Research Article Open Access

Bang-Xian Han and Andrea Mondino*

Angles between Curves in Metric Measure Spaces

https://doi.org/10.1515/agms-2017-0003 Received April 13, 2017; accepted July 11, 2017

Abstract: The goal of the paper is to study the angle between two curves in the framework of metric (and metric measure) spaces. More precisely, we give a new notion of angle between two curves in a metric space. Such a notion has a natural interplay with optimal transportation and is particularly well suited for metric measure spaces satisfying the curvature-dimension condition. Indeed one of the main results is the validity of the cosine formula on $\mathrm{RCD}^*(K, N)$ metric measure spaces. As a consequence, the new introduced notions are compatible with the corresponding classical ones for Riemannian manifolds, Ricci limit spaces and Alexandrov spaces.

Keywords: angle, metric measure space, Wasserstein space, curvature dimension condition, Ricci curvature

MSC: 53C23, 30L99, 31E05

1 Introduction

The 'angle' between two curves is a fundamental concept of mathematics, which aims to quantify the infinitesimal distance between two crossing curves at a crossing point. Such a notion is classical in Euclidean and in Riemannian geometries where a global (respectively infinitesimal) scalar product is given: the cosine of the angle between two crossing curves is by definition the scalar product of the velocity vectors. If the space is not given an infinitesimal scalar product, it is a challenging problem to define angles in a sensible way. In this paper, we will study this problem in a metric (measure) sense. More precisely, consider a metric space (X, d), a point $p \in X$, and two geodesics γ , η such that $\gamma_0 = \eta_0 = p$. Our task is to propose a meaningful definition of the angle between the curves γ , η at the point p, denoted by $\angle \gamma p \eta$, and to establish some interesting properties.

We recall some examples first. Assume that γ and η are geodesics, and the space (X, d) is an Alexandrov space, with upper or lower curvature bounds. From the monotonicity implied by the Alexandrov condition, it is known (see for instance [7]) that the angle $\angle \gamma p \eta$ is well defined by the cosine formula:

$$\angle \gamma p \eta = \lim_{s,t \to 0} \arccos \frac{s^2 + t^2 - \mathrm{d}^2(\gamma_s,\eta_t)}{2st} = \lim_{t \to 0} \arccos \frac{2t^2 - \mathrm{d}^2(\gamma_t,\eta_t)}{2t^2}.$$

In order to define the angle for geodesics in a more general framework, a crucial observation is that a geodesic can be seen as gradient flow of the distance function, i.e. a geodesic γ 'represents' the gradient of $-d(\gamma_0, \gamma_1) d(\gamma_1, \cdot)$ on each point γ_t . Inspired by the seminal work of De Giorgi on gradient flows [15], given an arbitrary metric space (X, d) with a geodesic $\gamma : [0, 1] \to X$ and a Lipschitz function $f : X \to \mathbb{R}$, we say that γ represents ∇f at time 0, or γ represents the gradient of the function f at the point $p = \gamma(0)$ if the following

A.Mondino@warwick.ac.uk

Bang-Xian Han: University of Bonn, Germany, E-mail: hanbangxian@gmail.com

^{*}Corresponding Author: Andrea Mondino: University of Warwick, Coventry, United Kingdom, E-mail:

inequality holds

$$\underline{\lim_{t\to 0}}\frac{f(\gamma_t)-f(\gamma_0)}{t}\geq \frac{1}{2}|\operatorname{lip}(f)|^2+\frac{1}{2}|\dot{\gamma}|^2,$$

where $|\dot{\gamma}| = d(\gamma_0, \gamma_1)$ is the (constant, metric) speed of the geodesic γ . Notice that the opposite inequality

$$\overline{\lim_{t\to 0}} \frac{f(\gamma_t) - f(\gamma_0)}{t} \le \frac{1}{2} |\operatorname{lip}(f)|^2 + \frac{1}{2} |\dot{\gamma}|^2$$

is always true by Leibniz rule and Cauchy-Schwartz inequality. Hence γ represents ∇f at time 0 if and only if the equality holds. It is easily seen that the geodesic γ always represents the gradient of $f_{\gamma}(\cdot) := -\mathrm{d}(\gamma_0, \gamma_1)\,\mathrm{d}(\gamma_1, \cdot)$ at the point γ_0 (see for instance Lemma 3.5). We then say that the angle $\angle \gamma p\eta$ between two geodesics γ , η with $\gamma_0 = \eta_0 = p$ exists if the limit $\lim_{t\downarrow 0} \frac{f_{\gamma}(\eta_t) - f_{\gamma}(\eta_0)}{t}$ exists. In this case we set

$$\angle \gamma p \eta := \arccos\left(\frac{1}{|\dot{\gamma}||\dot{\eta}|} \lim_{t\downarrow 0} \frac{f_{\gamma}(\eta_t) - f_{\gamma}(\eta_0)}{t}\right). \tag{1.1}$$

Notice that in case (X, d) is the metric space associated to a smooth Riemannian manifold (M, g), the definition (1.1) reduces to the familiar notion of angle

$$\angle \gamma p \eta = \arccos g_p \left(\frac{\nabla f_{\gamma}(p)}{|\dot{\gamma}_0|}, \frac{\dot{\eta}_0}{|\dot{\eta}_0|} \right) = \arccos g_p \left(\frac{\dot{\gamma}_0}{|\dot{\gamma}_0|}, \frac{\dot{\eta}_0}{|\dot{\eta}_0|} \right).$$

Besides the case of Alexandrov spaces, a class of spaces where the angle is particularly well behaved is the one of Lipschitz-infinitesimally Hillbertian spaces. By definition, a metric measure space $(X, \mathbf{d}, \mathbf{m})$ is Lipschitz-infinitesimally Hillbertian if for any pair of Lipschitz functions $f, g: X \to \mathbb{R}$ both the limits for $\varepsilon \to 0$ of $\frac{|\operatorname{lip}(f+\varepsilon g)|^2(x)-|\operatorname{lip}(f)|^2(x)}{2\varepsilon}$ and $\frac{|\operatorname{lip}(g+\varepsilon f)|^2(x)-|\operatorname{lip}(g)|^2(x)}{2\varepsilon}$ exist and are equal for \mathfrak{m} -a.e. $x \in X$, where $\operatorname{lip}(f)$ is the local Lipschitz constant of f (for the standard definition see (2.1)). A remarkable example of Lipschitz-infinitesimally Hillbertian spaces is given by the RCD $^*(K,N)$ -spaces, a class of metric measure spaces satisfying Ricci curvature lower bound by $K \in \mathbb{R}$ and dimension upper bound by $N \in (1,\infty)$ in a synthetic sense such that the Laplacian is linear, and which include as notable subclasses the Alexandrov spaces with curvature bounded below and the Ricci limit spaces (i.e. pointed measured Gromov-Hausdorff limits of sequences of Riemannian manifolds with uniform lower Ricci curvature bounds).

In the class of Lipschitz-infinitesimally Hillbertian spaces, the second author [26] introduced a notion of 'angle between three points'; more precisely for every fixed pair of points $p, q \in X$, for \mathfrak{m} -a.e. $x \in X$ the angle $\angle pxq$ given by the formula

$$[0, \pi] \ni \angle pxq := \arccos\left(\lim_{\varepsilon \to 0} \frac{|\operatorname{lip}(r_p + \varepsilon r_q)|^2(x) - |\operatorname{lip}(r_p)|^2(x)}{2\varepsilon}\right), \tag{1.2}$$

is well defined, unique, and symmetric in p and q. Here $r_p(\cdot) := \operatorname{d}(p, \cdot)$ is the distance function from p. A first result of the present paper is to relate the angle between three points with the angle between two geodesics: in Theorem 3.9 we prove that if the angle $\angle pxq$ exists in the sense of [26] then also the angle between the geodesics γ^{xp} , γ^{xq} joining x to p and x to q exists and coincides with the angle between the three points, i.e. $\angle \gamma^{xp}x\gamma^{xq} = \angle pxq$. In particular it follows that in a Lipschitz-infinitesimally Hilbertian geodesic space the angle between two geodesics in well defined in an a.e. sense.

