as a system of n^2 equations (for every σ) linear non-homogeneous with respect to $\Gamma_{\sigma\alpha}^{\beta}$. This problem has been solved for 2-dimensional space in [1] for 3-dimensional space in [2] (not published doctoral thesis of the author). In the present paper we give main stages of this solution.

2. Some scalar concomitants of the tensor field T

Let \mathcal{T} denote the set of all non-zero tensor fields of type (1.1) defined on M .

We shall consider scalar fields which are concomitants of $T \in \mathcal{T}$ [3]

$$(2.1) \quad \begin{cases} (1) \\ S(T) := T_{\alpha}^{\alpha}, S^{(2)}(T) := T_1^1 T_2^2 - T_1^2 T_2^1 + T_1^1 T_3^3 - T_1^3 T_3^1 + T_2^2 T_3^3 - T_2^3 T_3^2, \\ (3) \\ S(T) := \det(T_{\lambda}^{\mu}) = 3! T_{[1}^1 T_2^2 T_3^3], \end{cases}$$

(1) $S(T)$ is called the trace of first order of the matrix (T_{λ}^{μ}) of the tensor field T and is denoted by $\text{tr } T$, $S^{(2)}(T)$ is called the trace of second order of the matrix (T_{λ}^{μ}) and it is seen from the definition that $S^{(3)}(T) = \det(T_{\lambda}^{\mu})$.

The traces $S^{(i)}(T)$ ($i = 1, 2, 3$) are scalar fields (1) $S(T) \in \mathcal{F}(M)$.

The partial derivatives of these scalar field

$$(2.2) \quad U_{\sigma}^{(i)}(T) := \partial_{\sigma} S^{(i)}(T) \quad (i = 1, 2, 3; \sigma = 1, 2, 3).$$

are co-vector fields. The vanishing of these fields, i.e.

(i) $U_6(T) = 0$ is an invariant property holding in every chart $\xi \in \text{atl } M$.

Let T_Q denote T_λ^μ , where Q is a joint index: $Q = (\mu, \lambda)$, ($Q = 1, 2, \dots, 9$). By (2.1) and (2.2) we have

$$(2.3) \quad U^{(i)}(T) = \partial_6^{(i)} S^{(i)}(T) = \sum_{Q=1}^9 \frac{\partial S^{(i)}(T)}{\partial T_Q} \partial_6 T_Q \quad (i = 1, 2, 3).$$

Let us consider the matrix

$$(2.4) \quad A(T) := \left(\frac{\partial S^{(i)}(T)}{\partial T_Q} \right) \quad (i = 1, 2, 3; \quad Q = 1, 2, \dots, 9),$$

where i denotes the row and Q the column of the matrix $A(T)$. Let

$$(2.5) \quad r_{A(T)} := \text{the rank of } A(T).$$

It is known that $r_{A(T)}$ is an absolute invariant.

Further, let r_A be the maximal rank of the matrix $A(T)$ where T runs through \mathcal{T} , i.e.

$$(2.6) \quad r_A = \max_T r_{A(T)}.$$

From a theorem in [3] it follows that

Theorem 1. Under the above assumptions

$$r_A = 3.$$

Moreover, let us consider the following scalar fields which are concomitants of given tensor field $T \in \mathcal{T}$ [2]

$$(2.7) \quad \begin{aligned} (1) \quad \phi^{(1)}(T) &:= T_\alpha^\alpha, & (2) \quad \phi^{(2)}(T) &:= T_\alpha^\beta T_\beta^\alpha, \\ (3) \quad \phi^{(3)}(T) &:= T_\alpha^\beta T_\beta^\gamma T_\gamma^\alpha. \end{aligned}$$

Partial derivatives of these scalar fields

$$(2.8) \quad \overset{(i)}{V}_\sigma(T) = \partial_\sigma \overset{(i)}{\phi}(T) \quad (i = 1, 2, 3),$$

are covariant vector fields. The vanishing of these fields, i.e. $\overset{(i)}{V}_\sigma(T) = 0$ is an invariant property holding in every chart $\xi \in \text{atl } M$.

From (2.7) and (2.8) we obtain

$$(2.9) \quad \overset{(i)}{V}_\sigma(T) = \partial_\sigma \overset{(i)}{\phi}(T) = \sum_{Q=1}^9 \frac{\partial \overset{(i)}{\phi}(T)}{\partial T_Q} \partial_\sigma T_Q \quad (i = 1, 2, 3),$$

where T_Q denotes the coordinates of the tensor T^μ_λ .

