

Research Article

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Limit Cycles Coming from Some Uniform Isochronous Centers

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Abstract: This article is about the weak 16th Hilbert problem, i.e. we analyze how many limit cycles can bifurcate from the periodic orbits of a given polynomial differential center when it is perturbed inside a class of polynomial differential systems. More precisely, we consider the uniform isochronous centers

$$\dot{x} = -y + x^2y(x^2 + y^2)^n, \quad \dot{y} = x + xy^2(x^2 + y^2)^n,$$

of degree $2n + 3$ and we perturb them inside the class of all polynomial differential systems of degree $2n + 3$. For $n = 0, 1$ we provide the maximum number of limit cycles, 3 and 8 respectively, that can bifurcate from the periodic orbits of these centers using averaging theory of first order, or equivalently Abelian integrals. For $n = 2$ we show that at least 12 limit cycles can bifurcate from the periodic orbits of the center.

Keywords: Periodic Solution, Uniform Isochronous Centers, Averaging Theory, Weak Hilbert Problem

MSC 2010: 37G15, 37C80, 37C30

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1 Introduction and Statement of the Main Results

The second part of the *16th Hilbert's problem* asks for the maximum number $H(n)$ of limit cycles that planar polynomial differential systems of degree n can have, see for instance [7, 8, 11], and the references quoted therein. The problem on the number $H(n)$ remains open, even for $n = 2$.

A weaker problem than the 16th Hilbert's problem, known now as the *weak 16th Hilbert's problem* was proposed by Arnold [2], who asked for the maximum number $Z(m, n)$ of isolated zeros of Abelian integrals of all polynomial 1-forms of degree n over algebraic ovals of degree m , for more details on the weak 16th Hilbert's problem see [4, 9, 19] and the hundreds of references quoted in these articles. Unfortunately the weak 16th Hilbert's problem is also extremely hard to study. On the other hand, the weak 16th Hilbert's problem is a particular case of the problem of studying the maximum number of limit cycles that can bifurcate from the periodic orbits of a center of a polynomial differential system of degree $m - 1$ when it is perturbed inside the class of all polynomial differential systems of degree n . Of course $Z(m, n) \leq H(\max(n, m - 1))$.

In this paper we provide lower bounds for the maximum number of limit cycles that can bifurcate from the periodic solutions of a polynomial differential uniform isochronous center of degree 3, 5 and 7 when it is perturbed inside the class of all polynomial differential systems of the same degree. The main result is based on the averaging theory of first order. But here the main work is to study the maximum number of simple zeros

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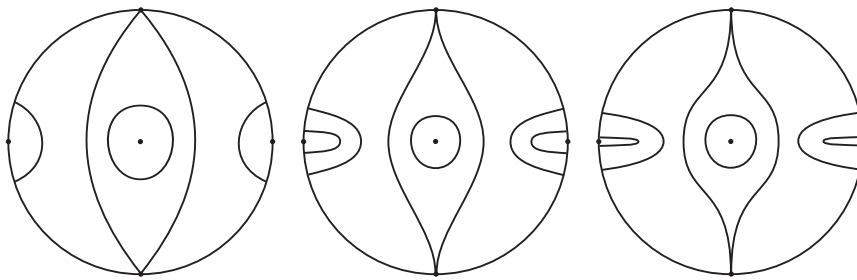


Figure 1. Phase portrait of the uniform isochronous center (1.1) for $n = 0$, $n = 1$, and $n = 2$, respectively.

of the obtained averaged functions, because not always the standard study of *Extended Chebyshev systems* (ET-systems) can be applied (see Appendix B). The study is based on some new results that can be applied when the family of functions that define $\mathcal{F}$ is not an ET-system. Some delicate study using qualitative theory on some differential equations is also needed to complete the study.

More precisely, we consider the polynomial differential system

$$\begin{cases} \dot{x} = -y + x^2y(x^2 + y^2)^n, \\ \dot{y} = x + xy^2(x^2 + y^2)^n, \end{cases} \quad (1.1)$$

of degree $2n + 3$ with $n \geq 0$, having a uniform isochronous center at the origin of coordinates, which in polar coordinates (r, θ) , where $x = r \sin \theta$ and $y = r \cos \theta$, becomes

$$\begin{cases} \dot{r} = r^{2n+3} \cos \theta \sin \theta, \\ \dot{\theta} = 1. \end{cases}$$

Since $\dot{\theta} = 1$, the origin of equation (1.1) is a uniform isochronous center, which taking as independent variable the variable θ writes

$$\frac{dr}{d\theta} = r' = r^{2n+3} \cos \theta \sin \theta.$$

An easy computation shows that the periodic solutions $r(\theta, r_0)$ surrounding the center $r = 0$ such that $r(0, r_0) = r_0$ are

$$r(\theta, r_0) = r_0(1 - (n+1)r_0^{2(n+1)} \sin^2 \theta)^{-\frac{1}{2n+2}}, \quad (1.2)$$

with $0 < r_0 < (n+1)^{-\frac{1}{2n+2}}$. The global phase portraits, in the Poincaré disc, of system (1.1) for $n = 0, 1, 2$ are shown in Figure 1.

Our purpose is to provide a lower bound for the maximum number of limit cycles that can bifurcate from the periodic solutions $r(\theta, r_0)$ surrounding the uniform isochronous center at $r = 0$ of degree 3, 5, 7 when we perturb it inside the class of all polynomial differential systems of degree 3, 5, 7, respectively. In other words, we study the number of limit cycles of the following three polynomial differential systems:

$$\begin{cases} \dot{x} = -y + x^2y + \varepsilon \sum_{i+j=0}^3 a_{ij}x^i y^j, \\ \dot{y} = x + xy^2 + \varepsilon \sum_{i+j=0}^3 b_{ij}x^i y^j, \end{cases} \quad (1.3)$$

$$\begin{cases} \dot{x} = -y + x^2y(x^2 + y^2) + \varepsilon \sum_{i+j=0}^5 a_{ij}x^i y^j, \\ \dot{y} = x + xy^2(x^2 + y^2) + \varepsilon \sum_{i+j=0}^5 b_{ij}x^i y^j, \end{cases} \quad (1.4)$$

$$\begin{cases} \dot{x} = -y + x^2y(x^2 + y^2)^2 + \varepsilon \sum_{i+j=0}^7 a_{ij}x^i y^j, \\ \dot{y} = x + xy^2(x^2 + y^2)^2 + \varepsilon \sum_{i+j=0}^7 b_{ij}x^i y^j, \end{cases} \quad (1.5)$$

where ε is a small parameter.

Our main result is the following.

Theorem 1.1. *For $|\varepsilon| \neq 0$ sufficiently small using averaging theory of first order we obtain that*

- (a) *system (1.3) can have up to 3 limit cycles and there are perturbations that only 0, 1, 2 and 3 limit cycles bifurcate from the center;*
- (b) *system (1.4) can have up to 8 limit cycles and there are perturbations that only 0, 1, 2, ..., 7 and 8 limit cycles bifurcate from the center;*
- (c) *there are perturbations that only 0, 1, 2, ..., 11 and 12 limit cycles bifurcate from the center of system (1.5).*

In fact, in the plane $\mathbb{R}^2$ the averaging theory of first order, or the generalized Abelian integrals, or the Melnikov function provide the same information because all these methods are based on the study of the first term in ε of the Poincaré return map. Some concrete applications of that theory to planar differential systems of low degree can be seen in [6, 16]. In higher dimension, the averaging theory can be also used, for example, for the study of the Hopf bifurcation, see [12, 13].

As we will see, by using the averaging theory of first order, the limit cycles of the perturbed system, which emerge from the period annulus of the isochronous center of system (1.1), correspond to the zeros of a linear combination of the functions $f_0, f_1, \dots, f_{(n^2+7n+6)/2}$, $n = 0, 1, 2$. The proof of Theorem 1.1 for the case $n = 0$ is easy and it is done in Section 2. But the difficulty arises evidently as n increases. For $n = 1$, as the collection of functions $f_0, \dots, f_7$ is not an ET-system, part of our efforts have been focused on determining the numbers of simple zeros of Wronskian determinants $W_6(s)$ and $W_7(s)$, which have the expressions $\sum_{i=0}^k a_i(s)E^i(s)K^{k-i}(s)$ ($k = 2, 3$), where a_i is a polynomial of high degree, E and K are respectively the elliptic integrals of the first kind and second kind:

$$E(x) = \int_0^{\pi/2} \sqrt{1 - x \sin^2 \theta} d\theta, \quad K(x) = \int_0^{\pi/2} \frac{1}{\sqrt{1 - x \sin^2 \theta}} d\theta.$$

The proof is done using qualitative analysis and algebraic calculations. It turns out that all the Wronskian determinants but $W_6(s)$ do not vanish and the later has a unique zero which is simple. So the conditions of the classic Chebyshev criterion are not satisfied. According to the result of the recent paper [15], the maximum number of zeros of the linear combination of $f_0, \dots, f_7$ is less than or equal to 8. Consequently, another part of our efforts has been focused on proving that the possible upper bound 8 can be reached. To show this, we construct a function which has a zero of multiplicity 7 as well as an extra simple zero. Then, under suitable perturbation this function possesses 8 simple zeros. This is done in Section 3. For $n = 2$, the corresponding functions $f_0, f_1, \dots, f_{12}$, which contain several hypergeometric functions, is neither an Extended Complete Chebyshev system, nor a system satisfying the condition of [15]. We do not know how to find out the maximum number of zeros of all the possible linear combination of $f_0, f_1, \dots, f_{12}$. Instead, we provide a lower bound for this number or zeros. This is done in Section 4.

2 Proof of Theorem 1.1 (a)

This section is devoted to the proof of statement (a) of Theorem 1.1 by using Theorem A.1 (see Appendix A).

First, we make the polar coordinate transformation and change system (1.3) to

$$\frac{dr}{d\theta} = R_0(\theta, r) + \varepsilon R_1(\theta, r) + O(\varepsilon^2), \quad (2.1)$$

where $R_0(\theta, r) = r^3 \cos \theta \sin \theta$ and

$$\begin{aligned} R_1(\theta) = & a_{00}C + b_{00}S + r[a_{10}C^2 + (a_{01} + b_{10})CS + b_{01}S^2] + r^2[a_{20}C^3 \\ & + (a_{11} - b_{00} + b_{20})C^2S + (a_{00} + a_{02} + b_{11})CS^2 + b_{02}S^3] \\ & + r^3[a_{30}C^4 + (a_{21} - b_{10} + b_{30})C^3S + (a_{10} + a_{12} - b_{01} + b_{21})C^2S^2 \\ & + (a_{01} + a_{03} + b_{12})CS^3 + b_{03}S^4] + r^4[-b_{20}C^4S + (a_{20} - b_{11})C^3S^2 \\ & + (a_{11} - b_{02})C^2S^3 + a_{02}CS^4] + r^5[-b_{30}C^5S + (a_{30} - b_{21})C^4S^2 \\ & + (a_{21} - b_{12})C^3S^3 + (a_{12} - b_{03})C^2S^4 + a_{03}CS^5] \end{aligned} \quad (2.2)$$

with $C = \cos \theta$ and $S = \sin \theta$.

Since equation $(2.1)_{\varepsilon=0}$ has the periodic solutions $r(\theta, r_0)$ satisfying $r_0 = r(0, r_0)$ for $0 < r_0 < 1$ given in (1.2), according to the averaging theory described in Appendix A, we solve the variational differential equation

$$\frac{dM}{d\theta} = \frac{\partial}{\partial r} R_0(\theta, r(\theta, r_0))M,$$

with $M_{r_0}(0) = 1$ and get the fundamental solution

$$M_{r_0}(\theta) = (1 - r_0^2 \sin^2 \theta)^{-3/2}.$$

Next we go to study the maximum number of zeros of the function

$$\mathcal{F}(r_0) = \int_0^{2\pi} M_{r_0}^{-1}(\theta) R_1(\theta, r(\theta, r_0)) d\theta$$

with $r_0 \in (0, 1)$. Using expression (2.2), we perform the computation and obtain

$$\begin{aligned} \mathcal{F}(r_0) = & \frac{\pi}{r_0} \left((a_{10} - a_{12} + 3a_{30} + b_{01} + b_{03} - 3b_{21})r_0^2 + (b_{21} + b_{03} - b_{01})r_0^4 \right. \\ & \left. + 2(a_{12} - a_{30} - b_{03} + b_{21})(1 - \sqrt{1 - r_0^2}) - 2(a_{30} - b_{21})r_0^2 \sqrt{1 - r_0^2} \right). \end{aligned}$$

We denote

$$\begin{aligned} \alpha_0 &= \pi(a_{10} - a_{12} + 3a_{30} + b_{01} + b_{03} - 3b_{21}), \\ \alpha_1 &= \pi(b_{21} + b_{03} - b_{01}), \\ \alpha_2 &= 2\pi(a_{12} - a_{30} - b_{03} + b_{21}), \\ \alpha_3 &= -2\pi(a_{30} - b_{21}), \end{aligned}$$

and

$$f_0(s) = 1 - s^2, \quad f_1(s) = (1 - s^2)^2, \quad f_2(s) = 1 - s, \quad f_3(s) = s(1 - s^2),$$

where $s = (1 - r_0^2)^{1/2} \in (0, 1)$. Then

$$\begin{aligned} r_0 \mathcal{F}(r_0) &= \alpha_0 f_0(s) + \alpha_1 f_1(s) + \alpha_2 f_2(s) + \alpha_3 f_3(s) \\ &= (1 - s)(\alpha_0 + \alpha_1 + \alpha_2 + (\alpha_0 + \alpha_1 + \alpha_3)s + (\alpha_3 - \alpha_1)s^2 - \alpha_1 s^3). \end{aligned} \quad (2.3)$$

It is not hard to check that $\alpha_0, \alpha_1, \alpha_2$ and α_3 are independent constants and hence the four numbers $\alpha_0 + \alpha_1 + \alpha_2, \alpha_0 + \alpha_1 + \alpha_3, \alpha_3 - \alpha_1$ and α_1 can be chosen freely. Thus it follows from (2.3) that $\mathcal{F}(r_0)$ can have 0, 1, 2, 3 (and no more) simple zeros in the interval $(0, 1)$.