An important class of metric spaces are the spaces of probability measures over metric spaces endowed with the quadratic transportation distance: given a metric space (X, d) denote by $\mathcal{W}_2 := (\mathcal{P}_2(X), W_2)$ the corresponding Wasserstein space. By using ideas similar to the ones above, together with Otto Calculus (see [28]) and the calculus tools developed by Ambrosio-Gigli-Savaré [2] and Gigli [17], in Subsection 3.3 we study in detail the angle between two geodesics in \mathcal{W}_2 . In particular if the underlying space (X, d, \mathfrak{m}) is an RCD^{*}(K, N) space, we get the angle $\angle pxq$ between three points as the limit of the angle between the geodesics in \mathcal{W}_2 obtained by joining geodesically diffused approximations of Dirac masses centered at p, x and y (see Proposition 3.15 for the precise statement; see also Proposition 3.17 for a more detailed link with the optimal transport picture).

Besides the case of Alexandrov spaces, another class of spaces where the notion of angle is quite well understood is given by Ricci limit spaces. Indeed it was proved by Honda [23] that if (X, d, \mathfrak{m}) is a Ricci-limit

space, then for \mathfrak{m} -a.e. $p \in X$ the angle between two geodesics is well defined and it satisfies the following single-variable cosine formula:

$$\cos \angle \gamma p \eta = \lim_{t \to 0} \frac{2t^2 - d^2(\gamma_t, \eta_t)}{2t^2}.$$
(1.3)

One of the main goals of the present paper is to extend the validity of the formula (1.3) to metric measure spaces satisfying Ricci curvature lower bounds in a synthetic sense, the so-called $RCD^*(K, N)$ -spaces (for the definition and basic properties of such spaces see Section 2 and references therein). This is the content of the next theorem (corresponding to Theorem 4.4 in the body of the manuscript), which is one of the main results of the paper.

Theorem 1.1 (Cosine formula for angles in RCD^{*}(K, N) spaces). Let (X, d, m) be an RCD^{*}(K, N) space and fix p, $q \in X$. Then for m-a.e. $x \in X$ there exist unique geodesics from x to p and from x to q denoted by γ^{xp} , $\gamma^{xq} \in \text{Geo}(X)$ and

$$\angle \gamma^{xp} x \gamma^{xq} = \angle pxq = \lim_{t \to 0} \arccos \frac{2t^2 - d^2(\gamma_t^{xp}, \gamma_t^{xq})}{2t^2}, \quad \text{for m-a.e. } x.$$
 (1.4)

The proof of Theorem 1.1 is independent and different from the one given by Honda [23] for Ricci limit spaces: indeed Honda argues by getting estimates on the smooth approximating manifolds and then passes to the limit, while our proof for $RCD^*(K, N)$ spaces goes by arguing directly on the non smooth space (X, d, m). More precisely, we perform a blow up argument centered at x and use that for m-a.e. x the tangent cone is unique and euclidean [21, 27]. From the technical point of view we also make use of the fine convergence results for Sobolev functions proved in [4, 20], and we prove estimates on harmonic approximations of distance functions (see in particular Proposition 4.3). Harmonic approximations of distance functions are well known for smooth Riemannian manifolds with lower Ricci curvature bounds, and are indeed one of the key technical tools in the Cheeger-Colding theory of Ricci limit spaces [12–14]; on the other hand for non-smooth $RCD^*(K, N)$ -spaces it seems they have not yet appeared in the literature, and we expect them to be a useful technical tool in the future development of the field.

As a consequence of Theorem 1.1, we get that our definition of angle between two geodesics agrees (at least in a.e. sense) with the Alexandrov's definition in case (X, d) is an Alexandrov space, and with the Honda's definition [23] in case (X, d, \mathfrak{m}) is a Ricci limit space.

2 Preliminaries

2.1 Metric measure spaces

Let (X, d) be a complete metric space. A continuous map $\gamma : [0, 1] \mapsto X$ will be called *curve*. The space of curves defined on [0, 1] with values in X is denoted by C([0, 1], X). The space C([0, 1], X) equipped with the uniform distance is a complete metric space.

We define the length of γ by

$$l[\gamma] := \sup_{\tau} \sum_{i=1}^{n} d(\gamma_{t_{i-1}}, \gamma_{t_i})$$

where $\tau := \{0 = t_0, t_1, ..., t_n = 1\}$ is a partition of [0, 1], and the sup is taken over all finite partitions. The space (X, d) is said to be a *length space* if for any $x, y \in X$ we have

$$d(x, y) = \inf_{\gamma} l[\gamma],$$

where the infimum is taken over all $\gamma \in C([0, 1], X)$ connecting x and y. A geodesic from x to y is a curve γ such that:

$$d(\gamma_s, \gamma_t) = |s - t| d(\gamma_0, \gamma_1), \quad \forall t, s \in [0, 1], \quad \gamma_0 = x, \gamma_1 = y.$$

The space of all geodesics on X will be denoted by Geo(X). It is a closed subset of C([0, 1], X).

Given $p \in [1, +\infty]$ and a curve γ , we say that γ belongs to $AC^p([0, 1], X)$ if

$$d(\gamma_s, \gamma_t) \leq \int_s^t G(r) dr, \quad \forall t, s \in [0, 1], s < t,$$

for some $G \in L^p([0, 1])$. In particular, the case p = 1 corresponds to absolutely continuous curves, whose class is denoted by AC([0, 1], X). It is known that for $\gamma \in AC([0, 1], X)$, there exists an a.e. minimal function G satisfying this inequality, called metric derivative and denoted by $|\dot{\gamma}|$. The metric derivative of γ can be computed for a.e. $t \in [0, 1]$ as

$$|\dot{\gamma}_t| := \lim_{h \to 0} \frac{\mathrm{d}(\gamma_{t+h}, \gamma_t)}{|h|}.$$

It is known that (see for example [7]) the length of a curve $\gamma \in AC([0,1],X)$ can be computed as

$$l[\gamma] := \int_0^1 |\dot{\gamma}_t| \, \mathrm{d}t.$$

In particular, on a length space X we have

$$d(x, y) = \inf_{\gamma} \int_{0}^{1} |\dot{\gamma}_{t}| dt$$

where the infimum is taken among all $\gamma \in AC([0, 1], X)$ which connect x and y.

Given $f: X \mapsto \mathbb{R}$, the local Lipschitz constant $lip(f): X \mapsto [0, \infty]$ is defined as

$$\operatorname{lip}(f)(x) := \overline{\lim_{y \to x} \frac{|f(y) - f(x)|}{\operatorname{d}(x, y)}} \tag{2.1}$$

if x is not isolated, 0 otherwise, while the (global) Lipschitz constant is defined as

$$\operatorname{Lip}(f) := \sup_{x \neq y} \frac{|f(y) - f(x)|}{\operatorname{d}(x, y)}.$$

If (X, d) is a length space, we have $Lip(f) = \sup_{x} lip(f)(x)$.