Now let us form the matrix

$$(2.10) \quad B(T) = \left(\frac{\partial \overset{(i)}{\phi}(T)}{\partial T_Q} \right) \quad \begin{matrix} (i = 1, 2, 3; \\ Q = 1, 2, \dots, 9), \end{matrix}$$

where i denotes the row and Q the column of the matrix $B(T)$.

Let $r_B(T)$ denote the rank of this matrix

$$(2.11) \quad r_B(T) = \text{rank } B(T) \quad \text{for } T \in J.$$

The maximal rank of $B(T)$ will be denoted by r_B :

$$(2.12) \quad r_B = \max_T r_B(T).$$

In [2] it was proved that the scalar field $\overset{(i)}{\phi}(T)$ defined by (2.7) can be expressed by means of the traces $\overset{(i)}{S}(T)$ defined in (2.1), in the following way:

$$(2.13) \quad \begin{cases} \begin{pmatrix} 1 \\ \phi \end{pmatrix} = \begin{pmatrix} 1 \\ S \end{pmatrix} (= T_{\alpha}^{\alpha}) \\ \begin{pmatrix} 2 \\ \phi \end{pmatrix} = \begin{pmatrix} 1 \\ S^2 \end{pmatrix} - 2 \begin{pmatrix} 2 \\ S \end{pmatrix} \\ \begin{pmatrix} 3 \\ \phi \end{pmatrix} = -3 \begin{pmatrix} 1 \\ S \end{pmatrix} \begin{pmatrix} 2 \\ S \end{pmatrix} + 3 \begin{pmatrix} 3 \\ S \end{pmatrix} + S \begin{pmatrix} 1 \\ 3 \end{pmatrix}. \end{cases}$$

The mapping $S \rightarrow \phi$ defined above, where $S = \begin{pmatrix} 1 & 2 & 3 \\ S & S & S \end{pmatrix}$, $\phi = \begin{pmatrix} 1 & 2 & 3 \\ \phi & \phi & \phi \end{pmatrix}$ is a bijection. The inverse mapping can be expressed as follows

$$(2.14) \quad \begin{cases} \begin{pmatrix} 1 \\ S \end{pmatrix} = \begin{pmatrix} 1 \\ \phi \end{pmatrix} \\ \begin{pmatrix} 2 \\ S \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 \\ \phi \end{pmatrix}^2 - \frac{1}{2} \begin{pmatrix} 2 \\ \phi \end{pmatrix} \\ \begin{pmatrix} 3 \\ S \end{pmatrix} = \frac{1}{6} \begin{pmatrix} 1 \\ \phi \end{pmatrix}^3 - \frac{1}{2} \begin{pmatrix} 1 \\ \phi \end{pmatrix} \begin{pmatrix} 2 \\ \phi \end{pmatrix} + \frac{1}{3} \begin{pmatrix} 3 \\ \phi \end{pmatrix}. \end{cases}$$

From Theorem 1 and from the fact that the mapping $S \rightarrow \phi$ is a bijection it follows that

Theorem 2. Under the assumptions above we have

$$r_B = r_A = 3.$$

3. The matrices of the system of equations (1.2)

In the system of equations (1.2) we take the pairs (μ, λ) in the following order: (11), (12), (13), (21), (22), (23), (31), (32), (33), and for the unknowns $\Gamma_{6\alpha}^{\beta}$ the pairs (α, β) in the same order. Let $\mathcal{M}(T)$ denote the fundamental matrix of the system (1.2) and its elements by $M_{\lambda\beta}^{\alpha\mu}$, where the pair (μ, λ) denotes the row, (α, β) the column of the matrix \mathcal{M} . From (1.2) we obtain

$$(3.1) \quad M_{\lambda\beta}^{\alpha\mu}(T) = \delta_{\lambda}^{\alpha} T_{\beta}^{\mu} - \delta_{\beta}^{\mu} T_{\lambda}^{\alpha}.$$

From this formula it is seen that the matrix $\mathcal{M}(T)$ is skew-symmetric with respect to the pairs (μ, λ) and (α, β) , i.e. $M_{\lambda}^{\alpha}{}_{\beta}^{\mu}(T) = -M_{\beta}^{\mu}{}_{\lambda}^{\alpha}(T)$. Writing out the augmented matrix $\mathcal{M}_6^*(T)$ of the system (1.2) we obtain from (3.1):