Using Theorem A.1, statement (a) of Theorem 1.1 is proved.

3 Proof of Theorem 1.1 (b)

In this section we will study the number of limit cycles of system (1.4) by using averaging theory of first order. We will only prove that this maximum number is 8 because according to the proof, the reader can easily see that system (1.4) can have 0, 1, 2, $\dots$, 8 limit cycles. First, let us state and prove the following lemma.

Lemma 3.1. *The maximum number of limit cycles of system (1.4) which emerge from the period annulus around the center of system (1.4)_{ε=0}, by using averaging theory of first order, is equal to the maximum number of simple zeros of the function*

$$G(s) = b_0 f_0(s) + b_1 f_1(s) + \cdots + b_7 f_7(s), \quad s \in (0, 1), \quad (3.1)$$

where $b_0, b_1, \dots, b_7$ are independent arbitrary constants and

$$\begin{cases} f_0(s) = (1-s)^2, \\ f_1(s) = (1-s)(1-s^2), \\ f_2(s) = (1-s^2)^2, \\ f_3(s) = (1-s)(1-s^2)^2, \\ f_4(s) = (1-s^2)^3, \\ f_5(s) = (1-s^2)^{5/2} g_1(s), \\ f_6(s) = (1-s^2)^{1/2} (g_1(s) - g_2(s)) - \frac{1}{2} (1-s^2) g_2(s), \\ f_7(s) = (1-s^2)^{3/2} (g_1(s) - g_2(s)), \end{cases} \quad (3.2)$$

with

$$g_1(s) = 2E(1-s^2), \quad g_2(s) = 2s^2 K(1-s^2). \quad (3.3)$$

Proof. Under the polar coordinate transformation, system (1.4) can be changed to

$$\frac{dr}{d\theta} = R_0(\theta, r) + \varepsilon R_1(\theta, r) + O(\varepsilon^2), \quad (3.4)$$

where $R_0(\theta, r) = r^5 \cos \theta \sin \theta$ and

$$\begin{aligned} R_1(\theta) = & (a_{00}C + b_{00}S) + r(a_{10}C^2 + (a_{01} + b_{10})CS + b_{01}S^2) \\ & + r^2[a_{20}C^3 + (a_{11} + b_{20})C^2S + (b_{11} + a_{02})CS^2 + b_{02}S^3] + r^3[a_{30}C^4 \\ & + (a_{21} + b_{30})C^3S + (a_{12} + b_{21})C^2S^2 + (a_{03} + b_{12})CS^3 + b_{03}S^4] \\ & + r^4[a_{40}C^5 - b_{00}C^2S + (a_{31} + b_{40})C^4S + a_{00}CS^2 + (a_{22} + b_{31})C^3S^2 \\ & + (a_{13} + b_{22})C^2S^3 + (a_{04} + b_{13})CS^4 + b_{04}S^5] + r^5[a_{50}C^6 - b_{10}C^3S \\ & + (a_{41} + b_{50})C^5S + (a_{10} - b_{01})C^2S^2 + (a_{32} + b_{41})C^4S^2 + a_{01}CS^3 \\ & + (a_{23} + b_{32})C^3S^3 + (a_{14} + b_{23})C^2S^4 + (a_{05} + b_{14})CS^5 + b_{05}S^6] \\ & + r^6[-b_{20}C^4S + (a_{20} - b_{11})C^3S^2 + (a_{11} - b_{02})C^2S^3 + a_{02}CS^4] \\ & + r^7[-b_{30}C^5S + (a_{30} - b_{21})C^4S^2 + (a_{21} - b_{12})C^3S^3 + (a_{12} - b_{03})C^2S^4 \\ & + a_{03}CS^5] + r^8[-b_{40}C^6S + (a_{40} - b_{31})C^5S^2 + (a_{31} - b_{22})C^4S^3 \\ & + (a_{22} - b_{13})C^3S^4 + (a_{13} - b_{04})C^2S^5 + a_{04}CS^6] \\ & + r^9[-b_{50}C^7S + (a_{50} - b_{41})C^6S^2 + (a_{41} - b_{32})C^5S^3 + (a_{32} - b_{23})C^4S^4 \\ & + (a_{23} - b_{14})C^3S^5 + (a_{14} - b_{05})C^2S^6 + a_{05}CS^7], \end{aligned}$$

with $C = \cos \theta$, $S = \sin \theta$.

Equation (3.4)_{ε=0} has the periodic solutions $r(\theta, r_0) = r_0(1 - 2r_0^4 \sin^2 \theta)^{-1/4}$ satisfying $r_0 = r(0, r_0)$ for $0 < r_0 < 2^{-1/4}$. We solve the variational differential equation

$$\frac{dM}{d\theta} = \frac{\partial}{\partial r} R_0(\theta, r(\theta, r_0))M,$$

with $M_{r_0}(0) = 1$ and get the fundamental solution

$$M_{r_0}(\theta) = (1 - 2r_0^4 \sin^2 \theta)^{-5/4}.$$

Next, a straightforward calculation leads to

$$\begin{aligned}\mathcal{F}(r_0) &= \int_0^{2\pi} M_{r_0}^{-1}(\theta) R_1(\theta, r(\theta, r_0)) d\theta \\ &= \int_0^{2\pi} r_0(1 - 2r_0^4 S^2)(c_{00} + c_{02}S^2 + c_{04}S^4 + c_{06}S^6 + c_{40}C^4 + c_{60}C^6 \\ &\quad + c_{22}C^2S^2 + c_{42}C^4S^2 + c_{62}C^6S^2 + c_{24}C^2S^4 + c_{44}C^4S^4 + c_{26}C^2S^6 + Y(C, S)) d\theta, \quad (3.5)\end{aligned}$$

for $r_0 \in (0, 2^{-1/4})$, where $Y(C, S) = \sum \alpha_{i,j} C^i S^j$ is a polynomial in C, S with i or j being an odd number, which leads to

$$\int_0^{2\pi} r_0(1 - 2r_0^4 S^2) Y(C, S) d\theta = 0$$

and

$$\begin{aligned}c_{00} &= a_{10} := e_0, \\ c_{02} &= -a_{10} + b_{01} := e_1, \\ c_{04} &= b_{03}r^2(\theta, r_0) := e_2r^2(\theta, r_0), \\ c_{06} &= b_{05}r^4(\theta, r_0) := e_3r^4(\theta, r_0), \\ c_{40} &= a_{30}r^2(\theta, r_0) := e_4r^2(\theta, r_0), \\ c_{60} &= a_{50}r^4(\theta, r_0) := e_5r^4(\theta, r_0), \\ c_{22} &= (a_{12} + b_{21})r^2(\theta, r_0) + (a_{10} - b_{01})r^4(\theta, r_0) := e_6r^2(\theta, r_0) - e_1r^4(\theta, r_0), \\ c_{24} &= (a_{14} + b_{23})r^4(\theta, r_0) + (a_{12} - b_{03})r^6(\theta, r_0) := e_7r^4(\theta, r_0) + e_8r^6(\theta, r_0), \\ c_{26} &= (a_{14} - b_{05})r^8(\theta, r_0) := e_9r^8(\theta, r_0), \\ c_{42} &= (a_{32} + b_{41})r^4(\theta, r_0) + (a_{30} - b_{21})r^6(\theta, r_0) := e_{10}r^4(\theta, r_0) + e_{11}r^6(\theta, r_0), \\ c_{44} &= (a_{32} - b_{23})r^8(\theta, r_0) := e_{12}r^8(\theta, r_0), \\ c_{62} &= (a_{50} - b_{41})r^8(\theta, r_0) := e_{13}r^8(\theta, r_0).\end{aligned}$$

It is not hard to check that the constants $e_0, e_1, \dots, e_{13}$ are independent. Computing (3.5), we get

$$\mathcal{F}(r_0) = I_1(r_0) + I_2(r_0) + I_3(r_0) + I_4(r_0),$$

where

$$\begin{aligned}I_1(r_0) &= \alpha_1 r_0 + \alpha_2 r_0^5, \\ I_2(r_0) &= \frac{1}{15r_0^5}((\alpha_3 + \alpha_4 r_0^4 + \alpha_5 r_0^8)\bar{g}_1(r_0) - (\alpha_3 + (\alpha_3 + \alpha_4)r_0^4)\bar{g}_2(r_0)), \\ I_3(r_0) &= -\frac{\pi}{16r_0^7}((2\alpha_6 + 2\alpha_7 r_0^4 + \alpha_8 r_0^8)\bar{g}_3(r_0) + (2\alpha_6 r_0^4 + (\alpha_6 + 2\alpha_7)r_0^8 + (\alpha_6 + \alpha_7 + \alpha_8)r_0^{12})), \\ I_4(r_0) &= -\frac{1}{30r_0^5}((4\alpha_9 + \alpha_{10}r_0^4 - (7\alpha_9 + \alpha_{10})r_0^8)\bar{g}_1(r_0) - (4\alpha_9 + (4\alpha_9 + \alpha_{10})r_0^4)\bar{g}_2(r_0)),\end{aligned}$$

with

$$\begin{aligned}\alpha_1 &= \pi(2e_0 + e_1), & \alpha_2 &= \frac{\pi}{8}(-16e_0 - 14e_1 + 5e_3 + 5e_5 + e_7 + e_{10}), \\ \alpha_3 &= -e_2 - e_4 + e_6, & \alpha_4 &= -3e_2 + 7e_4 - 2e_6, \\ \alpha_5 &= 16e_2 + 6e_4 + 4e_6, & \alpha_6 &= e_9 - e_{12} + e_{13}, \\ \alpha_7 &= 2e_{12} - 4e_{13}, & \alpha_8 &= e_{13}, \\ \alpha_9 &= -e_8 + e_{11}, & \alpha_{10} &= 3e_8 - 13e_{11},\end{aligned}$$

and

$$\begin{aligned}\bar{g}_1(r_0) &= E(2r_0^4) + \sqrt{1-2r_0^4} E(1-1/(1-2r_0^4)), \\ \bar{g}_2(r_0) &= (1-2r_0^4)K(2r_0^4) + \sqrt{1-2r_0^4} K(1-1/(1-2r_0^4)), \\ \bar{g}_3(r_0) &= \sqrt{1-2r_0^4} - 1.\end{aligned}$$

Using the expression of each α_i , one can easily check that $\alpha_1, \alpha_2, \dots, \alpha_{10}$ are independent constants.

To simplify the computation, we let $s = (1-2r_0^4)^{1/2}$, $s \in (0, 1)$. Using the definition of the elliptic functions, we have

$$\begin{aligned}sE(1-1/s^2) + E(1-s^2) &= 2E(1-s^2), \\ sK(1-1/s^2) + s^2K(1-s^2) &= 2s^2K(1-s^2).\end{aligned}$$

Hence we obtain

$$240r_0^7\mathcal{F}(r_0) = G(s) = b_0f_0(s) + b_1f_1(s) + \dots + b_7f_7(s), \quad s \in (0, 1),$$

where $f_i(s)$, $i = 0, 1, \dots, 7$, are the functions defined in (3.2), and the constants $b_0, b_1, \dots, b_7$ in (3.1) are independent constants each of which is a linear combination of $\alpha_1, \alpha_2, \dots, \alpha_{10}$.

By Theorem A.1, the lemma is proved. $\square$

Next, we denote by $W_i(s)$ the Wronskian determinant for the functions $f_0, f_1, \dots, f_i$ depending on s :

$$W_i(s) = W(f_0, \dots, f_i)(s), \quad i = 0, 1, \dots, 7.$$

In what follows we will show that all the Wronskian determinants have no zeros except W_6 which vanishes at a unique zero, which is simple.