We are not only interested in metric structures, but also in the interaction between metric and measure. For the metric measure space (X, d, m), basic assumptions used in this paper are:

Assumption 2.1. The metric measure space (X, d, m) satisfies:

- (*X*, d) is a complete and separable length space,
- m is a non-negative Borel measure with respect to d and finite on bounded sets,
- $\operatorname{supp} \mathfrak{m} = X$.

In this paper, we will often assume that the metric measure space (X, d, m) satisfies the $RCD^*(K, N)$ condition, for some $K \in \mathbb{R}$ and $N \in [1, \infty]$ (when $N = \infty$ it is denoted by $RCD(K, \infty)$). The $RCD(K, \infty)$ and $RCD^*(K, N)$ conditions are refinements of the curvature-dimensions proposed by Lott-Sturm-Villani (see [25] and [30, 31] for $CD(K, \infty)$), and Bacher-Sturm (see [9] for $CD^*(K, N)$) in order to isolate the non-smooth 'Riemannian' structures from the 'Finslerian' ones. More precisely, the RCD conditions are obtained by reinforcing the corresponding CD conditions by adding the requirement that the Sobolev space $W^{1,2}(X, d, m)$ is a Hilbert space (see the next subsection for more details). It is then clear that the following relations hold

$$\mathrm{RCD}^{\star}(K,N)\subset\mathrm{CD}^{\star}(K,N)$$
 and $\mathrm{RCD}(K,\infty)\subset\mathrm{CD}(K,\infty)$;

moreover one has that

$$\mathrm{RCD}^{\star}(K,N) \subset \mathrm{RCD}(K,\infty)$$
 and $\mathrm{CD}^{\star}(K,N) \subset \mathrm{CD}(K,\infty)$.

It is known that, for finite $N \in [1, \infty)$, a $CD^*(K, N)$ space (X, d, m) satisfies the following properties:

- (X, d, m) is locally doubling and therefore a locally compact space, [9];
- (*X*, d, m) supports a local Poincaré inequality, [29].

For more details about $RCD(K, \infty)$ and $RCD^*(K, N)$ spaces, we refer to [1, 3, 5, 16].

2.2 Optimal transport and Sobolev functions

The set of Borel probability measures on (X, d) will be denoted by $\mathcal{P}(X)$. We also use $\mathcal{P}_2(X) \subseteq \mathcal{P}(X)$ to denote the set of measures with finite second moment, i.e. $\mu \in \mathcal{P}_2(X)$ if $\mu \in \mathcal{P}(X)$ and $\int d^2(x, x_0) d\mu(x) < +\infty$ for some (and thus every) $x_0 \in X$. For $t \in [0, 1]$, the evaluation map $e_t : C([0, 1], X) \to X$ is given by

$$e_t(\gamma) := \gamma_t, \quad \forall \gamma \in C([0, 1], X).$$

The space $\mathcal{P}_2(X)$ is naturally endowed with the quadratic transportation distance W_2 defined by:

$$W_2^2(\mu, \nu) := \inf_{\pi} \int_{X \times X} d^2(x, y) \, d\pi(x, y), \tag{2.2}$$

where the inf is taken among all couplings $\pi \in \mathcal{P}(X \times X)$ with marginals μ and ν , i.e. $(P_1)_{tt}\pi = \mu$ and $(P_2)_{tt}\pi = \nu$ where P_i , i = 1, 2 are the projection maps onto the first and second coordinate respectively. The metric space $(\mathcal{P}_2(X), W_2)$ will be denoted by W_2 . Let us recall that the infimum in the Kantorovich problem (2.2) is always attained by an optimal coupling π . We denote the set of optimal couplings between μ and ν by $Opt(\mu, \nu)$. Below we recall some fundamental properties of the metric space W_2 we will use throughout the paper.

Proposition 2.2 (Geodesics in the Wasserstein space). Let (X, d) be a metric space and fix $\mu_0, \mu_1 \in \mathcal{P}_2(X)$. Then the curve $(\mu_t)_{t\in[0,1]}\subset \mathcal{W}_2$ is a constant speed geodesic connecting μ_0 and μ_1 , i.e. it satisfies

$$W_2(\mu_s, \mu_t) = |s - t| W_2(\mu_0, \mu_1), \quad \forall s, t \in [0, 1]$$
(2.3)

if and only if there exists $\Pi \in \mathcal{P}(Geo(X)) \subseteq \mathcal{P}(C([0,1],X))$, called optimal dynamical plan (or simply optimal plan), such that

$$\mu_t = (e_t)_{\sharp} \Pi \ \forall t \in [0, 1] \ and \ (e_0, e_1)_{\sharp} \Pi \in Opt(\mu_0, \mu_1).$$

The set of optimal dynamical plans from μ_0 to μ_1 is denoted by OptGeo(μ_0, μ_1). Moreover, if X is a geodesic space, then W_2 is also geodesic.

Absolutely continuous curves in W_2 are characterized by the following theorem:

Theorem 2.3 (Superposition principle, [24]). Let (X, d) be a complete and separable metric space and let $(\mu_t) \in AC^2([0,1], \mathcal{P}_2(X))$. Then there exists a measure $\Pi \in \mathcal{P}(C([0,1],X))$ concentrated on $AC^2([0,1],X)$ such that:

$$(\mathbf{e}_t)_{\sharp}\Pi = \mu_t, \quad \forall t \in [0, 1]$$

$$\int |\dot{\gamma}_t|^2 \, \mathrm{d}\Pi(\gamma) = |\dot{\mu}_t|^2, \quad \textit{for a.e. } t \in [0, 1].$$

Moreover, the infimum of the energy $\int_0^1 \int |\dot{\gamma}_t|^2 d\Pi'(\gamma) dt$ among all the $\Pi' \in \mathcal{P}(\mathcal{C}([0,1],X))$ satisfying $(e_t)_H \Pi' = \mu_t$ for every $t \in [0, 1]$ is attained by such Π .

Definition 2.4 (Test plan). Let (X, d, \mathfrak{m}) be a metric measure space and $\Pi \in \mathcal{P}(C([0, 1], X))$. We say that $\Pi \in \mathcal{P}(C([0, 1], X))$. $\mathcal{P}(C([0,1],X))$ has *bounded compression* provided there exists C>0 such that

$$(e_t)_{\sharp}\Pi \leq C\mathfrak{m}, \quad \forall t \in [0,1].$$

We say that Π is a *test plan* if it has bounded compression, is concentrated on $AC^2([0, 1], X)$ and

$$\int_0^1 \int |\dot{\gamma}_t|^2 d\Pi(\gamma) dt < +\infty.$$

The notion of Sobolev function is given in duality with that of test plan:

Definition 2.5 (The Sobolev class $S^2(X)$). Let (X, d, m) be a metric measure space. A Borel function $f: X \to \mathbb{R}$ belongs to the Sobolev class $S^2(X)$ (resp. $S^2_{loc}(X)$) provided there exists a non-negative function $G \in L^2(X, \mathfrak{m})$ (resp. $L^2_{loc}(X, \mathfrak{m})$) such that

$$\int |f(\gamma_1) - f(\gamma_0)| \, \mathrm{d}\Pi(\gamma) \le \int \int_0^1 G(\gamma_s) |\dot{\gamma}_s| \, \mathrm{d}s \, \mathrm{d}\Pi(\gamma), \quad \forall \text{ test plan } \Pi.$$

In this case, G is called a 2-weak upper gradient of f, or simply weak upper gradient.

It is known, see e.g. [2], that there exists a minimal function G in the m-a.e. sense among all the weak upper gradients of f. We denote such minimal function by |Df| or $|Df|_X$ to emphasize which space we are considering and call it *minimal weak upper gradient*. Notice that if f is Lipschitz, then $|Df| \le \text{lip}(f)$ m-a.e., because lip(f)is a weak upper gradient of f.