$$(3.2) \quad \mathcal{M}_6^*(T) := \begin{bmatrix} 0 & T_2^1 & T_3^1 & -T_1^2 & 0 & 0 & -T_1^3 & 0 & 0 & \partial_6 T_1^1 \\ -T_2^1 & 0 & 0 & T_1^1 - T_2^2 & T_2^1 & T_3^1 & -T_2^3 & 0 & 0 & \partial_6 T_2^1 \\ -T_3^1 & 0 & 0 & T_3^2 & 0 & 0 & T_1^1 - T_3^3 & T_2^1 & T_3^1 & \partial_6 T_3^1 \\ T_1^2 & T_2^2 - T_1^1 & T_3^2 & 0 & -T_1^2 & 0 & 0 & -T_1^3 & 0 & \partial_6 T_1^2 \\ 0 & -T_2^1 & 0 & T_1^2 & 0 & T_3^2 & 0 & -T_2^3 & 0 & \partial_6 T_2^2 \\ 0 & -T_3^1 & 0 & 0 & -T_3^2 & 0 & T_1^2 & T_2^2 - T_3^3 & T_3^2 & \partial_6 T_3^2 \\ T_1^3 & T_2^3 & T_3^3 - T_1^1 & 0 & 0 & -T_1^2 & 0 & 0 & -T_1^3 & \partial_6 T_1^3 \\ 0 & 0 & -T_2^1 & T_1^3 & T_2^3 & T_3^3 - T_2^2 & 0 & 0 & -T_2^3 & \partial_6 T_2^3 \\ 0 & 0 & -T_3^1 & 0 & 0 & -T_3^2 & T_1^3 & T_2^3 & 0 & \partial_6 T_3^3 \end{bmatrix}$$

The basic matrix $\mathcal{M}(T)$ consists of the first nine columns of the matrix $\mathcal{M}_6^*(T)$. Let us introduce the following notation

$$(3.3) \quad r_{\mathcal{M}(T)} := \text{the rank of } \mathcal{M}(T), \quad r_{\mathcal{M}_6^*(T)} := \text{the rank of } \mathcal{M}_6^*(T).$$

Since the matrix $\mathcal{M}(T)$ is skew symmetric, we have

$$r_{\mathcal{M}(T)} \neq 1, 3, 5, 7, 9.$$

In [2] it is proved that for the non-zero tensor field T_{λ}^{μ} the rank $r_{\mathcal{M}(T)}$ of the matrix $\mathcal{M}(T)$ can assume the following values

$$(3.4) \quad r_{\mathcal{M}(T)} = 0, 4, 6,$$

and the following equivalences hold

$$(3.5) \quad \left\{ \begin{array}{l} \text{a) } T_{\lambda}^{\mu} = \alpha \delta_{\lambda}^{\mu} (\alpha \neq 0) \Leftrightarrow r_B = 1 (T_{\lambda}^{\mu} \neq 0) \\ \text{b) } r_B = 1 (T_{\lambda}^{\mu} \neq 0) \Leftrightarrow r_m = 0 (T_{\lambda}^{\mu} \neq 0) \\ \text{c) } r_B = 2 \Leftrightarrow r_m = 4 \\ \text{d) } r_B = 3 \Leftrightarrow r_m = 6, \end{array} \right.$$

where δ_{λ}^{μ} is the Kronecker delta.

4. The existence of linear connection

The following theorem holds:

Theorem 4. A necessary and sufficient condition in order that there exist a linear connection ∇ , with respect to which the tensor field $T \in \mathcal{T}$ is parallel, is the vanishing of three covector fields $\stackrel{(i)}{V}_6(T)$ defined in (2.8), i.e.

$$\stackrel{(i)}{V}_6(T) = 0 \quad (i = 1, 2, 3).$$

Outline of the proof. Necessity follows from the theorem proved by A. Lichnerowicz [7] and A. Zajtš [8]: if a tensor field T is parallel (covariantly constant) with respect to the connection ∇ and the field K is its concomitant, then K is also parallel. Hence in particular for the scalar concomitant $K(T)$ we have

$$\nabla_6 K(T) = 0 \Leftrightarrow \partial_6 K(T) = 0.$$

Thus for scalar concomitants (2.7) we have by (2.8):

$$\stackrel{(i)}{V}_6(T) = 0 \text{ for every } i = 1, 2, 3.$$

Conversely, assume that $\overset{(i)}{V}_6(T) = 0$ ($i = 1, 2, 3$). We have to show that $r_{\mathcal{M}(T)} = r_{\mathcal{M}_6^*(T)}$. From the assumption by (2.9) it follows that

$$(4.1) \quad \sum_{Q=1}^9 \frac{\partial \overset{(i)}{\phi}(T)}{\partial T_Q} \partial_\sigma T_Q = 0 \quad (i = 1, 2, 3).$$

a) In the case $r_{\mathcal{M}(T)} = 6$ we show that $B(T) \cdot \mathcal{M}_6^*(T) = [0]$ (see [6], [5]).