By direct calculation we obtain

$$\left\{ \begin{array}{l} W_0(s) = (1-s)^2, \\ W_1(s) = (1-s)^4, \\ W_2(s) = 2(1-s)^6, \\ W_3(s) = -12(1-s)^8, \\ W_4(s) = 288(1-s)^{10}, \\ W_5(s) = Y_5(s)(Z_{50}(s)g_2(s) + Z_{51}(s)g_1(s)), \\ W_6(s) = Y_6(s)(Z_{60}(s)g_2^2(s) + Z_{61}(s)g_2(s)g_1(s) + Z_{62}(s)g_1^2(s)), \\ W_7(s) = Y_7(s)(Z_{70}(s)g_2^3(s) + Z_{71}(s)g_2^2(s)g_1(s) + Z_{72}(s)g_2(s)g_1^2(s) + Z_{73}(s)g_1^3(s)), \end{array} \right. \quad (3.6)$$

where

$$\begin{aligned}Y_5(s) &= 288s^{-3}(1-s)^{15/2}(1+s)^{-5/2}, \\ Z_{50}(s) &= 1-20s-33s^2-20s^3+s^4, \\ Z_{51}(s) &= -2(1-5s+10s^2+5s^3+10s^4-5s^5+s^6), \\ Y_6(s) &= 432s^{-6}(1-s)^2(1+s)^{-8}, \\ Z_{60}(s) &= \sqrt{1-s^2}(-330-761s+3720s^2+25036s^3+63490s^4+100713s^5 \\ &\quad +102410s^6+66145s^7+23800s^8+3760s^9-770s^{10}-269s^{11}) \\ &\quad -2(210+637s-2490s^2-22210s^3-67910s^4-129477s^5-160950s^6 \\ &\quad -135498s^7-74890s^8-25715s^9-2910s^{10}+728s^{11}+240s^{12}), \\ Z_{61}(s) &= \sqrt{1-s^2}(660+1544s-2430s^2-6587s^3+18350s^4+65033s^5+107880s^6 \\ &\quad +106430s^7+64350s^8+11893s^9-4790s^{10}-2177s^{11}+1580s^{12}+536s^{13}) \\ &\quad +4(210+644s-150s^2-2133s^3+1670s^4+9718s^5+23630s^6+29805s^7 \\ &\quad +27970s^8+14822s^9+3430s^{10}-2877s^{11}-300s^{12}+736s^{13}+240s^{14}),\end{aligned}$$

$$\begin{aligned}
Z_{62}(s) &= s\sqrt{1-s^2}(-44-1410s+405s^2+6880s^3-6789s^4-63230s^5-111946s^6 \\
&\quad -111330s^7-57735s^8-6880s^9+7555s^{10}-270s^{11}-1146s^{12}-80s^{13}+4s^{14}) \\
&\quad -2s(28+1590s+1873s^2-8750s^3-4018s^4+81350s^5+209051s^6+273750s^7 \\
&\quad +208594s^8+80650s^9-4307s^{10}-8050s^{11}+1712s^{12}+1560s^{13}+32s^{14}), \\
Y_7(s) &= -3888s^{-9}(1-s)^{-7/2}(1+s)^{-27/2}, \\
Z_{70}(s) &= \sqrt{1-s^2}(12600+126840s+317850s^2+52804s^3-990100s^4-1480411s^5 \\
&\quad +1111220s^6+6912788s^7+11963620s^8+12059608s^9+7333520s^{10} \\
&\quad +1499322s^{11}-1510040s^{12}-1134068s^{13}+73250s^{14}+339916s^{15}+124820s^{16} \\
&\quad +15361s^{17}+60s^{18})-8\sqrt{1-s^2}(1-s^2)^3s^9(1+s)^{10}(3150+31710s+95130s^2 \\
&\quad +77561s^3-209720s^4-618425s^5-351890s^6+1120612s^7+3107560s^8 \\
&\quad +4115235s^9+3350990s^{10}+1407368s^{11}-195310s^{12}-611035s^{13}-252580s^{14} \\
&\quad +65059s^{15}+103320s^{16}+36240s^{17}+3600s^{18}), \\
Z_{71}(s) &= 3\sqrt{1-s^2}(-8400-85120s-232660s^2-173008s^3+437450s^4+811463s^5 \\
&\quad -2067190s^6-10578234s^7-22263600s^8-30680904s^9-30782260s^{10} \\
&\quad -22474150s^{11}-10890340s^{12}-2297236s^{13}+798740s^{14}+530744s^{15} \\
&\quad -174630s^{16}-249833s^{17}-82630s^{18}-10282s^{19}-80s^{20})-24(2100+21280s \\
&\quad +69900s^2+98437s^3-32560s^4-271689s^5+257650s^6+2845680s^7 \\
&\quad +7627330s^8+12561491s^9+14784900s^{10}+12847024s^{11}+8032670s^{12} \\
&\quad +3146115s^{13}+365000s^{14}-280461s^{15}-58190s^{16}+94628s^{17}+76050s^{18} \\
&\quad +24320s^{19}+2400s^{20}), \\
Z_{72}(s) &= 3s\sqrt{1-s^2}(1120+43540s+298396s^2+428670s^3-487118s^4-1615600s^5 \\
&\quad +1598463s^6+14431800s^7+34113716s^8+48793180s^9+48574944s^{10} \\
&\quad +34190000s^{11}+15087426s^{12}+1926380s^{13}-1797616s^{14}-499770s^{15} \\
&\quad +470882s^{16}+285160s^{17}+48663s^{18}-1040s^{19}+84s^{20}+80s^{21}) \\
&\quad +24s(280+13530s+98799s^2+200770s^3-31141s^4-611730s^5+35248s^6 \\
&\quad +4717630s^7+13875311s^8+23495150s^9+27804240s^{10}+24063400s^{11} \\
&\quad +14574499s^{12}+5157020s^{13}+149657s^{14}-658770s^{15}-66644s^{16} \\
&\quad +208280s^{17}+108756s^{18}+14970s^{19}+320s^{20}), \\
Z_{73}(s) &= s^2\sqrt{1-s^2}(480-10712s-177820s^2-762730s^3-87610s^4+2806841s^5 \\
&\quad +909430s^6-14409046s^7-40794440s^8-62886616s^9-63914340s^{10} \\
&\quad -42197910s^{11}-14285100s^{12}+1389476s^{13}+2447740s^{14}-233714s^{15} \\
&\quad -691850s^{16}-196931s^{17}+7070s^{18}+1990s^{19}-760s^{20}-8s^{21}) \\
&\quad -8s^2(270+538s+49880s^2+295952s^3+320270s^4-754291s^5 \\
&\quad -1276480s^6+4132165s^7+18820380s^8+36641394s^9+45116550s^{10} \\
&\quad +37467621s^{11}+19916970s^{12}+4809620s^{13}-1224770s^{14}-907079s^{15} \\
&\quad +291730s^{16}+324433s^{17}+58270s^{18}+1472s^{19}+180s^{20}).
\end{aligned}$$

Lemma 3.2. Let g_1 and g_2 be the two functions defined in (3.3) and let $h(s) = g_1(s)/g_2(s)$. Then $h(s) > 0$, $h'(s) < 0$, $s \in (0, 1)$ and

$$\lim_{s \rightarrow 0^+} h(s) = +\infty, \quad \lim_{s \rightarrow 1^-} h(s) = 1, \quad \lim_{s \rightarrow 0^+} h'(s) = -\infty, \quad \lim_{s \rightarrow 1^-} h'(s) = -1. \quad (3.7)$$

Proof. It follows directly from the definition of the elliptic integral that $g_i(s) > 0$ ($i = 1, 2$), $s \in (0, 1)$ and hence $h(s) > 0$, $s \in (0, 1)$. A direct computation shows that

$$\begin{aligned} g_1(s) &= 1 - \frac{1}{2}s^2 \log s + \frac{1}{4}s^2(4 \log 2 - 1) + o(s^2), \\ g_2(s) &= -s^2 \log s + 2s^2 \log 2 + o(s^2), \end{aligned}$$

where $s \rightarrow 0^+$. Thus the first and the third equalities of (3.7) hold.

Similarly, we find that

$$\begin{cases} g_1(s) = \frac{\pi}{2} \left(1 - \frac{1}{2}(1-s) + \frac{1}{16}(1-s)^2 + O((1-s)^3) \right), \\ g_2(s) = \frac{\pi}{2} \left(1 - \frac{3}{2}(1-s) + \frac{5}{16}(1-s)^2 + O((1-s)^3) \right), \end{cases} \quad (3.8)$$

as $s \rightarrow 1^-$. This verifies the second and the fourth equalities of (3.7).

Next we go to prove that $h'(s) < 0$, $s \in (0, 1)$. By straightforward calculation we find

$$dh/ds = (1 - 2h + h^2 s^2)/(s - s^3).$$

Hence $h(s)$ is a solution of system

$$\dot{h} = s^2 h^2 - 2h + 1, \quad \dot{s} = s - s^3. \quad (3.9)$$

System (3.9) has two invariant straight lines $s = 0$ and $s = 1$ as well as two singularities at $S_1(0, 1/2)$ and $S_2(1, 1)$, where S_1 is a saddle and S_2 is a saddle-node of system (3.9). Moreover, system (3.9) has two horizontal isocline curves

$$h_+(s) = \frac{1}{1 - \sqrt{1-s^2}} \quad \text{and} \quad h_-(s) = \frac{1}{1 + \sqrt{1-s^2}}, \quad (3.10)$$

satisfying

$$h'_+(s) < 0, \quad h'_-(s) > 0, \quad h_+(0) = +\infty, \quad h_-(0) = \frac{1}{2}, \quad h_+(1) = h_-(1) = 1.$$

Obviously,

$$h'(s) = s^2(h(s) - h_+(s)) \frac{h(s) - h_-(s)}{s - s^3}. \quad (3.11)$$

In view of (3.8) and (3.10), it follows that

$$h_-(s) < h(s) < h_+(s), \quad s \rightarrow 1^-. \quad (3.12)$$

We assert that

$$h_-(s) < h(s) < h_+(s), \quad s \in (0, 1). \quad (3.13)$$

Indeed, if there exists some point $s_0 \in (0, 1)$ such that $h(s_0) \geq h_+(s_0)$, then by (3.11) we find $h'(s_0) \geq 0$. By the monotonicity of $h_+(s)$ we know that $h(s) > h_+(s)$ for all $s_0 < s < 1$. This contradicts (3.12). Hence $h(s) < h_+(s)$ for $s \in (0, 1)$. If there exists some point $s_0 \in (0, 1)$ such that $h(s_0) = h_-(s_0)$, then by (3.11) we know that $h'(s_0) = 0$. Since $h'_-(s_0) > 0$, it follows $h(s) < h_-(s)$ for $s \rightarrow s_0^+$. Using this fact, we find that the curve $h = h(s)$ cannot go across the curve $h = h_-(s)$ at any point $s_1 > s_0$ because once $h(s_1) = h_-(s_1)$, it must hold that $h(s) < h_-(s)$ for $s \rightarrow s_1^+$. This also contradicts (3.12). Hence $h(s) > h_-(s)$, $s \in (0, 1)$.

Finally, combining (3.13) and (3.11), we conclude that $dh/ds < 0$, $s \in (0, 1)$. $\square$

Lemma 3.3. *The function $W_5(s)$ does not vanish in the open interval $(0, 1)$.*

Proof. Using Sturm's theorem (see [18]) and $Z_{51}(0) = -2$, we find that $Z_{51}(s) < 0$ for all $s \in (0, 1)$. Hence we have

$$W_5(s) = Y_5(s)Z_{51}(s)g_2(s) \left(\frac{Z_{50}(s)}{Z_{51}(s)} + \frac{g_1(s)}{g_2(s)} \right), \quad s \in (0, 1). \quad (3.14)$$

A direct computation leads to $Z_{50}(s_0) = 0$, where

$$s_0 = 5 + 3\sqrt{15}/2 - \sqrt{231 + 60\sqrt{15}}/2 \approx 0.0463551.$$

Again, by Sturm's theorem we find that $Z_{50}(s) > 0$ for $s \in (0, s_0)$ and $Z_{50}(s) < 0$ for $s \in (s_0, 1)$. Further,

$$\frac{d}{ds} \left(\frac{Z_{50}(s)}{Z_{51}(s)} \right) = \frac{2p_9(s)}{Z_{51}^2(s)},$$

where

$$p_9(s) = 15 + 86s - 290s^2 - 364s^3 - 575s^4 - 274s^5 + 190s^6 + 68s^7 - 65s^8 + 2s^9.$$

Using Sturm's theorem, we get that $p_9(s) > 0$, $s \in (0, 1/5)$. This fact, combined with $Z_{50}(0)/Z_{51}(0) = -1/2$, yields that

$$\frac{Z_{50}(s)}{Z_{51}(s)} > 0 \text{ for } s \in (s_0, 1) \quad \text{and} \quad \frac{Z_{50}(s)}{Z_{51}(s)} \in \left(-\frac{1}{2}, 0\right) \text{ for } s \in (0, s_0). \quad (3.15)$$

Since by Lemma 3.2 we have $g_1(s)/g_2(s) > 1$, it follows from (3.14) and (3.15) that $W_5(s) \neq 0$ for all $s \in (0, 1)$. $\square$

Next, we will determine the sign of the functions $W_6(s)$ and $W_7(s)$. In order to make the computation easier we need to make the transformation of variable $r = ((1-s)/(1+s))^{1/2}$ or equivalently, $s = (1-r^2)/(1+r^2)$. We also need the following lemma. Let

$$\bar{h}(r) = \frac{g_1(s)}{g_2(s)} \Big|_{s=(1-r^2)/(1+r^2)}, \quad r \in (0, 1). \quad (3.16)$$

Lemma 3.4. *The function $h = \bar{h}(r)$ is the solution of the differential system*

$$\dot{h} = ((r-1)^2 h - r^2 - 1)((r+1)^2 h - r^2 - 1), \quad \dot{r} = r(r^4 - 1), \quad (3.17)$$

satisfying $\bar{h}'(r) > 0$ for $r \in (0, 1)$, $\bar{h}(0) = 1$, and $\lim_{r \rightarrow 1^-} \bar{h}(r) = +\infty$.