It is known that the locality holds for |Df|, i.e. |Df| = |Dg| m-a.e. on the set $\{f = g\}$, moreover $S_{loc}^2(X, d, m)$ is a vector space and the inequality

$$|D(\alpha f + \beta g)| \le |\alpha||Df| + |\beta||Dg|, \quad \mathfrak{m}\text{-a.e.}, \tag{2.4}$$

holds for every $f, g \in S^2_{loc}(X, d, m)$ and $\alpha, \beta \in \mathbb{R}$. Moreover, the space $S^2_{loc} \cap L^{\infty}_{loc}(X, d, m)$ is an algebra, with the inequality

$$|D(fg)| \le |f||Dg| + |g||Df|, \quad \text{m-a.e.},$$
 (2.5)

being valid for any $f,g\in S^2_{loc}\cap L^\infty_{loc}(X,\operatorname{\mathtt{d}},\mathfrak{m})$.

The Sobolev space $W^{1,2}(X, \mathbf{d}, \mathbf{m})$, also denoted by $W^{1,2}(X)$ for short, is defined as

$$W^{1,2}(X) := S^2(X, d, \mathfrak{m}) \cap L^2(X, \mathfrak{m})$$

and is endowed with the norm

$$||f||_{W^{1,2}(X)}^2 := ||f||_{L^2(X,\mathfrak{m})}^2 + |||Df|||_{L^2(X,\mathfrak{m})}^2.$$

 $W^{1,2}(X)$ is always a Banach space, but in general it is not a Hilbert space. (X, d, \mathfrak{m}) is said infinitesimally Hilbertian if $W^{1,2}(X)$ is a Hilbert space.

On an infinitesimally Hilbertian space, we have a natural pointwise inner product $\langle \nabla \cdot, \nabla \cdot \rangle$: $[W^{1,2}(X)]^2 \mapsto L^1(X)$ defined by

$$\langle \nabla f, \nabla g \rangle := \frac{1}{4} \Big(|\mathrm{D}(f+g)|^2 - |\mathrm{D}(f-g)|^2 \Big).$$

In order to prove the cosine formula we will use properties of harmonic functions in open sets of a m.m. space. Let us define the relevant quantities and recall the properties we will use; for simplicity, as always we assume the space (X, d) to be proper, complete and separable, and the measure \mathfrak{m} to be finite on bounded sets (this indeed is the geometric case correspoding to $RCD^*(K, N)$ spaces, for $N < \infty$ we will be interested in). For the general case see for instance [6, 17, 19].

Definition 2.6 (Sobolev classes in Ω). Let (X, d, m) be a m.m. space and let $\Omega \subset X$ be an open subset. The space $S^2(\Omega)$ is the space of Borel functions $f:\Omega\to\mathbb{R}$ such that $\chi f\in S^2(X)$ for any Lipschitz function $\chi: X \to [0, 1]$ such that supp $\chi \subset \Omega$, where χf is taken 0 by definition on $X \setminus \Omega$. Let $W^{1,2}(\Omega) := L^2(\Omega) \cap S^2(\Omega)$ be the corresponding Sobolev space endowed with the natural norm, and denote by $W_0^{1,2}(\Omega) \subset W^{1,2}(X)$ the closure of compactly supported Lipschitz functions on Ω .

Definition 2.7 (Measure valued Laplacian). Let (X, d, m) be a m.m. space, $\Omega \subset X$ an open subset and f: $\Omega \to \mathbb{R}$ a Borel function. We say that f is in the domain of the Laplacian in Ω , and write $f \in D(\Lambda, \Omega)$ provided $f \in S^2(\Omega)$ and there exists a locally finite Borel measure μ on Ω such that for any $\varphi \in LIP(X)$ with compact support contained in Ω it holds

$$\int_X \varphi \, \mathrm{d}\mu = -\int_X \langle \nabla \varphi, \nabla f \rangle \, \mathrm{d}\mathfrak{m}.$$

In this case the measure μ is unique and we denote it by $\Delta f_{\perp}\Omega$, or simply Δf . If $\Delta f_{\perp}\Omega \ll \mathfrak{m}$, we denote its density with respect to \mathfrak{m} by $\Delta f_{\perp}\Omega$ or simply by Δf .

A function $f \in D(\Delta, \Omega)$ is said to be *harmonic in* Ω , or simply harmonic, if $\Delta f \cup \Omega = 0$.

For simplicity we state the next proposition for $RCD^*(K, N)$ space, though it is valid more generally for doubling spaces supporting a weak-local 1-2 Poincaré inequality (see [6] for details).

Proposition 2.8. Let (X, d, \mathfrak{m}) be a $\mathrm{RCD}^*(K, N)$ space, for some $K \in \mathbb{R}$ and $N \in [1, \infty)$, and let $\Omega \subset X$ be a bounded open set. Then the following properties hold.

- i) Regularity. Let $f: \Omega \to \mathbb{R}$ be harmonic in Ω . Then f admits a continuous representative (actually even locally Lipschitz).
- ii) Comparison. If $f, g \in D(\Delta, \Omega)$ are such that $f \in W_0^{1,2}(\Omega)$, $|g| \le C$ m-a.e. on Ω for some $C \in \mathbb{R}$ and $\Delta(f+g) \ge 0$ then $f \le 2C$ m-a.e. on Ω .
- *ii)* Existence and uniqueness of harmonic functions. Assume that $\mathfrak{m}(X \setminus \Omega) > 0$ and let $f \in W^{1,2}(X)$. Then there exists a unique harmonic function g on Ω such that $f g \in W^{1,2}_0(\Omega)$.
- iv) Strong maximum principle. Let $f: \Omega \to \mathbb{R}$ be harmonic in Ω and assume that its continuous representative has a maximum at a point $x_0 \in \Omega$. Then f is constant on the connected component of Ω containing x_0 .

In order to state the Laplacian Comparison Theorem, let us introduce the coefficients $\tilde{\sigma}_{K,N}(\cdot):[0,\infty)\to\mathbb{R}$ defined by

$$\tilde{\sigma}_{K,N}(\theta) := \left\{ \begin{array}{ll} \theta \sqrt{\frac{K}{N}} \cot \alpha \left(\theta \sqrt{\frac{K}{N}}\right), & \text{if } K > 0, \\ 1 & \text{if } K = 0, \\ \theta \sqrt{-\frac{K}{N}} \coth \left(\theta \sqrt{-\frac{K}{N}}\right), & \text{if } K < 0. \end{array} \right.$$

Theorem 2.9 (Laplacian comparison, [17]). *Let* (X, d, \mathfrak{m}) *be an* $RCD^*(K, N)$ *space for some* $K \in \mathbb{R}$ *and* $N \in (1, \infty)$. *Then*

$$\frac{\mathrm{d}^2(x_0,\cdot)}{2}\in\mathrm{D}(\Delta,X)\quad with\quad \Delta\frac{\mathrm{d}^2(x_0,\cdot)}{2}\leq N\,\tilde{\sigma}_{K,N}(\mathrm{d}(x_0,\cdot))\,\mathfrak{m}\quad\forall x_0\in X$$

and

$$d(x_0,\cdot)\in D(\Delta,X\setminus\{x_0\}) \quad with \quad \Delta d(x_0,\cdot) \perp X\setminus\{x_0\} \leq \frac{N\,\tilde{\sigma}_{K,N}(d(x_0,\cdot))-1}{d(x_0,\cdot)}\,\mathfrak{m} \quad \forall x_0\in X.$$