In fact, in the last column of the product $B(T) \cdot \mathcal{M}_6^*(T)$ we obtain

$$\sum_{Q=1}^9 \frac{\partial \overset{(i)}{\phi}(T)}{\partial T_Q} \partial_\sigma T_Q.$$

These expressions are equal by means of (4.1). For the remaining columns of this product we also obtain zero (see [6]). This implies that

$$r_{\mathcal{M}_6^*(T)} = r_{\mathcal{M}(T)} = 6.$$

b) For the case $r_{\mathcal{M}} = 4$ the proof is given in [2].

c) If $r_{\mathcal{M}(T)} = 0$, from (3.1) and (3.5) it follows that for the non-zero field T_λ^μ the following equivalences hold: $(r_{\mathcal{M}(T)} = 0) \Leftrightarrow (r_B(T) = 1) \Leftrightarrow T_\lambda^\mu = \alpha \delta_\lambda^\mu$, where $\alpha = \alpha(\xi^1, \xi^2, \xi^3)$ and δ_λ^μ is Kronecker's symbol. Then $\overset{(i)}{V}_6(T) = 0$ (for $i = 1, 2, 3$), i.e. $\overset{(i)}{V}_6(\alpha \delta_\lambda^\mu) = 0$ implies $\alpha = \text{const}$, and consequently $r_{\mathcal{M}_6^*(T)} = r_{\mathcal{M}(T)} = 0$ for $\alpha = \text{const}$.

From (3.4) we obtain the following corollary.

C o r o l l a r y 1. If $r_{\mathcal{M}(T)} = m$ (where $m = 0, 4, 6$) then for any non-zero field T with respect to $\Gamma_{\sigma\alpha}^B$ the solutions of the system (1.2) depend upon $3 \cdot (9-m)$ arbitrary functions.

Making use of Theorems 1 and 2 we can express Theorem 4 in the following equivalent form.

Theorem 5. A sufficient and necessary condition in order that there exist a linear connection ∇ with respect to which a non-zero tensor field $T \in \mathcal{T}$ is parallel is the vanishing of the covector fields $\stackrel{(1)}{U}_6(T)$ defined in (2.2), i.e.

$$\stackrel{(1)}{U}_6(T) = 0 \quad (i = 1, 2, 3).$$

5. The relation to the results of Y.C. Wong

Let M_3 be a three-dimensional connected differential manifold of class C^∞ with a non-zero smooth tensor field $T_\lambda^\mu(p)$, $p \in M_3$. Clearly, any system of 3rd linearly independent vectors lying in the space M_p tangent to the manifold M_3 at the point p is a basis of M_p .

The following theorem is a special case of the theorem proved by Y.C. Wong ([10], p.72): A necessary and sufficient condition in order that there exist a linear connection ∇ with respect to which the tensor field T_λ^μ is parallel (i.e. satisfying the system of equation (1.2)) is that at every point $p \in M_3$ there exist a basis (in general non-holonomic)

$$(5.1) \quad b_q^\sigma(p) \quad (q, \sigma = 1, 2, 3)$$

in which the coordinates of the tensor $T_\lambda^\mu(p)$ are constants C_λ^μ which are not simultaneously zero and such that the following identities hold:

$$(5.2) \quad T_\alpha^\beta(p) b_\beta^\mu(p) b_\lambda^{\alpha-1}(p) = C_\lambda^\mu,$$

where $b_\lambda^{\alpha-1}(p)$ is the system inverse to $b_\beta^\mu(p)$, i.e.

$$(5.3) \quad b_q^\mu(p) b_\lambda^{\alpha-1}(p) = \delta_\lambda^\mu.$$

From the above-mentioned theorem and from Theorem 4 we obtain the following corollary.

C o r o l l a r y. A necessary and sufficient condition in order that in a three-dimensional space for the tensor field $T_{\lambda}^{\mu}(p)$ at every point p there exist a basis (5.1) in which the coordinates of the tensor $T_{\lambda}^{\mu}(p)$ are constants C_{λ}^{μ} which are simultaneously non-zero and satisfy (5.2) and (5.3) is the vanishing of the covector fields $V_{\epsilon}^{(1)}(T)$, i.e. the relations $V_{\epsilon}^{(1)}(T) = 0$ ($i = 1, 2, 3$) (or in the equivalent form $U_{\epsilon}^{(1)}(T) = 0$, where $U_{\epsilon}^{(1)}(T)$ is defined in (2.2)).

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