Proof. The conclusion follows from the proof of Lemma 3.2 by direct calculation. $\square$

Lemma 3.5. *The function $W_6(s)$ has a unique zero in $(0, 1)$ and this zero is simple.*

Proof. Let $s = (1-r^2)/(1+r^2)$ for $0 < r < 1$. Then it follows from the definition of $W_6(s)$ that

$$\overline{W}_6(r) := W_6\left(\frac{1-r^2}{1+r^2}\right) = \overline{Y}_6(r)\bar{g}_2^2(r)(C_{60}(r) + C_{61}(r)\bar{h}(r) + C_{62}(r)\bar{h}^2(r)),$$

where $\bar{h}(r)$ is the function defined in (3.16) and

$$\begin{aligned} \overline{Y}_6(r) &= \frac{2Y_6(s)}{(1+r^2)^{16}} \Big|_{s=(1-r^2)/(1+r^2)}, \\ \bar{g}_2(r) &= g_2(s) \Big|_{s=(1-r^2)/(1+r^2)}, \\ C_{60}(r) &= -(1+r^2)^4(-620235 - 386944r + 63082r^2 + 4352r^3 + 1114260r^4 \\ &\quad + 747808r^5 - 136770r^6 - 21280r^7 + 349425r^8 + 201312r^9 - 5852r^{10} \\ &\quad + 97568r^{11} + 90000r^{12} + 59232r^{13} + 29692r^{14} - 39392r^{15} - 39225r^{16} \\ &\quad - 6880r^{17} - 10350r^{18} + 18912r^{19} + 24540r^{20} - 1152r^{21} - 1242r^{22} \\ &\quad + 2304r^{23} + 2835r^{24}), \\ C_{61}(r) &= 2(1+r^2)^2(107415 + 181136r - 33728r^2 + 11072r^3 - 399945r^4 - 731304r^5 \\ &\quad + 149536r^6 - 47384r^7 + 1468275r^8 + 1717984r^9 - 340256r^{10} + 125584r^{11} \\ &\quad + 2454435r^{12} + 1436304r^{13} - 95680r^{14} + 188016r^{15} + 156165r^{16} \\ &\quad - 122704r^{17} + 100096r^{18} - 61184r^{19} - 145275r^{20} - 30376r^{21} - 23136r^{22} \\ &\quad + 33384r^{23} + 39345r^{24} - 1632r^{25} - 2592r^{26} + 4464r^{27} + 5985r^{28}), \end{aligned}$$

$$\begin{aligned}
C_{62}(r) = & (1 - r^2)(-835065 - 346016r - 836827r^2 - 363808r^3 + 3987923r^4 \\
& + 1598928r^5 + 4009041r^6 + 1699296r^7 - 8764869r^8 - 3404464r^9 \\
& - 8875247r^{10} - 3660192r^{11} + 1910023r^{12} + 319232r^{13} + 2054749r^{14} \\
& + 561472r^{15} + 3501349r^{16} + 1423232r^{17} + 3542143r^{18} + 1373728r^{19} \\
& - 275927r^{20} + 31056r^{21} - 405789r^{22} - 8224r^{23} + 15321r^{24} + 43728r^{25} \\
& + 44043r^{26} + 4512r^{27} + 2493r^{28} + 6624r^{29} + 9135r^{30}).
\end{aligned}$$

Define

$$\bar{w}_6(r, h) = C_{60}(r) + C_{61}(r)h + C_{62}(r)h^2.$$

We will show that on the curve $C := \{(r, h) \mid \bar{w}_6(r, h) = 0, r \in (0, 1)\}$, there is a unique point P at which vector field (3.17) is tangent to C . We call P the *contact point* with the vector field (3.17). In fact, by direct computation we obtain

$$D(r, h) := (\partial \bar{w}_6 / \partial r, \partial \bar{w}_6 / \partial h) \cdot (\dot{r}, \dot{h}) = -2 \sum_{i=0}^3 d_i(r)h^i,$$

where

$$\begin{aligned}
d_0(r) = & (1 + r^2)^4(-107415 + 12336r + 2451586r^2 + 1530176r^3 - 4561843r^4 \\
& - 2896872r^5 - 3880856r^6 - 2845432r^7 + 4366665r^8 + 2321984r^9 - 1985574r^{10} \\
& - 1627056r^{11} - 175627r^{12} - 500432r^{13} - 524832r^{14} + 797392r^{15} + 938867r^{16} \\
& + 960688r^{17} + 476566r^{18} - 543968r^{19} - 529425r^{20} - 119176r^{21} - 195912r^{22} \\
& + 200040r^{23} + 275163r^{24} - 24288r^{25} - 27378r^{26} + 31248r^{27} + 39375r^{28}), \\
d_1(r) = & -(1 + r^2)^2(-1049895 - 889424r - 726170r^2 - 1522240r^3 + 7709074r^4 \\
& + 7536696r^5 + 1889478r^6 + 6432616r^7 - 31052618r^8 - 28741640r^9 - 6577218r^{10} \\
& - 16058648r^{11} - 6172022r^{12} + 328016r^{13} - 22042946r^{14} - 12811872r^{15} \\
& + 33763160r^{16} + 23233568r^{17} - 4741326r^{18} + 5413264r^{19} + 1388182r^{20} \\
& - 2485928r^{21} + 3589522r^{22} - 1915736r^{23} - 3612822r^{24} - 869704r^{25} - 768726r^{26} \\
& + 678984r^{27} + 822606r^{28} - 77664r^{29} - 113814r^{30} + 122832r^{31} + 170415r^{32}), \\
d_2(r) = & -(1 - r^4)(1777545 + 1046176r + 3311818r^2 + 1457408r^3 - 16459052r^4 \\
& - 8289520r^5 - 15874018r^6 - 7006192r^7 + 62475516r^8 + 28826544r^9 + 35342330r^{10} \\
& + 15708784r^{11} - 49795012r^{12} - 19466016r^{13} - 8698050r^{14} - 3515776r^{15} \\
& - 24983266r^{16} - 12853376r^{17} - 14258354r^{18} - 5372832r^{19} + 31346828r^{20} \\
& + 11380816r^{21} + 2668682r^{22} + 500624r^{23} - 4087764r^{24} - 691664r^{25} - 471570r^{26} \\
& + 404016r^{27} + 455268r^{28} - 51264r^{29} - 120294r^{30} + 84960r^{31} + 121905r^{32}), \\
d_3(r) = & -(1 - r^2)^3(-835065 - 346016r - 836827r^2 - 363808r^3 + 3987923r^4 \\
& + 1598928r^5 + 4009041r^6 + 1699296r^7 - 8764869r^8 - 3404464r^9 - 8875247r^{10} \\
& - 3660192r^{11} + 1910023r^{12} + 319232r^{13} + 2054749r^{14} + 561472r^{15} + 3501349r^{16} \\
& + 1423232r^{17} + 3542143r^{18} + 1373728r^{19} - 275927r^{20} + 31056r^{21} - 405789r^{22} \\
& - 8224r^{23} + 15321r^{24} + 43728r^{25} + 44043r^{26} + 4512r^{27} + 2493r^{28} + 6624r^{29} + 9135r^{30}).
\end{aligned}$$

Using Sturm's theorem, we find $d_3(r) > 0, r \in (0, 1)$. Further, the resultant of $\bar{w}_6(r, h)$ and $D(r, h)$ with respect to h is a polynomial in the variable r of degree 166, which, applying again Sturm's theorem, can be proved to have a unique simple zero $r_0 \in (0, 1)$ with $9/10 < r_0 < 91/100$. Hence there exists a unique h_0 such that

$$\bar{w}_6(r_0, h_0) = D(r_0, h_0) = 0.$$

This confirms that on the curve C there is a unique point (r_0, h_0) at which the vector field (3.17) is tangent to C .

By direct computation we have

$$C_{61}^2 - 4C_{60}C_{62} = 3600(1+r^2)^4 p_{56}(r),$$

where $p_{56}(r)$ is a polynomial of degree 56. Again, we can apply Sturm's theorem to prove that $p_{56}(r) > 0$ and $C_{62}(r) < 0$ in $r \in (0, 1)$. Let

$$C_- = \left\{ h = \bar{h}_-(r) = \frac{-C_{61} - \sqrt{C_{61}^2 - 4C_{60}C_{62}}}{2C_{62}} \right\}$$

and

$$C_+ = \left\{ h = \bar{h}_+(r) = \frac{-C_{61} + \sqrt{C_{61}^2 - 4C_{60}C_{62}}}{2C_{62}} \right\}$$

be the two branches of the curve C in the (r, h) -plane. A calculation shows that

$$\bar{h}_-(0) = 1, \quad \bar{h}_+(0) = -\frac{179}{241}, \quad \bar{h}(0) = 1, \quad \bar{h}'_-(0) = \frac{64}{231}, \quad \bar{h}'_+(0) = \frac{96160}{1916673}, \quad \bar{h}'(0) = 0,$$

$\bar{h}_+(1) = 1/2$ and when $r \rightarrow 1^-$,

$$\bar{h}_-(r) = \frac{15}{1-r} + \cdots, \quad \bar{h}(r) = \frac{1}{(\log 4 - \log(1-r))(r-1)^2} + \cdots,$$

where the dots denote the terms which are infinitesimal being compared to the former one.

It follows that

$$\bar{h}_+(r) < \bar{h}(r) < \bar{h}_-(r) \quad \text{as } r \rightarrow 0^+,$$

$$\bar{h}_+(r) < \bar{h}_-(r) < \bar{h}(r) \quad \text{as } r \rightarrow 1^-.$$

Obviously, the curve $\Gamma = \{h = \bar{h}(r)\}$ intersects C_- in at least one point (r^*, h^*) . By an observation on the direction of vector field (3.17) at the two endpoints of the segment of curve $\{(r, h) \mid h = \bar{h}_-(r), r \in (0, r^*]\}$, we find that there exists a point P at which the vector field (3.17) is tangent to the curve C_- (see Figure 2). Since the contact point P is unique, the curve Γ cannot intersect C_- in other points. Moreover, the curve Γ has no common point with C_+ , otherwise a second contact point would emerge. Therefore the function $\bar{w}_6(r, \bar{h}(r))$ has a unique zero in the interval $(0, 1)$. This yield that $\bar{W}_6(r)$ has a unique zero in the interval $(0, 1)$.

Finally, since $r_0 < 91/100$ and $\bar{h}(91/100) \approx 30.54045135 < h_-(91/100) \approx 35.81140037$, it follows that $r_0 < r^*$. This means that (r^*, h^*) is not the contact point of C_- with the vector field. Therefore, the unique zero of $\bar{W}_6(r)$ is simple and thus the required conclusion holds. $\square$

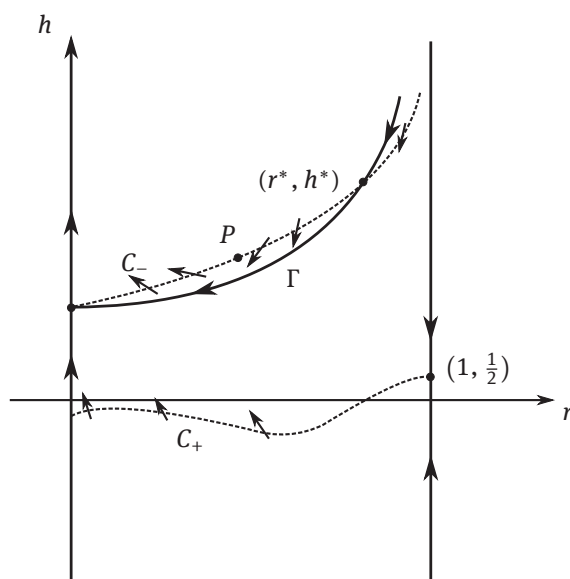


Figure 2. The curve Γ has a unique common point with C_- .