2.3 Pointed measured Gromov-Hausdorff convergence and convergence of functions

In order to study the convergence of possibly non-compact metric measure spaces, it is useful to fix reference points. We then say that $(X, d, \mathfrak{m}, \bar{x})$ is a pointed metric measure space, p.m.m.s. for short, if (X, d, \mathfrak{m}) is a m.m.s. as before and $\bar{x} \in X$ plays the role of reference point. Recall that, for simplicity, we always assume supp $\mathfrak{m} = X$. We will adopt the following definition of convergence of p.m.m.s. (see [7], [20] and [32]):

Definition 2.10 (Pointed measured Gromov-Hausdorff convergence). A sequence $(X_j, \mathbf{d}_j, \mathfrak{m}_j, \bar{x}_j)$ is said to converge in the pointed measured Gromov-Hausdorff topology (p-mGH for short) to $(X_{\infty}, \mathbf{d}_{\infty}, \mathfrak{m}_{\infty}, \bar{x}_{\infty})$ if there exists a separable metric space (Z, \mathbf{d}_Z) and isometric embeddings $\{\iota_j : (X_j, \mathbf{d}_j) \to (Z, \mathbf{d}_Z)\}_{i \in \mathbb{N}}$ such that for every $\varepsilon > 0$ and R > 0 there exists j_0 such that for every $j > j_0$

$$\iota_{\infty}(B_R^{X_{\infty}}(\bar{x}_{\infty})) \subset B_{\varepsilon}^Z[\iota_j(B_{R+\varepsilon}^{X_j}(\bar{x}_j))]$$
 and $\iota_j(B_R^{X_j}(\bar{x}_j)) \subset B_{\varepsilon}^Z[\iota_{\infty}(B_{R+\varepsilon}^{X_{\infty}}(\bar{x}_{\infty}))],$

where $B_{\varepsilon}^{Z}[A] := \{z \in Z : d_{Z}(z, A) < \varepsilon\}$ for every subset $A \subset Z$, and

$$\lim_{j\to\infty}\int_{Z}\varphi\,\mathrm{d}((\iota_{j})_{\sharp}(\mathfrak{m}_{j}))=\int_{Z}\varphi\,\mathrm{d}((\iota_{\infty})_{\sharp}(\mathfrak{m}_{\infty}))\qquad\forall\varphi\in C_{b}(Z),$$

where $C_b(Z)$ denotes the set of real valued bounded continuous functions with bounded support in Z.

Sometimes in the following, for simplicity of notation, we will identify the spaces X_j with their isomorphic copies $\iota_j(X_j) \subset Z$. It is obvious that this is in fact a notion of convergence for isomorphism classes of p.m.m.s., moreover it is induced by a metric (see e.g. [20] for details).

Next, following [20], we recall various notions of convergence of functions defined on p-mGH converging spaces.

Definition 2.11 (Pointwise convergence of scalar valued functions). Let $(X_j, \mathbf{d}_j, \mathfrak{m}_j, \bar{x}_j)$, $j \in \mathbb{N} \cup \{\infty\}$ be a p-mGH converging sequence of p.m.m.s. and let $f_j : X_j \mapsto \mathbb{R}, j \in \mathbb{N} \cup \{\infty\}$ be a sequence of functions. We say that f_j converge pointwise to f_∞ provided:

$$f_i(x_i) \to f_{\infty}(x_{\infty})$$
 for every sequence of points $x_i \in X_i$ such that $\iota_i(x_i) \to \iota_{\infty}(x_{\infty})$ in (Z, d_Z) .

If for any $\epsilon > 0$ there exists $N \in \mathbb{N}$ such that $|f_j(x_j) - f_\infty(x_\infty)| \le \epsilon$ for every $j \ge N$ and every $x_j \in X_j$, $x_\infty \in X_\infty$ with $\mathrm{d}_Z(\iota_j(x_j), \iota_\infty(x_\infty)) \le \frac{1}{N}$, we say that $f_j \to f_\infty$ uniformly.

Definition 2.12 (L^2 weak and strong convergence). Let $(X_j, d_j, \mathfrak{m}_j, \bar{x}_j)$, $j \in \mathbb{N} \cup \{\infty\}$ be a p-mGH converging sequence of pointed metric measure spaces and let $f_i \in L^2(X_i, \mathfrak{m}_i)$, $j \in \mathbb{N} \cup \{\infty\}$ be a sequence of functions.

• We say that (f_j) converges weakly in L^2 to f_{∞} provided $(\iota_j)_{\sharp}(f_j \mathfrak{m}_j) \rightharpoonup (\iota_{\infty})_{\sharp}(f_{\infty} \mathfrak{m})$ weakly as Radon measures, i.e.

$$\int_{X_j} f_j(x) \ \varphi(\iota_j(x)) \ \mathrm{d}\mathfrak{m}_j(x) \to \int_{X_\infty} f_\infty(x) \ \varphi(\iota_\infty(x)) \ \mathrm{d}\mathfrak{m}_\infty(x), \qquad \forall \varphi \in C_b(Z),$$

and

$$\sup_{j\in\mathbb{N}}\int_{X_j}|f_j|^2\,\mathrm{d}\mathfrak{m}_j<\infty.$$

• We say that (f_j) converges *strongly* in L^2 to f_∞ provided it converges weakly in L^2 to f_∞ and moreover

$$\lim_{j\to\infty}\int_{X_i}|f_j|^2\,\mathrm{d}\mathfrak{m}_j=\int_{X_\infty}|f_\infty|^2\,\mathrm{d}\mathfrak{m}_\infty.$$

Definition 2.13 ($W^{1,2}$ weak and strong convergence). Let $(X_j, \mathbf{d}_j, \mathfrak{m}_j, \bar{x}_j)$, $j \in \mathbb{N} \cup \{\infty\}$ be a p-mGH converging sequence of pointed metric measure spaces and let $f_j \in W^{1,2}(X_j, \mathbf{d}_j, \mathfrak{m}_j)$, $j \in \mathbb{N} \cup \{\infty\}$ be a sequence of functions. We say that (f_i) converges *weakly in* $W^{1,2}$ to f_∞ if f_i are L^2 -weakly convergent to f and

$$\sup_{j\in\mathbb{N}}\int_{X_i}|\mathrm{D}f_j|^2\,\mathrm{d}\mathfrak{m}_j<\infty.$$

Strong convergence in $W^{1,2}$ is defined by requiring L^2 -strong convergence of the functions and that

$$\lim_{j\to\infty}\int_{X_i}|\mathrm{D}f_j|^2\,\mathrm{d}\mathfrak{m}_j=\int_{X_{i+1}}|\mathrm{D}f_\infty|^2\,\mathrm{d}\mathfrak{m}_\infty.$$

The next result proved in [4, Corollary 5.5] (see also [20, Corollary 6.10]) will be useful in the sequel.

Proposition 2.14. Let $(X_j, d_j, \mathfrak{m}_j, \bar{x}_j), j \in \mathbb{N} \cup \{\infty\}$ be a p-mGH converging sequence of pointed metric measure spaces. If for every $j \in \mathbb{N}$ one has $f_j \in W^{1,2}(X_i)$, $f_j \in D(\Delta_j, X_j)$ with $\Delta_j f_j$ uniformly bounded in L^2 , and (f_j) converges strongly in L^2 to f_∞ , then $f_\infty \in D(\Delta_\infty, X_\infty)$ and (f_j) converges to f_∞ strongly in $W^{1,2}$.