Lemma 3.6. *The function $W_7(s)$ does not vanish in the open interval $(0, 1)$.*

Proof. By taking transformation $s = (1 - r^2)/(1 + r^2)$, $0 < r < 1$, we obtain from

$$W_7(s) = Y_7(s)g_2^3(s)\left(Z_{70}(s) + Z_{71}(s)\frac{g_1(s)}{g_2(s)} + Z_{72}(s)\frac{g_1^2(s)}{g_2^2(s)} + Z_{73}(s)\frac{g_1^3(s)}{g_2^3(s)}\right)$$

that

$$\overline{W}_7(r) := W_7(s)|_{s=(1-r^2)/(1+r^2)} = \overline{Y}_7(r)\overline{g}_2^3(r)w_7(r),$$

with

$$w_7(r) = C_{70}(r) + C_{71}(r)\bar{h}(r) + C_{72}(r)\bar{h}^2(r) + C_{73}(r)\bar{h}^3(r),$$

where $\bar{h}(r)$ is the function defined in (3.16) and

$$\begin{aligned}\overline{Y}_7(r) &= \frac{16rY_7(s)}{(1+r^2)^{45}}\Big|_{s=(1-r^2)/(1+r^2)}, \\ \overline{g}_2(r) &= g_2(s)\Big|_{s=(1-r^2)/(1+r^2)}, \\ C_{70}(r) &= 4r(2301810 + 59701740r^2 + 727558755r^4 - 364073500959r^6 \\ &\quad + 3428595727383r^8 - 10544549722741r^{10} + 3730074158113r^{12} \\ &\quad + 49576965802069r^{14} - 40961684822285r^{16} - 396767446632771r^{18} \\ &\quad + 1609108209115716r^{20} - 3203874112868486r^{22} + 4488939441796380r^{24} \\ &\quad - 4090928490421940r^{26} + 3532364976473268r^{28} - 553179108109916r^{30} \\ &\quad + 1013150621056664r^{32} + 1960313990634764r^{34} + 589939846153370r^{36} \\ &\quad + 1620210743086006r^{38} + 485382645874034r^{40} + 531165876327274r^{42} \\ &\quad + 41364680501774r^{44} - 94823180435898r^{46} - 177772367419966r^{48} \\ &\quad - 169241753304026r^{50} - 122994867094516r^{52} - 75347197109816r^{54} \\ &\quad - 38883154071124r^{56} - 17219939255972r^{58} - 6420781546876r^{60} \\ &\quad - 1958359101836r^{62} - 448781050130r^{64} - 53625628864r^{66} + 12582124635r^{68} \\ &\quad + 10801790529r^{70} + 4279675055r^{72} + 1249653091r^{74} + 292463001r^{76} \\ &\quad + 55639117r^{78} + 8432979r^{80} + 963045r^{82} + 73560r^{84} + 2790r^{86}), \\ C_{71}(r) &= -3(1+r^2)^{23}(62214075 + 32618040r + 57649725r^2 + 31801560r^3 \\ &\quad - 311083290r^4 - 161979460r^5 - 284045278r^6 - 156667936r^7 \\ &\quad + 1118662427r^8 + 566796672r^9 + 1015532421r^{10} + 545264968r^{11} \\ &\quad + 1483781736r^{12} + 882038312r^{13} + 1421172920r^{14} + 856116840r^{15} \\ &\quad + 32362070r^{16} - 26403192r^{17} + 175436346r^{18} + 24968536r^{19} \\ &\quad - 2404284r^{20} - 55354400r^{21} + 31425964r^{22} - 122819896r^{23} \\ &\quad - 129246866r^{24} - 82418248r^{25} - 147649550r^{26} - 83415400r^{27} \\ &\quad - 50931800r^{28} - 40785832r^{29} - 66284296r^{30} - 5366728r^{31} \\ &\quad - 2020081r^{32} - 3163552r^{33} - 1758167r^{34} + 617456r^{35} + 624838r^{36} \\ &\quad + 127460r^{37} + 489090r^{38} + 19640r^{39} + 50775r^{40} + 14760r^{41} + 40425r^{42}), \\ C_{72}(r) &= 3(-1+r^2)(1+r^2)^{21}(-113149575 - 48972840r - 219419550r^2 \\ &\quad - 96894000r^3 + 524903715r^4 + 226248140r^5 + 1216381628r^6 \\ &\quad + 542101276r^7 - 1447999469r^8 - 628726896r^9 - 3888634006r^{10} \\ &\quad - 1779992536r^{11} - 2876628303r^{12} - 1259091264r^{13} - 2141852592r^{14} \\ &\quad - 754743200r^{15} + 1112715114r^{16} + 583398240r^{17} + 4296755236r^{18} \\ &\quad + 1852262832r^{19} + 2357760878r^{20} + 907216248r^{21} + 687064168r^{22})\end{aligned}$$

$$\begin{aligned}
& + 234254008r^{23} + 208154838r^{24} + 151069392r^{25} - 302393084r^{26} \\
& - 153692000r^{27} - 210572686r^{28} - 204720480r^{29} - 215766192r^{30} \\
& - 104223744r^{31} - 131334803r^{32} - 41765176r^{33} - 20357126r^{34} \\
& - 11754496r^{35} - 11547809r^{36} + 729436r^{37} + 904508r^{38} + 33740r^{39} \\
& + 444615r^{40} + 42400r^{41} + 114450r^{42} + 18360r^{43} + 50925r^{44}), \\
C_{73}(r) = & -(-1 + r^2)^2(1 + r^2)^{20}(164085075 + 58271640r + 318975300r^2 \\
& + 115256400r^3 - 920074365r^4 - 320089980r^5 - 2079306888r^6 \\
& - 746616756r^7 + 2728744989r^8 + 941220096r^9 + 7172532884r^{10} \\
& + 2616432392r^{11} + 779348149r^{12} + 546885328r^{13} - 4995512160r^{14} \\
& - 1385204368r^{15} - 4685978914r^{16} - 1569573552r^{17} - 4525683128r^{18} \\
& - 2222726784r^{19} + 166924638r^{20} - 650651016r^{21} + 4249280720r^{22} \\
& + 1445886696r^{23} + 2224407802r^{24} + 1065394144r^{25} + 814165128r^{26} \\
& + 244092272r^{27} + 320088074r^{28} - 95701872r^{29} - 346329440r^{30} \\
& - 166958928r^{31} - 180571809r^{32} - 99999912r^{33} - 42394764r^{34} \\
& - 27896336r^{35} - 30911889r^{36} + 462756r^{37} + 365688r^{38} - 212820r^{39} \\
& + 107865r^{40} + 50400r^{41} + 137700r^{42} + 21960r^{43} + 61425r^{44}).
\end{aligned}$$

The number of zeros of $\bar{W}_7(r)$ in $(0, 1)$ equals the number of intersection points of the curve

$$C = \{C_{70}(r) + C_{71}(r)h + C_{72}(r)h^2 + C_{73}(r)h^3 = 0\}$$

with the curve $\Gamma = \{h = \bar{h}(r)\}$ in the (r, h) -plane. In what follows we will study the relative positions of C and Γ . To this end, since Γ is not an algebraic curve, we need to establish another auxiliary algebraic curve which is easier for computation.

First, using Sturm's theorem, we find that $C_{73}(r) \neq 0$, $r \in (0, 1)$. This means that

$$\bar{w}_7(r, h) := C_{70}(r) + C_{71}(r)h + C_{72}(r)h^2 + C_{73}(r)h^3$$

is a cubic polynomial of h for each fixed $r \in (0, 1)$. Let

$$A = C_{72}^2 - 3C_{71}C_{73}, \quad B = C_{71}C_{72} - 9C_{70}C_{73}, \quad C = C_{71}^2 - 3C_{70}C_{72},$$

and $\Delta = B^2 - 4AC$. It is not hard to see that Δ has exactly two zeros r_1, r_2 in $(0, 1)$ with $39/50 < r_1 < 79/100$, $91/100 < r_2 < 23/25$. If $r \in (0, r_1) \cup (r_2, 1)$ then $\Delta > 0$; if $r \in (r_1, r_2)$ then $\Delta < 0$. Therefore, the curve C has three branches C_1 (the lower branch), C_2 (the middle branch) and C_3 (the upper branch) with the property that C_2 and C_3 have the same endpoints $E_1(r_1, h_1)$ and $E_2(r_2, h_2)$. See Figure 3.

Second, we claim that $C_2 \cup C_3$ lies over the curve Γ . To show this we introduce an auxiliary algebraic curve $Y = \{h = \Phi(r)\}$ with

$$\begin{aligned}
\Phi(r) = & \frac{1}{2}(5 + 4r^2 + 6r^4 + 8r^6 + 10r^8 + 12r^{10} + 13r^{12} + 15r^{14} + 16r^{16} \\
& + 18r^{18} + 20r^{20} + 21r^{22} + 22r^{24} + 38r^{26} + 25r^{28} + 26r^{30} + 28r^{32} + 30r^{34} \\
& + 30r^{36} + 32r^{38} + 34r^{40} + 34r^{42} + 36r^{44} + 38r^{46} + 38r^{48} + 40r^{50}),
\end{aligned}$$

where $r \in (0, 1)$. By direct computations as well as by applying Sturm's theorem we obtain

$$\bar{w}_7(r, \Phi(r)) = p_{238}(r) < 0,$$

where $p_{238}(r)$ is a polynomial of degree 238. Thus the curve C does not intersect Y . Moreover, since the straight line $r = 9/10$ intersects the curve C and Y at the points $(9/10, c_1^*)$, $(9/10, c_2^*)$, $(9/10, c_3^*) \in C$ and $(9/10, \phi_1^*) \in Y$, respectively, where

$$c_1^* \approx 0.1592878, \quad c_2^* \approx 30.7373179, \quad c_3^* \approx 40.3056908, \quad \phi_1^* \approx 26.9337561,$$

we conclude that $C_2 \cup C_3$ lies over the curve Y , and C_1 lies below the curve Y . See Figure 3.

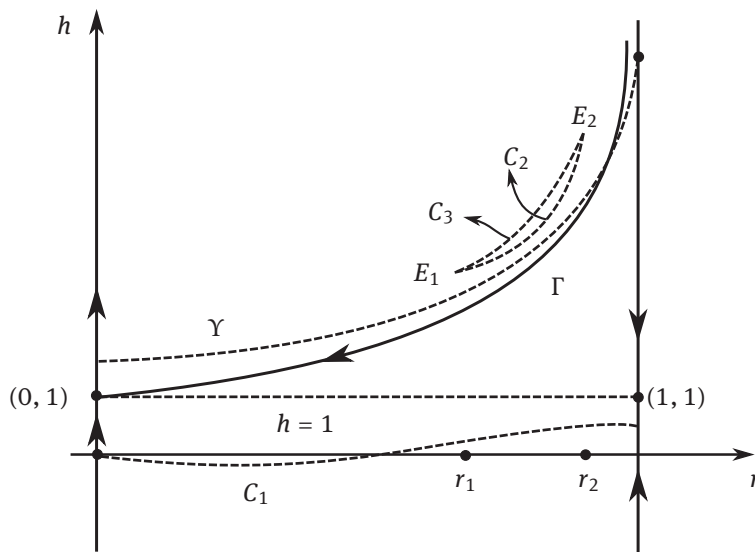


Figure 3. The relative positions of the curve Γ and C .

On the other hand, using (3.17), we obtain by computation that

$$\begin{aligned}
 (\dot{h} - \Phi'(r)\dot{r})|_{h=\Phi(r)} = & \frac{1}{4}(-9 + 42r^2 + 3r^4 - 4r^8 - 8r^{10} + 18r^{12} - 2r^{14} + 14r^{16} - 6r^{18} - 36r^{20} + 10r^{22} \\
 & + 43r^{24} - 826r^{26} + 280r^{28} + 748r^{30} - 23r^{32} - 82r^{34} + 70r^{36} + 62r^{38} \\
 & - 132r^{40} + 102r^{42} + 9r^{44} - 124r^{46} + 103r^{48} + 42r^{50} + 4248r^{52} + 4446r^{54} \\
 & - 81r^{56} + 82r^{58} + 115r^{60} + 184r^{62} - 20r^{64} + 124r^{66} + 108r^{68} + 72r^{70} \\
 & + 128r^{72} + 116r^{74} + 1308r^{78} - 1056r^{80} + 64r^{82} + 136r^{84} + 152r^{86} - 12r^{88} \\
 & + 144r^{90} + 156r^{92} - 8r^{94} + 152r^{96} + 160r^{98} - 4r^{100} + 160r^{102} - 1600r^{104}),
 \end{aligned}$$

which has a unique zero in the interval $r_0 \in (0, 1)$ with $0 < r_0 < 1/2$. Therefore, there exists a unique contact point on the curve Υ with the vector field (3.17). Taking this into account and noting the fact that

$$\Phi(0) - \bar{h}(0) > 0, \quad \Phi(1/2) - \bar{h}(1/2) > 1 > 0, \quad \Phi(23/25) - \bar{h}(23/25) > 2/5 > 0,$$

it is clear that the curve $\Upsilon|_{r \in (0, 23/25)}$ lies over $\Gamma|_{r \in (0, 23/25)}$, otherwise there would exist at least two contact points on $\Upsilon|_{r \in (0, 23/25)}$ with the vector field (3.17), which leads to a contradiction.