2.4 Euclidean tangent cones to $RCD^*(K, N)$ spaces

Let us first recall the notion of measured tangents. Let (X, d, m) be a m.m.s., $\bar{x} \in X$ and $r \in (0, 1)$; we consider the rescaled and normalized p.m.m.s. $(X, r^{-1}d, m_r^{\bar{x}}, \bar{x})$ where the measure $m_r^{\bar{x}}$ is given by

$$\mathfrak{m}_r^{\bar{x}} := \left(\int_{B_r(\bar{x})} 1 - \frac{1}{r} d(\cdot, \bar{x}) d\mathfrak{m} \right)^{-1} \mathfrak{m}. \tag{2.6}$$

Then we define:

Definition 2.15 (Tangent cone and regularity). Let (X, d, m) be a m.m.s. and $\bar{x} \in X$. A p.m.m.s. (Y, d_Y, n, y) is called a *tangent* to (X, d, m) at $\bar{x} \in X$ if there exists a sequence of rescalings $r_i \downarrow 0$ so that $(X, r_i^{-1}d, \mathfrak{m}_{r_i}^{\bar{X}}, \bar{x}) \rightarrow$ $(Y, d_Y, \mathfrak{n}, y)$ as $j \to \infty$ in the p-mGH sense. We denote the collection of all the tangents of (X, d, \mathfrak{m}) at $\bar{x} \in X$ by $Tan(X, d, m, \bar{X})$. A point $\bar{X} \in X$ is called *regular* if the tangent is unique and euclidean, i.e. if $Tan(X, d, m, \bar{X}) = Tan(X, d, m, \bar{X})$ $\{(\mathbb{R}^n, d_E, \mathcal{L}_n, 0^n)\}$, where d_E is the Euclidean distance and \mathcal{L}_n is the properly rescaled Lebesgue measure of \mathbb{R}^n .

The a.e. regularity was settled for Ricci-limit spaces by Cheeger-Colding [12–14]; for an $RCD^*(K, N)$ -space (X, d, m), it was proved in [21] that for m-a.e. $x \in X$ there exists a blow-up sequence converging to a Euclidean space. The \mathfrak{m} -a.e. uniqueness of the blow-up limit, together with the rectifiability of an $\mathrm{RCD}^*(K, N)$ -space, was then established in [27]. More precisely the following holds:

Theorem 2.16 (m-a.e. infinitesimal regularity of RCD $^*(K, N)$ -spaces). Let (X, d, m) be an RCD $^*(K, N)$ -space for some $K \in \mathbb{R}$, $N \in (1, \infty)$. Then \mathfrak{m} -a.e. $x \in X$ is a regular point, i.e. for \mathfrak{m} -a.e. $x \in X$ there exists $n = n(x) \in [1, N] \cap \mathbb{N}$ such that, for any sequence $r_i \downarrow 0$, the rescaled pointed metric measure spaces $(X, r_i^{-1}d, \mathfrak{m}_{r_i}^X, x)$ converge in the p-mGH sense to the pointed Euclidean space $(\mathbb{R}^n, d_E, \mathcal{L}_n, 0^n)$.

3 Definition of angle

3.1 Angle between three points

In [26], the second author proposed a notion of angle between three points $p, x, q \in X$ in a metric space (X, d). In general such an angle is not unique, the possible causes of non-uniqueness being a lack of regularity of the distance function (e.g., x is in the cut locus of p or q) or a lack of infinitesimal strict convexity of the distance function (for more details we refer to [26, Sections 1,2]). For simplicity, here we only treat the case when the angle is unique. Given two points $p, q \in X$, consider the distance functions

$$r_p(\cdot) := \mathrm{d}(p, \cdot), \quad r_q(\cdot) := \mathrm{d}(q, \cdot).$$
 (3.1)

Definition 3.1. We say that the angle $\angle pxq$ exists if and only if the limit for $\varepsilon \to 0$ of the quantity $\frac{|\text{lip}(r_p+\varepsilon r_q)|^2(x)-|\text{lip}(r_p)|^2(x)}{2\varepsilon}$ exists. In this case we set

$$[0, \pi] \ni \angle pxq := \arccos\left(\lim_{\varepsilon \to 0} \frac{|\operatorname{lip}(r_p + \varepsilon r_q)|^2(x) - |\operatorname{lip}(r_p)|^2(x)}{2\varepsilon}\right). \tag{3.2}$$

Note that if (X, d) is a smooth Riemannian manifold and x is not in the cut locus of p and q, then $\angle pxq$ is the angle based at x between $\nabla r_p(x)$ and $\nabla r_q(x)$; in other words $\angle pxq$ is the angle based at x "in direction of p and q". As already mentioned, for a general triple pxq in a general metric space (X, d) the angle $\angle pxq$ may not exist; moreover, even if both $\angle pxq$ and $\angle qxp$ exist they may not be equal in general. On the other hand, such a definition satisfies some natural properties one expects from the geometric picture: the angle is invariant under a constant rescaling of the metric d, moreover for any two points $x, p \in X$ the angle $\angle pxp$ always exists and, if (X, d) is a length space, is equal to 0.

We now discuss an important class of metric measure spaces (X, d, m) where the angle exists and is symmetric in an a.e. sense, the so called Lipschitz-infinitesimally Hilbertian spaces.

Definition 3.2. A metric measure space (X, d, m) is said to be Lipschitz-infinitesimally Hilbertian if for any pair of Lipschitz functions $f,g \in LIP(X)$ both the limits for $\varepsilon \to 0$ of $\frac{(|\text{lip}(f+\varepsilon g)|^2(x)-|\text{lip}(f)|^2(x)}{2\varepsilon}$ and $\frac{|\text{lip}(g+\varepsilon f)|^2(x)-|\text{lip}(g)|^2(x)}{2\varepsilon}$ exist and are equal for m-a.e. $x \in X$, i.e.

$$\lim_{\varepsilon \to 0} \frac{|\operatorname{lip}(f + \varepsilon g)|^2(x) - |\operatorname{lip}(f)|^2(x)}{2\varepsilon} = \lim_{\varepsilon \to 0} \frac{|\operatorname{lip}(g + \varepsilon f)|^2(x) - |\operatorname{lip}(g)|^2(x)}{2\varepsilon}, \quad \text{m-a.e. } x.$$
 (3.3)

It is clear that if (X, d, m) is Lipschitz-infinitesimally Hilbertian then, given $p, q \in X$, for m-a.e. $x \in X$ both the angles $\angle pxq$, $\angle qxp$ exist and $\angle pxq = \angle qxp$.

Remark 3.3. The concept of Lipschitz-infinitesimally Hilbertian space was proposed in [26] as a variant of the notion of infinitesimally Hilbertian space introduced in [3, 17], using the language of minimal weak upper gradients; let us mention that Lipschitz-infinitesimally Hilbertian always implies infinitesimally Hilbertian, but the converse is not clear in general. An important class of spaces where also the converse implication holds is the one of locally doubling spaces satisfying a weak Poincaré inequality. Indeed, by a celebrated result of Cheeger [11], we have that for every $f \in \text{LIP}(X)$ it holds lip(f) = |Df| m-a.e., in other words the local Lipschitz constant is equal to the minimal weak upper gradient m-a.e. In particular for $\text{CD}^{\star}(K, N)$ spaces, $K \in \mathbb{R}$, $N \in [1, \infty)$ the two notions are equivalent. For more details we refer to [26, Remark 3.3].

It follows that $\mathrm{RCD}^{\star}(K, N)$ -spaces are Lipschitz-infinitesimally Hilbertian, for $N < \infty$; let us recall that the class of $\mathrm{RCD}^{\star}(K, N)$ -spaces include finite dimensional Alexandrov spaces with curvature bounded below and Ricci limit spaces as remarkable sub-classes.