In summary, according to the relative positions of Υ and Γ as well as the relative positions of Υ and $C_2 \cup C_3$, we find that the claim is true.

Third, we claim that C_1 lies below the curve Γ . This claim is easy to confirm due to the following facts:

- Γ lies over the straight line $h = 1$ (by Lemma 3.4);
- C_1 does not intersect the line $h(r) = 1$ because $C_{70}(r) + C_{71}(r) + C_{72}(r) + C_{73}(r) \neq 0$, $r \in (0, 1)$ (by applying Sturm's theorem);
- C_1 is a continuous curve passing through the point $(0, 0)$ (because $C_{70}(0) + C_{71}(0)h + C_{72}(0)h^2 + C_{73}(0)h^3 = 0$ implies $h = 0$).

Finally, taking into account the above results, we conclude that the curve C has no common points with the curve Γ . Thus $\bar{W}_7(r) \neq 0$, i.e., $W_7(s) \neq 0$. $\square$

Proof of Theorem 1.1 (b). It follows from equation (3.6), Lemma 3.3, Lemma 3.5 and Lemma 3.6 that $w_i(s) \neq 0$, $i = 0, 1, 2, 3, 4, 5, 7$ and $w_6(s)$ has a simple zero in $(0, 1)$. Very recently, Novaes and Torregrosa [15] proved that if the analytical functions $f_0, f_1, \dots, f_n : I \rightarrow \mathbb{R}$ satisfy

- all Wronskian determinants $W_i(s) = W(f_0, \dots, f_i)(s)$ but W_{n-1} have no zero in the interval I , and
- W_{n-1} has a unique simple zero in I ,

then any linear combination of $f_0, f_1, \dots, f_n$ can possess at most $n + 1$ zeros in I , counting with multiplicities. But Novaes and Torregrosa do not prove that the upper bound can be reached in the general cases.

In what follows we will show that the upper bound 8 can be reached in our system.

Let $s_0 = 200/10001$, $E_0 = E(1 - s_0^2)$, $K_0 = K(1 - s_0^2)$ and

$$k = 4224932006353520086857838671137556(AE_0^2 + BE_0K_0 + CK_0^2),$$

where

$$A = 162632756824646343526934358039550191813769181219039360950399,$$

$$B = -26652192151547736499563618692313149392412558977363562257000,$$

$$C = 5303055781903243156556160927943006794119698485540000000.$$

Consider the function

$$G(s) = a_0f_0(s) + a_1f_1(s) + \dots + a_6f_6(s) + kf_7(s), \quad s \in (0, 1). \quad (3.18)$$

By direct calculation we get the power series of G around the point s_0 :

$$G(s) = e_0 + e_1(s - s_0) + \dots + e_6(s - s_0)^6 + e_7(s - s_0)^7 + \dots,$$

where e_i is the linear combination of $a_0, a_1, \dots, a_6$. We solve the equations

$$e_0 = 0, \quad e_1 = 0, \quad \dots, \quad e_6 = 0,$$

and find the values of $a_0, a_1, \dots, a_6$ which have the form

$$a_i = \sum_{j=0}^3 q_{ij} E_0^j K_0^{3-j}, \quad i = 0, 1, \dots, 6,$$

where each q_{ij} is an integer which occupies many digits. We will not write down here the explicit expression of a_i for the sake of brevity. By the way, we would like to point out that our purpose of choosing such a k in (3.18) is to make the expression of a_i to be relative simple. It turns out that

$$G(s) = e_7(s - s_0)^7 + O((s - s_0)^8), \quad s \rightarrow s_0, \quad (3.19)$$

where

$$e_7 = -\frac{1000600150020001500060001}{264828047171937480000}(A_0E_0^3 - A_1E_0^2K_0 + A_2E_0K_0^2 - A_3K_0^3),$$

with

$$A_0 = 9434365215900096702757086133640555232723441933485420521056595192874$$

$$214200423786458280868229,$$

$$A_1 = 1732705885903853693884640417850773768026603157714652167368522666300$$

$$981344222008726672767200,$$

$$A_2 = 1550165493805896510978302557036974594023078666779498920543561841897$$

$$54798606490809063800000,$$

$$A_3 = 3092940994614887237716575160316676355701108658523634572794702093798$$

$$4340166396000000000,$$

and

$$e_7 \approx -8.7875569 \times 10^{97}.$$

On the other hand, at the endpoint $s = 1$ we have

$$G(s) = B(1 - s) + o((1 - s)), \quad (3.20)$$

where

$$\begin{aligned} B &= 12800263824853240746592933248429440000532\pi \\ &\quad \cdot (5792157212337693345948517518844704378216313299609167E_0^2 \\ &\quad - 1020168724968577415929102393676106281813032946881000E_0K_0 \\ &\quad + 203589533638142582450403592018967753470820000000K_0^2) \\ &\approx 1.6008470 \times 10^{91}. \end{aligned}$$

Equations (3.19) and (3.20) mean that (i) G has a zero at s_0 with multiplicity 7, (ii) there exists an ε_0 with $s_0 < \varepsilon_0 < 1$ such that $G(s)$ is negative in $(s_0, \varepsilon_0]$, and (iii) $G(s)$ is positive near the endpoint $s = 1$.

Fixing the numbers $a_0, a_1, \dots, a_6$ and k , we consider the function

$$G_\varepsilon(s) = G(s) + \sum_{i=0}^7 \varepsilon_i f_i(s), \quad s \in (0, 1). \quad (3.21)$$

We note that f_i can be extended analytically to $(0, 1]$. Thus there exists an $M > 0$ such that

$$G_\varepsilon(\varepsilon_0) < \frac{1}{2}G(\varepsilon_0) < 0, \quad G_\varepsilon(s) > \frac{1}{2}B(1-s), \quad \text{when } s \rightarrow 1^-,$$

for all $|\varepsilon_i| < M, i = 0, 1, \dots, 7$. Moreover, near s_0 we find

$$\sum_{i=0}^7 \varepsilon_i f_i(s) = \mu_0 + \mu_1(s - s_0) + \dots + \mu_7(s - s_0)^7 + \dots,$$

where $\mu_i = \mu_i(\varepsilon_0, \varepsilon_1, \dots, \varepsilon_7)$ is linear combination of $\varepsilon_0, \varepsilon_1, \dots, \varepsilon_7$. One can directly check that the matrix of the coefficients of $\mu_0, \mu_1, \dots, \mu_7$ with respect to $\varepsilon_0, \varepsilon_1, \dots, \varepsilon_7$ has rank 8, and hence $\mu_0, \mu_1, \dots, \mu_7$ are independent.

Consequently, since f_i is analytic at s_0 and G has a zero at s_0 with multiplicity 7, it follows that there exists some small $|\varepsilon_i| \ll M$ ($i = 0, 1, \dots, 7$) (and hence μ_i is small) such that G_ε has exactly 7 simple zeros in a neighborhood of s_0 . In view of (3.21) G has an extra zero in $(\varepsilon_0, 1)$. According to the result of [15], this zero is simple. That is to say, G_ε has 8 simple zeros.

Finally, using Lemma 3.1 and averaging theory of first order, we see that systems (1.4) have at most 8 limit cycles, and the upper bound can be reached. The proof is finished. $\square$

4 Proof of Theorem 1.1 (c)

The goal of this section is to investigate the number of limit cycles of system (1.5) which bifurcate from the period annulus of the isochronous center. Before we prove our result, we should first recall the concept of hypergeometric function.

Let $H(a, b, c, z)$ be the ordinary hypergeometric function which is defined for $|z| < 1$ by the power series

$$H(a, b, c, z) = \sum_{k=0}^{\infty} \frac{(a)_k (b)_k}{(c)_k} \frac{z^k}{k!}, \quad (4.1)$$

where

$$(a)_k = \begin{cases} 1, & k = 0, \\ a(a+1) \cdots (a+k-1), & k > 0. \end{cases}$$

It is undefined (or infinite) if c equals a non-positive integer. Many of the common mathematical functions can be expressed in terms of the hypergeometric function. For example, $(1-z)^{-a} = H(a, 1, 1, z)$ and

$$H\left(\frac{1}{2}, \frac{1}{2}, 1, m\right) = \frac{2K(m)}{\pi}, \quad H\left(-\frac{1}{2}, \frac{1}{2}, 1, m\right) = \frac{2E(m)}{\pi}.$$

For more information on hypergeometric functions, the reader is referred to [1, Chapter 15].

Lemma 4.1. *The maximum number of limit cycles of system (1.5) which emerge from the period annulus of center of system (1.5) _{$\varepsilon=9$} , by using averaging theory of first order, is equal to the maximum number of the simple zeros of the function*

$$G(s) = c_0 g_0(s) + c_1 g_1(s) + \cdots + c_{12} g_{12}(s),$$

where $c_0, c_1, \dots, c_{12}$ are independent arbitrary constants, and

$$\begin{aligned} g_0(s) &= s, \\ g_1(s) &= s^2, \\ g_2(s) &= s\sqrt{1-s}, \\ g_3(s) &= sH(1, 1/2, 2, s), \\ g_4(s) &= sH(1, 5/2, 4, s), \\ g_5(s) &= sH(1, 3/2, 3, s), \\ g_6(s) &= s^{4/3}H(-2/3, 1/2, 1, s), \\ g_7(s) &= s^{5/3}H(-1/3, 1/2, 1, s), \\ g_8(s) &= s^{5/3}H(-1/3, 3/2, 2, s), \\ g_9(s) &= s^{5/3}H(-1/3, 5/2, 3, s), \\ g_{10}(s) &= s^{4/3}H(-2/3, 3/2, 2, s), \\ g_{11}(s) &= s^{4/3}\left(-(1-s)H\left(\frac{1}{3}, \frac{5}{2}, 3, s\right) - 2H\left(-\frac{2}{3}, \frac{5}{2}, 3, s\right)\right), \\ g_{12}(s) &= s^{2/3}\left((4-9s)H\left(\frac{2}{3}, \frac{3}{2}, 3, s\right) - (1-s)(4+33s)H\left(\frac{2}{3}, \frac{5}{2}, 3, s\right)\right). \end{aligned}$$

Proof. As usual we take the polar coordinate transformation to change system (1.5) to

$$\frac{dr}{d\theta} = R_0(\theta, r) + \varepsilon R_1(\theta, r) + O(\varepsilon^2), \quad (4.2)$$

where $R_0(\theta, r) = r^7 \cos \theta \sin \theta$ and $R_1(\theta, r) = R_{11}(\theta, r) + R_{12}(\theta, r)$ with

$$\begin{aligned} R_{11}(\theta, r) &= r(a_{10}C^2 + b_{01}S^2) + r^3[a_{30}C^4 + (a_{12} + b_{21})C^2S^2 + b_{03}S^4] \\ &\quad + r^5[a_{50}C^6 + (a_{32} + b_{41})C^4S^2 + (a_{14} + b_{23})C^2S^4 + b_{05}S^6] \\ &\quad + r^7[a_{70}C^8 + (a_{10} - b_{01})C^2S^2 + (a_{52} + b_{61})C^6S^2 + (a_{34} + b_{43})C^4S^4 \\ &\quad + (a_{16} + b_{25})C^2S^6 + b_{07}S^8] + r^9[(a_{30} - b_{21})C^4S^2 + (a_{12} - b_{03})C^2S^4] \\ &\quad + r^{11}[(a_{50} - b_{41})C^6S^2 + (a_{32} - b_{23})C^4S^4 + (a_{14} - b_{05})C^2S^6] \\ &\quad + r^{13}[(a_{70} - b_{61})C^8S^2 + (a_{52} - b_{43})C^6S^4 + (a_{34} - b_{25})C^4S^6 + (a_{16} - b_{07})C^2S^8], \end{aligned}$$

and $R_{12}(\theta, r)$ is a polynomial of degree 13 in r of the form $\sum_{i,j,k} d_{i,j,k} C^i S^j r^k$ where i or j are odd numbers. As before, here $C = \cos \theta$, $S = \sin \theta$. We do not write down the explicit expression of $R_{12}(\theta, r)$ because it is too long and, as we will see, it does not play any role in further calculation.

The equation (4.2) _{$\varepsilon=0$} has the periodic solutions $r(\theta, r_0) = r_0(1 - 3r_0^6 \sin^2 \theta)^{-1/6}$ satisfying $r_0 = r(0, r_0)$ for $0 < r_0 < 3^{-1/6}$. The corresponding variational differential equation

$$\frac{dM}{d\theta} = \frac{\partial}{\partial r} R_0(\theta, r(\theta, r_0))M,$$

with $M_{r_0}(0) = 1$ has the fundamental solution

$$M_{r_0}(\theta) = (1 - 3r_0^6 \sin^2 \theta)^{-7/6}.$$

Next we go to study the maximum number of zeros of the function

$$\mathcal{F}(r_0) = \int_0^{2\pi} M_{r_0}^{-1}(\theta) R_1(\theta, r(\theta, r_0)) d\theta = r_0^7 \int_0^{2\pi} r^{-7}(\theta, r_0) R_1(\theta, r(\theta, r_0)) d\theta,$$

when $r_0 \in (0, 3^{-1/6})$.

One can check directly that

$$\int_0^{2\pi} r^{-7}(\theta, r_0) R_{12}(\theta, r(\theta, r_0)) d\theta = 0.$$

It turns out that

$$\mathcal{F}(r_0) = r_0^7 \int_0^{2\pi} r^{-7}(\theta, r_0) R_{11}(\theta, r(\theta, r_0)) d\theta, \quad r_0 \in (0, 3^{-1/6}).$$

Further, taking the transformation $r_0 = (s/3)^{1/6}$, we have

$$\bar{\mathcal{F}}(s) := \mathcal{F}(r_0) = \int_0^{2\pi} (1 - s \sin^2 \theta)^{7/6} R_{11}(\theta, \bar{r}(\theta, s)) d\theta, \quad s \in (0, 1),$$

where $\bar{r}(\theta, s) = 3^{-1/6}(s/(1 - s \sin^2 \theta))^{1/6}$.

Using an algebraic manipulator, we obtain after a long calculation that

$$\bar{\mathcal{F}}(s) = \frac{\bar{c}_0 \bar{f}_0(s) + \bar{c}_1 \bar{f}_1(s) + \cdots + \bar{c}_{15} \bar{f}_{15}(s)}{s^{17/6}(1-s)^{2/3}}, \quad (4.3)$$

where

$$\begin{aligned} \bar{f}_0(s) &= s^3(1-s)^{2/3}, \\ \bar{f}_1(s) &= s^4(1-s)^{2/3}, \\ \bar{f}_2(s) &= s^3(1-s)^{7/6}, \\ \bar{f}_3(s) &= s^2((1-s)^{7/6} - (1-s)^{2/3}), \\ \bar{f}_4(s) &= 8(1-s)^{7/6} + (1-s)^{2/3}(s^2 + 4s - 8), \\ \bar{f}_5(s) &= 2s(1-s)^{7/6} - (1-s)^{2/3}(2s - s^2), \\ \bar{f}_6(s) &= s^{10/3}(1-s)^{2/3}H(-\frac{2}{3}, \frac{1}{2}, 1, s), \\ \bar{f}_7(s) &= s^{11/3}(1-s)^{2/3}H(-\frac{1}{3}, \frac{1}{2}, 1, s), \\ \bar{f}_8(s) &= s^{11/3}(1-s)^{2/3}H(-\frac{1}{3}, \frac{3}{2}, 2, s), \\ \bar{f}_9(s) &= s^{11/3}(1-s)^{2/3}H(-\frac{1}{3}, \frac{5}{2}, 3, s), \\ \bar{f}_{10}(s) &= s^{14/3}\left((1-s)^{2/3}H(\frac{2}{3}, \frac{3}{2}, 3, s) + H(\frac{2}{3}, \frac{3}{2}, 3, \frac{s}{s-1})\right), \\ \bar{f}_{11}(s) &= s^{13/3}\left((1-s)^{2/3}H(\frac{1}{3}, \frac{3}{2}, 3, s) + (1-s)^{1/3}H(\frac{1}{3}, \frac{3}{2}, 3, \frac{s}{s-1})\right), \\ \bar{f}_{12}(s) &= s^{14/3}\left(2H(\frac{2}{3}, \frac{3}{2}, 3, \frac{s}{s-1}) + (1-s)^{2/3}H(\frac{2}{3}, \frac{5}{2}, 4, s) - H(\frac{2}{3}, \frac{5}{2}, 4, \frac{s}{s-1})\right), \\ \bar{f}_{13}(s) &= s^{10/3}(1-s)^{2/3}\left(9\Gamma(-\frac{2}{3})\Gamma(\frac{7}{6})H(-\frac{2}{3}, \frac{3}{2}, -\frac{1}{6}, 1-s) \right. \\ &\quad \left. - 10\sqrt{3}\pi^{3/2}H(-\frac{2}{3}, \frac{3}{2}, 2, s) + 18(1-s)^{7/6}\Gamma(-\frac{7}{6})\Gamma(\frac{8}{3})H(\frac{1}{2}, \frac{8}{3}, \frac{13}{6}, 1-s)\right), \\ \bar{f}_{14}(s) &= s^{10/3}(1-s)^{1/3}\left(243(1-s)^{1/3}\Gamma(-\frac{2}{3})\Gamma(\frac{7}{6})H(-\frac{2}{3}, \frac{5}{2}, -\frac{1}{6}, 1-s) \right. \\ &\quad \left. - 360\sqrt{3}\pi^{3/2}(1-s)^{1/3}H(-\frac{2}{3}, \frac{5}{2}, 3, s) + 40\sqrt{3}\pi^{3/2}sH(\frac{1}{3}, \frac{3}{2}, 3, \frac{s}{s-1}) \right. \\ &\quad \left. + 20\sqrt{3}\pi^{3/2}s\left((1-s)^{1/3}H(\frac{1}{3}, \frac{5}{2}, 4, s) - H(\frac{1}{3}, \frac{5}{2}, 4, \frac{s}{s-1})\right) \right. \\ &\quad \left. + 324(1-s)^{3/2}\Gamma(-\frac{7}{6})\Gamma(\frac{11}{3})H(\frac{1}{2}, \frac{11}{3}, \frac{13}{6}, 1-s)\right), \\ \bar{f}_{15}(s) &= s^{11/3}\left(-240(1-s)^{2/3}H(-\frac{1}{3}, \frac{7}{2}, 4, s) + s\left(16H(\frac{2}{3}, \frac{3}{2}, 3, \frac{s}{s-1}) \right. \right. \\ &\quad \left. \left. - 16H(\frac{2}{3}, \frac{5}{2}, 4, \frac{s}{s-1}) + 5(1-s)^{2/3}H(\frac{2}{3}, \frac{7}{2}, 5, s) + 5H(\frac{2}{3}, \frac{7}{2}, 5, \frac{s}{s-1})\right)\right). \end{aligned}$$

Here $\Gamma(z)$ is the Gamma function defined by $\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$ and $\bar{c}_i$, for $i = 0, 1, 2, \dots, 15$, is the linear combination of a_{ij} and b_{ij} . We do not give the explicit expressions of $\bar{c}_i$ ($i = 0, 1, 2, \dots, 15$) because they are too long. We can check by direct calculation that $\bar{c}_1, \bar{c}_2, \dots, \bar{c}_{15}$ are independent.

From (4.1) we have

$$\begin{aligned}\bar{f}_3(s) &= s^3(1-s)^{2/3} \frac{\sqrt{1-s}-1}{s} = -\frac{1}{2}s^3(1-s)^{2/3}H(1, \frac{1}{2}, 2, s), \\ \bar{f}_4(s) &= -\frac{1}{2}s^3(1-s)^{2/3} \frac{8\sqrt{1-s}+s^2+4s-8}{-s^3/2} - \frac{1}{2}s^3(1-s)^{2/3}H(1, \frac{5}{2}, 4, s), \\ \bar{f}_5(s) &= -\frac{1}{4}s^3(1-s)^{2/3} \frac{s-2+2\sqrt{1-s}}{-s^2/4} = -\frac{1}{4}s^3(1-s)^{2/3}H(1, \frac{3}{2}, 3, s).\end{aligned}$$

Using Pfaff transformation (see [1, Chapter 15])

$$(1-z)^a H(a, b, c, z) = H(a, c-b, c, z/(z-1)),$$

as well as Gauss' contiguous relation

$$\begin{aligned}\frac{abz}{c}H(a+1, b+1, c+1, z) &= a(H(a+1, b, c, z) - H(a, b, c, z)) \\ &= b(H(a+1, b, c, z) - H(a, b, c, z)) \\ &= (c-1)(H(a, b, c-1, z) - H(a, b, c, z)) \\ &= \frac{(c-a)H(a-1, b, c, z) + (a-c+bz)H(a, b, c, z)}{1-z} \\ &= \frac{(c-b)H(a, b-1, c, z) + (b-c+az)H(a, b, c, z)}{1-z} \\ &= z \frac{(c-a)(c-b)H(a, b, c+1, z) + c(a+b-c)H(a, b, c, z)}{c(1-z)},\end{aligned}\quad (4.4)$$

we obtain that

$$\begin{aligned}\bar{f}_{10}(s) &= 2s^{14/3}(1-s)^{2/3}H(\frac{2}{3}, \frac{3}{2}, 3, s), \\ \bar{f}_{11}(s) &= 2s^{13/3}(1-s)^{2/3}H(\frac{1}{3}, \frac{3}{2}, 3, s), \\ \bar{f}_{12}(s) &= 2s^{14/3}(1-s)^{2/3}H(\frac{2}{3}, \frac{5}{2}, 4, s), \\ \bar{f}_{14}(s) &= 3s^{10/3}(1-s)^{2/3} \left(81\Gamma(-\frac{2}{3})\Gamma(\frac{7}{6})H(-\frac{2}{3}, \frac{5}{2}, -\frac{1}{6}, 1-s) \right. \\ &\quad \left. + 40\sqrt{3}\pi^{3/2} \left(-H(-\frac{2}{3}, \frac{5}{2}, 3, s) - 2(1-s)H(\frac{1}{3}, \frac{5}{2}, 3, s) \right) \right. \\ &\quad \left. + 108(1-s)^{7/6}\Gamma(-\frac{7}{6})\Gamma(\frac{11}{3})H(\frac{1}{2}, \frac{11}{3}, \frac{13}{6}, 1-s) \right), \\ \bar{f}_{15}(s) &= \frac{216}{35}s^{8/3}(1-s)^{2/3} \left((4-9s)H(\frac{2}{3}, \frac{3}{2}, 3, s) - (1-s)(4+33s)H(\frac{2}{3}, \frac{5}{2}, 3, s) \right).\end{aligned}$$

Further, applying the formula

$$\begin{aligned}H(a, b, c, z) &= \frac{\Gamma(c)\Gamma(c-a-b)}{\Gamma(c-a)\Gamma(c-b)}H(a, b, a+b-c+1, 1-z) \\ &\quad + (1-z)^{c-a-b} \frac{\Gamma(c)\Gamma(a+b-c)}{\Gamma(a)\Gamma(b)}H(c-a, c-b, c-a-b+1, 1-z),\end{aligned}$$

for $|\arg(1-z)| < \pi$, we obtain that

$$\begin{aligned}\bar{f}_{13}(s) &= s^{10/3}(1-s)^{2/3} \left(9\Gamma(-\frac{2}{3})\Gamma(\frac{8}{3})\Gamma(\frac{1}{2}) - 10\sqrt{3}\pi^{3/2} \right) H(-\frac{2}{3}, \frac{3}{2}, 2, s) \\ &= \sqrt{\pi}s^{10/3}(1-s)^{2/3} \left(9\Gamma(-\frac{2}{3})\Gamma(\frac{8}{3}) - 10\sqrt{3}\pi \right) H(-\frac{2}{3}, \frac{3}{2}, 2, s) \\ &= -20\sqrt{3}\pi^{3/2}s^{10/3}(1-s)^{2/3}H(-\frac{2}{3}, \frac{3}{2}, 2, s),\end{aligned}$$

$$\begin{aligned}
\bar{f}_{14}(s) &= \frac{3}{2}s^{10/3}(1-s)^{2/3}\left(\left(81\Gamma(-\tfrac{2}{3})\Gamma(\tfrac{11}{3})\Gamma(\tfrac{1}{2}) - 80\sqrt{3}\pi^{3/2}\right)H(-\tfrac{2}{3}, \tfrac{5}{2}, 3, s) \right. \\
&\quad \left. - 160\sqrt{3}\pi^{3/2}(1-s)H(\tfrac{1}{3}, \tfrac{5}{2}, 3, s)\right) \\
&= \frac{3\sqrt{\pi}}{2}s^{10/3}(1-s)^{2/3}\left(\left(81\Gamma(-\tfrac{2}{3})\Gamma(\tfrac{11}{3}) - 80\sqrt{3}\pi\right)H(-\tfrac{2}{3}, \tfrac{5}{2}, 3, s) \right. \\
&\quad \left. + 160\sqrt{3}\pi(s-1)H(\tfrac{1}{3}, \tfrac{5}{2}, 3, s)\right) \\
&= -240\sqrt{3}\pi^{3/2}s^{10/3}(1-s)^{2/3}\left((1-s)H(\tfrac{1}{3}, \tfrac{5}{2}, 3, s) + 2H(-\tfrac{2}{3}, \tfrac{5}{2}, 3, s)\right).
\end{aligned}$$