3.2 Angle between two geodesics

First of all observe that if (X, d) is a metric space and $\gamma \in \text{Geo}(X)$ is a geodesic, then $|\dot{\gamma}_t| = d(\gamma_0, \gamma_1)$ for a.e. $t \in [0, 1]$; we will denote such a constant simply by $|\dot{\gamma}|$. The next definition is inspired by the De Giorgi's metric concept of gradient flow [15].

Definition 3.4 (A geodesic representing the gradient of a Lipschitz function). Let $f \in \text{LIP}(X)$ be a Lipschitz function on (X, d). We say that $\gamma \in \text{Geo}(X)$ represents ∇f at time 0, or $\gamma \in \text{Geo}(X)$ represents the gradient of f at the point $x = \gamma_0$ if the following inequality holds

$$\underline{\lim_{t \to 0}} \frac{f(\gamma_t) - f(\gamma_0)}{t} \ge \frac{1}{2} \mathrm{lip}(f)^2(\gamma_0) + \frac{1}{2} |\dot{\gamma}|^2. \tag{3.4}$$

Notice that the opposite inequality is always true, indeed

$$\overline{\lim_{t\to 0}}\frac{f(\gamma_t)-f(\gamma_0)}{t}\leq \operatorname{lip}(f)(\gamma_0)\;|\dot{\gamma}|\leq \frac{1}{2}\operatorname{lip}(f)^2(\gamma_0)+\frac{1}{2}|\dot{\gamma}|^2.$$

Hence $\gamma \in \text{Geo}(X)$ represents ∇f at time 0 if and only if the equality holds. Note that, in the case of Riemannian manifolds, γ represents ∇f at time 0 if and only if $\dot{\gamma}_0 = \nabla f$.

It is easy to check that the geodesic $\gamma \in \text{Geo}(X)$ represents the gradient of $f \in \text{LIP}(X)$ at $x \in X$ if and only if for every $\alpha \in (0, 1)$ the rescaled geodesic $\tilde{\gamma} \in \text{Geo}(X)$ defined by $\tilde{\gamma}_t := \gamma_{\alpha t}, \forall t \in [0, 1]$, represents the gradient of the Lipschitz function αf at x. In the next lemma we give a simple but important example of a geodesic representing the gradient of a function.

Lemma 3.5. Let (X, d) be a metric space, fix $p \in X$ and let $r_p(\cdot) := d(p, \cdot)$. If for some $x \in X$ there exists a geodesic $\gamma^{xp} \in \text{Geo}(X)$ such that $\gamma_0 = x$ and $\gamma_1 = p$ then γ^{xp} represents the gradient of $f(\cdot) := -d(p, x) r_p(\cdot)$ at x.

Proof. For every $t \in (0, 1)$ it holds

$$\frac{f(\gamma_t^{xp}) - f(\gamma_0^{xp})}{t} = d(p, x) \frac{d(p, x) - d(p, \gamma_t^{xp})}{t} = d(p, x) \frac{d(x, \gamma_t^{xp})}{t} = d(p, x) \frac{td(x, p)}{t}$$
$$= d(p, x)^2.$$

On the other hand, by triangle inequality it is clear that $lip(r_p) \le 1$ and with an analogous argument as above it is easily checked that actually $lip(r_p)(x) = 1$. Therefore $lip(f)(x) = d(p, x) =: |\dot{\gamma}^{xp}|$ and the claim follows. \Box

We can now define the angle between two geodesics.

Definition 3.6 (Angle between two geodesics). Let (X, d) be a metric space and let $\gamma, \eta \in Geo(X)$ be two geodesics with $\gamma_0 = \eta_0 = p$. Let $f \in LIP(X)$ be a Lipschitz function such that γ represents the gradient of f at time 0. We say that the angle $\angle \eta p \gamma$ exists if and only if the limit as $t \downarrow 0$ of $\frac{f(\eta_t) - f(\eta_0)}{t}$ exists. In this case we set

$$[0, \pi] \ni \angle \eta p \gamma := \arccos\left(\frac{1}{|\dot{\eta}||\dot{\gamma}|} \lim_{t \downarrow 0} \frac{f(\eta_t) - f(\eta_0)}{t}\right). \tag{3.5}$$

Remark 3.7 (Locality of the angle between two geodesics). It is easily seen that the angle between the two geodesics γ , $\eta \in \text{Geo}(X)$ at the point $p = \gamma_0 = \eta_0$ depend just on the germs of the curves at p. To see that, fix arbitrary T_{γ} , $T_{\eta} \in (0, 1)$ and call $\tilde{\gamma}$, $\tilde{\eta}$ the restrictions of γ , η to $[0, T_{\gamma}]$, $[0, T_{\eta}]$ properly rescaled, i.e:

$$\tilde{\gamma}(t) := \gamma(T_{\gamma}t), \quad \tilde{\eta}(t) := \eta(T_{\eta}t), \quad \forall t \in [0, 1].$$

Of course we still have $\tilde{\gamma}$, $\tilde{\eta} \in \text{Geo}(X)$, and it is readily seen that $\tilde{\gamma}$ represents the gradient of $\tilde{f} := T_{\gamma}f$. It follows that $\angle \eta p \gamma$ exists if and only if $\angle \tilde{\eta} p \tilde{\gamma}$ exists, and in this case it holds

$$\angle \eta p \gamma := \arccos\left(\frac{1}{|\dot{\eta}||\dot{\gamma}|} \lim_{t \to 0} \frac{f(\eta_t) - f(\eta_0)}{t}\right) = \arccos\left(\frac{1}{|\dot{\bar{\eta}}||\dot{\tilde{\gamma}}|} \lim_{t \to 0} \frac{\tilde{f}(\tilde{\eta}_t) - \tilde{f}(\tilde{\eta}_0)}{t}\right)$$

$$= \angle \tilde{\eta} p \tilde{\gamma}.$$

Remark 3.8 (Dependence on the function f). Note also in the generality of metric spaces, the angle $\angle \gamma p\eta$ as given in Definition 3.6 may depend on the function f chosen in (3.5) (for instance this is the case of a tree with a vertex in p and two edges made by γ and η). In case (X, d, \mathfrak{m}) is an $RCD^*(K, N)$ -space we will see later in the paper that actually the angle between two geodesics is well defined for \mathfrak{m} -a.e. base point p just in terms of the geometric data, so it does not depend on the choice of f. In the general case of a metric space, a way to overcome the problem would be to fix a canonical Lipschitz function f such that γ represents ∇f at time 0. In view of Lemma 3.5, a natural choice is to consider $f_{\gamma}(\cdot) := -d(\gamma_0, \gamma_1)d(\gamma_1, \cdot)$. In case (X, d, \mathfrak{m}) is not an $RCD^*(K, N)$ space we will tacitly make such a choice so to have a good definition.

The next goal is to relate the angle between three points with the angle between two geodesics, i.e. relate Definitions 3.1 and 3.6.

Theorem 3.9. Let (X, d) be a metric space and let $p \neq x \neq q \in X$ satisfy the following assumptions:

- the angle $\angle pxq$ exists in the sense of Definition 3.1,
- there exist geodesics γ^{xp} , $\gamma^{xq} \in \text{Geo}(X)$ from x to p and from x to q respectively.

Then the angle $\angle \gamma^{xp} x \gamma^{xq}$ exists in the sense of Definition 3.6 and

$$\angle \gamma^{xp} x \gamma^{xq} = \angle pxq. \tag{3.6}$$

Note that if (X, d) is a geodesic Lipschitz-infinitesimally Hilbertian space then for every given $p, q \in X$ the two assumptions of Theorem 3.9 are satisfied for m-a.e. $x \in X$. This is in particular the case for $RCD^*(K, N)$ spaces (see Remark 3.3).