By the above equalities, we obtain from (4.3) that

$$\bar{\mathcal{F}}(s) = s^{-5/6}(\tilde{c}_0 f_0(s) + \tilde{c}_1 f_1(s) + \cdots + \tilde{c}_{15} f_{15}(s)), \quad (4.5)$$

where the functions f_i modulo a nonzero constant are $\bar{f}_i/(s^2(1-s)^{2/3})$:

$$\begin{aligned}
f_0(s) &= s, \\
f_1(s) &= s^2, \\
f_2(s) &= s\sqrt{1-s}, \\
f_3(s) &= sH(1, \tfrac{1}{2}, 2, s), \\
f_4(s) &= sH(1, \tfrac{5}{2}, 4, s), \\
f_5(s) &= sH(1, \tfrac{3}{2}, 3, s), \\
f_6(s) &= s^{4/3}H(-\tfrac{2}{3}, \tfrac{1}{2}, 1, s), \\
f_7(s) &= s^{5/3}H(-\tfrac{1}{3}, \tfrac{1}{2}, 1, s), \\
f_8(s) &= s^{5/3}H(-\tfrac{1}{3}, \tfrac{3}{2}, 2, s), \\
f_9(s) &= s^{5/3}H(-\tfrac{1}{3}, \tfrac{5}{2}, 3, s), \\
f_{10}(s) &= s^{8/3}H(\tfrac{2}{3}, \tfrac{3}{2}, 3, s), \\
f_{11}(s) &= s^{7/3}H(\tfrac{1}{3}, \tfrac{3}{2}, 3, s), \\
f_{12}(s) &= s^{8/3}H(\tfrac{2}{3}, \tfrac{5}{2}, 4, s), \\
f_{13}(s) &= s^{4/3}H(-\tfrac{2}{3}, \tfrac{3}{2}, 2, s), \\
f_{14}(s) &= -s^{4/3}\left((1-s)H(\tfrac{1}{3}, \tfrac{5}{2}, 3, s) + 2H(-\tfrac{2}{3}, \tfrac{5}{2}, 3, s)\right) \\
f_{15}(s) &= s^{2/3}\left((4-9s)H(\tfrac{2}{3}, \tfrac{3}{2}, 3, s) - (1-s)(4+33s)H(\tfrac{2}{3}, \tfrac{5}{2}, 3, s)\right),
\end{aligned}$$

It is also not hard to see that in (4.5), $\tilde{c}_0, \tilde{c}_1, \dots, \tilde{c}_{15}$ are independent constants.

Further, applying repeatedly Gauss' contiguous relation (4.4), we find that

$$f_{10} = 12f_7 - 12f_8, \quad f_{11} = 6f_6 - 6f_{13}, \quad f_{12} = 36f_8 - 36f_9. \quad (4.6)$$

It follows from (4.5) and (4.6) that

$$\mathcal{F}(r_0) = \bar{\mathcal{F}}(s) = s^{-5/6} \sum_{i=0}^{12} c_i g_i(s) = s^{-5/6} G(s),$$

where

$$g_0 = f_0, \quad g_1 = f_1, \quad \dots, \quad g_9 = f_9, \quad g_{10} = f_{13}, \quad g_{11} = f_{14}, \quad g_{12} = f_{15},$$

and $c_0, c_1, \dots, c_{12}$ are independent constants.

Finally, by Theorem A.1, we obtain the required conclusion. $\square$

Proof of Theorem 1.1 (c). First, we claim that the functions $g_0, g_1, \dots, g_{12}$ in Lemma 4.1 are linear independent. To show this, we write

$$\bar{G}(u) = d_0 g_0(u^3) + d_1 g_1(u^3) + \cdots + d_{12} g_{12}(u^3), \quad u > 0.$$

Then, near the point $u = 0$ we have

$$\begin{aligned} \frac{8957952\overline{G}(u)}{u^3} &= 8957952(d_0 + d_2 + d_3 + d_4 + d_5) + 8957952(d_{10} - 3d_{11} + d_6)u \\ &\quad - (348364800d_{12} - 8957952(d_7 + d_8 + d_9))u^2 \\ &\quad + 1119744(8d_1 - 4d_2 + 2d_3 + 5d_4 + 4d_5)u^3 \\ &\quad - 1492992(3d_{10} - 11d_{11} + 2d_6))u^4 + 82944(1400d_{12} \\ &\quad - 3(6d_7 + 9d_8 + 10d_9))u^5 - 559872(2d_2 - 2d_3 - 7d_4 - 5d_5)u^6 \\ &\quad - 124416(5d_{10} - 20d_{11} + 3d_6))u^7 \\ &\quad + 6912(5390d_{12} - 3(18d_7 + 30d_8 + 35d_9))u^8 \\ &\quad - 139968(4d_2 - 5d_3 - 7(3d_4 + 2d_5))u^9 - 6912(35d_{10} - 147d_{11} \\ &\quad + 20d_6)u^{10} + 2880(6860d_{12} - 3(20d_7 + 35d_8 + 42d_9))u^{11} \\ &\quad - 69984(5d_2 - 7d_3 - 33d_4 - 21d_5)u^{12} - 14112(9d_{10} - 39d_{11} \\ &\quad + 5d_6)u^{13} + 6720(1870d_{12} - 3(5d_7 + 9d_8 + 11d_9))u^{14} \\ &\quad - 4374(56d_2 - 3(28d_3 + 143d_4 + 88d_5))u^{15} - 1008(77d_{10} \\ &\quad - 341d_{11} + 42d_6)u^{16} + 176(50050d_{12} - 378d_7 - 693d_8 \\ &\quad - 858d_9)u^{17} - 2187(84d_2 - 11(12d_3 + 65d_4 + 39d_5))u^{18} + O(u^{19}) \\ &:= \alpha_0 + \alpha_1 u + \cdots + \alpha_{18} u^{18} + O(u^{19}). \end{aligned}$$

If $\overline{G}(u) \equiv 0$, then we have $\alpha_0 = \alpha_1 = \cdots = \alpha_{18} = 0$. By direct calculation, we get $d_0 = d_1 = \cdots = d_{12} = 0$. (By the way, it is remarkable that we cannot get from $\alpha_0 = \alpha_1 = \cdots = \alpha_{12} = 0$ that $d_0 = \cdots = d_{12} = 0$.) This shows that our claim holds.

Since g_i is an analytic function on $(0, 1)$ for $i = 0, 1, \dots, 12$, we know that, by applying [5, Lemma 4.5], by suitable choice of $c_0, c_1, \dots, c_{12}$, $c_0 g_0(s) + c_1 g_1(s) + \cdots + c_{12} g_{12}(s)$ can have $0, 1, 2, \dots, 12$ simple zeros in $(0, 1)$.

Consequently, according to Theorem A.1, there exist some coefficients a_{ij}, b_{ij} ($i + j = 0, 1, \dots, 7$) such that system (1.5) has $0, 1, 2, \dots, 12$ limit cycles. This completes the proof. $\square$

Remark 4.2. It seems very hard to find the smallest upper bound of the number of limit cycles of system (1.5) which emerge from the period annulus of the isochronous center for the case $n = 2$. In fact, the expressions of the Wronskian determinants W_k for $k \geq 9$ are too complicated to determine the number of simple zeros of them. On the other hand, by using the same method as the one in the proof of case $n = 1$, we can find some coefficients $c_0, c_1, \dots, c_{12}$ such that $c_0 g_0(s) + c_1 g_1(s) + \cdots + c_{12} g_{12}(s)$ has 12 simple zeros in a small interval $(s_0 - \varepsilon, s_0 + \varepsilon)$ and has extra zero in $(0, s_0 - \varepsilon)$ with $s_0 = 7/10$. However, we cannot prove that the extra zero is simple.

A The Averaging Theory of First Order

In this appendix we present the basic results from the averaging theory that we shall need for proving the main results of this paper.

We consider the problem of the bifurcation of T -periodic solutions from differential systems of the form

$$\mathbf{x}' = F_0(t, \mathbf{x}) + \varepsilon F_1(t, \mathbf{x}) + \varepsilon^2 F_2(t, \mathbf{x}, \varepsilon), \quad (\text{A.1})$$

with $\varepsilon = 0$ to $\varepsilon \neq 0$ sufficiently small. Here the functions $F_0, F_1 : \mathbb{R} \times \Omega \rightarrow \mathbb{R}^n$ and $F_2 : \mathbb{R} \times \Omega \times (-\varepsilon_0, \varepsilon_0) \rightarrow \mathbb{R}^n$ are $\mathcal{C}^2$ functions, T -periodic in the first variable, and Ω is an open subset of $\mathbb{R}^n$. The main assumption is that the unperturbed system

$$\mathbf{x}' = F_0(t, \mathbf{x}) \quad (\text{A.2})$$

has a submanifold of dimension n of periodic solutions. A solution of this problem is given using the averaging theory.

Let $\mathbf{x}(t, \mathbf{z}, \varepsilon)$ be the solution of the system (A.2) such that $\mathbf{x}(0, \mathbf{z}, \varepsilon) = \mathbf{z}$. We write the linearization of the unperturbed system along the periodic solution $\mathbf{x}(t, \mathbf{z}, 0)$ as

$$\mathbf{y}' = D_{\mathbf{x}}F_0(t, \mathbf{x}(t, \mathbf{z}, 0))\mathbf{y}. \quad (\text{A.3})$$

In what follows we denote by $M_{\mathbf{z}}(t)$ some fundamental matrix of the linear differential system (A.3).

We assume that there exists an open set V with $\text{Cl}(V) \subset \Omega$ such that for each $\mathbf{z} \in \text{Cl}(V)$, $\mathbf{x}(t, \mathbf{z}, 0)$ is T -periodic. The set $\text{Cl}(V)$ is *isochronous* for the system (A.1); i.e. it is a set formed only by periodic orbits, all of them having the same period. Then, an answer to the problem of the bifurcation of T -periodic solutions from the periodic solutions $\mathbf{x}(t, \mathbf{z}, 0)$ contained in $\text{Cl}(V)$ is given in the following result.

Theorem A.1 (Perturbations of an isochronous set). *Assume that there exists an open and bounded set V with $\text{Cl}(V) \subset \Omega$ such that for each $\mathbf{z} \in \text{Cl}(V)$, the solution $\mathbf{x}(t, \mathbf{z}, 0)$ is T -periodic. Consider the function*

$$\mathcal{F} : \text{Cl}(V) \rightarrow \mathbb{R}^n, \quad \mathcal{F}(\mathbf{z}) = \int_0^T M_{\mathbf{z}}^{-1}(t) F_1(t, \mathbf{x}(t, \mathbf{z}, 0)) dt.$$

If there exists $\alpha \in V$ with $\mathcal{F}(\alpha) = 0$ and $\det((d\mathcal{F}/d\mathbf{z})(\alpha)) \neq 0$, then there exists a T -periodic solution $\mathbf{x}(t, \varepsilon)$ of system (A.1) such that $\mathbf{x}(0, \varepsilon) \rightarrow \alpha$ as $\varepsilon \rightarrow 0$.

Theorem A.1 goes back to Malkin [14] and Roseau [17], for a shorter proof see Buica et al. [3].

B Extended Complete Chebyshev System

We say that the functions $(f_0, \dots, f_n)$ defined on the interval I form an *Extended Chebyshev system* or ET-system on I if and only if any nontrivial linear combination of these functions has at most n zeros counting their multiplicities and this number is reached. The functions $(f_0, \dots, f_n)$ are an *Extended Complete Chebyshev system* or an ECT-system on I if and only if, for any $k \in \{0, 1, \dots, n\}$, $(f_0, \dots, f_k)$ form an ET-system.

Theorem B.1. *Let $f_0, \dots, f_n$ be analytic functions defined on an open interval $I \subset \mathbb{R}$. Then $(f_0, \dots, f_n)$ is an ECT-system on I if and only if for each $k \in \{0, 1, \dots, n\}$ and all $y \in I$ the Wronskian*

$$W(f_0, \dots, f_k)(y) = \begin{vmatrix} f_0(y) & f_1(y) & \cdots & f_k(y) \\ f_0'(y) & f_1'(y) & \cdots & f_k'(y) \\ \vdots & \vdots & \ddots & \vdots \\ f_0^{(k)}(y) & f_1^{(k)}(y) & \cdots & f_k^{(k)}(y) \end{vmatrix}$$

is different from zero.

For a proof of Theorem B.1 see [10].

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