Proof. Let $f_p(\cdot) := -\mathrm{d}(p,x) \, r_p(\cdot)$ and $f_q(\cdot) := -\mathrm{d}(p,x) \, r_q(\cdot)$. Recall from Lemma 3.5 that γ^{xp} represents ∇f_p at $x = \gamma_0^{xp}$ in the sense of Definition 3.4, i.e.

$$\underline{\lim_{t \downarrow 0}} \frac{f_p(\gamma_t^{xp}) - f_p(\gamma_0^{xp})}{t} \ge \frac{\lim_{t \downarrow 0} (f_p)^2(x)}{2} + \frac{|\dot{\gamma}^{xp}|^2}{2}.$$
(3.7)

On the other hand

$$\overline{\lim_{t\downarrow 0}} \frac{(f_p + \varepsilon f_q)(\gamma_t^{xp}) - (f_p + \varepsilon f_q)(\gamma_0^{xp})}{t} \le \frac{\lim(f_p + \varepsilon f_q)^2(x)}{2} + \frac{|\dot{\gamma}^{xp}|^2}{2}.$$
 (3.8)

Subtracting (3.7) from (3.8) yields

$$\frac{\lim_{t\downarrow 0} \varepsilon \frac{f_q(\gamma_t^{xp}) - f_q(\gamma_0^{xp})}{t} \le \frac{\operatorname{lip}(f_p + \varepsilon f_q)^2(x) - \operatorname{lip}(f_p)^2(x)}{2}$$

$$= \operatorname{d}(p, x)^2 \frac{\operatorname{lip}(r_p + \varepsilon r_q)^2(x) - \operatorname{lip}(r_p)^2(x)}{2}.$$
(3.9)

If $\varepsilon > 0$, dividing both sides by $\varepsilon d(p, x)^2$ and letting $\varepsilon \downarrow 0$ we get

$$\frac{\lim_{t\downarrow 0} \frac{1}{\mathrm{d}(p,x)^2} \frac{f_q(\gamma_t^{xp}) - f_q(\gamma_0^{xp})}{t} \leq \lim_{\varepsilon\downarrow 0} \frac{\mathrm{lip}(r_p + \varepsilon r_q)^2(x) - \mathrm{lip}(r_p)^2(x)}{2\varepsilon} = \cos(\angle pxq),$$

where in the last identity we used the assumption that $\angle pxq$ exists. Analogously, if $\varepsilon < 0$, dividing both sides by $\varepsilon d(p, x)^2$ and letting $\varepsilon \uparrow 0$ we get

$$\underline{\lim_{t\downarrow 0}} \ \frac{1}{\mathrm{d}(p,x)^2} \frac{f_q(\gamma_t^{xp}) - f_q(\gamma_0^{xp})}{t} \ge \underline{\lim_{\varepsilon \uparrow 0}} \ \frac{\mathrm{lip}(r_p + \varepsilon r_q)^2(x) - \mathrm{lip}(r_p)^2(x)}{2\varepsilon} = \cos(\angle pxq).$$

The combination of the last two inequalities gives the existence of the limit for $t\downarrow 0$ of $\frac{1}{\mathrm{d}(p,x)^2}\frac{f_q(\gamma_t^{xp})-f_q(\gamma_0^{xp})}{t}$ and, more precisely,

$$\lim_{t\downarrow 0} \frac{1}{\mathrm{d}(p,x)^2} \frac{f_q(\gamma_t^{xp}) - f_q(\gamma_0^{xp})}{t} = \cos(\angle pxq).$$

Multiplying and dividing by d(q, x) the left hand side, we get

$$\lim_{t\downarrow 0} \frac{1}{\mathrm{d}(p,x)} \frac{\frac{\mathrm{d}(q,x)}{\mathrm{d}(p,x)} f_q(\gamma_t^{xp}) - \frac{\mathrm{d}(q,x)}{\mathrm{d}(p,x)} f_q(\gamma_0^{xp})}{t} = \cos(\angle pxq).$$

Since by Lemma 3.5 we know that γ^{xq} represents the gradient of $-\mathrm{d}(q,x)r_q(\cdot) = \frac{\mathrm{d}(q,x)}{\mathrm{d}(p,x)}f_q(\cdot)$ at $x = \gamma_0^{xq}$ in the sense of Definition 3.4, we get that the left hand side coincides with $\angle \gamma^{xp} x \gamma^{xq}$ and the thesis follows.

3.3 Angles in Wasserstein spaces

In the Wasserstein space, we have the notion of "Plans representing gradients" which is similar to the one of "geodesic representing the gradient" above.

Definition 3.10 (Plans representing gradients, see [17]). Let (X, d, m) be a metric measure space, $g \in S^2(X)$ and $\Pi \in \mathcal{P}(C([0, 1], X))$ be a test plan. We say that Π represents the gradient of g if

$$\underline{\lim_{t\downarrow 0}}\int \frac{g(\gamma_t)-g(\gamma_0)}{t}\,\mathrm{d}\Pi(\gamma)\geq \frac{1}{2}\int |\mathrm{D}g|^2(\gamma_0)\,\mathrm{d}\Pi(\gamma)+\frac{1}{2}\,\overline{\lim_{t\downarrow 0}}\,\frac{1}{t}\,\iint_0^t |\dot{\gamma}_s|^2\,\mathrm{d}s\,\Pi(\gamma).$$

Let $(\mu_t) \in AC^2([0,1], \mathcal{P}_2(X))$ be with uniformly bounded densities, Π be its lifting given by Theorem 2.3, and let $\varphi \in S^2(X)$. In case $(X, \mathbf{d}, \mathbf{m})$ is infinitesimally Hilbertian, it is proved in [18, Theorem 4.6] that Π represents the gradient of φ if and only if

$$\frac{\mathrm{d}}{\mathrm{d}t}\Big|_{t=0} \int_{Y} f \,\mathrm{d}\mu_{t} = \int_{Y} \langle \nabla f, \nabla \varphi \rangle \,\mathrm{d}\mu_{0}, \quad \forall f \in S^{2}(X). \tag{3.10}$$

If (3.10) holds, we also say that the velocity field of μ_t at time 0 is $\nabla \varphi$.

Combing the above technical tools with ideas from Otto's calculus [28], we can define the angle between two geodesics in W_2 .

Definition 3.11 (Angle between curves in W_2). Let (X, d, \mathfrak{m}) be an infinitesimally Hilbertian metric measure space, let (μ_t) , $(\nu_t) \in AC^2([0, 1], \mathcal{P}_2(X))$ be with bounded compression, and such that $\mu_0 = \nu_0 =: \eta$. Assume

there exist lifting test plans of (μ_t) and (ν_t) representing the gradients of f and g respectively, for some $f,g \in$ $S^2(X)$. Then the angle between $\mu = (\mu_t)$ and $\nu = (\nu_t)$ at t = 0 is defined by

$$[0,\pi] \ni \angle^{W} \mu \eta \nu := \arccos \left(\frac{\int_{X} \langle \nabla f, \nabla g \rangle \, \mathrm{d} \eta}{\||\mathrm{D}g|\|_{L^{2}(X,\eta)} \||\mathrm{D}f|\|_{L^{2}(X,\eta)}} \right).$$

The same definition makes sense if f, $g \in S^2_{loc}(X)$ provided (μ_t) , (ν_t) have uniformly bounded supports.

From the formula (3.10), we can see that the value of the angle does not depend on the choice of f, g, but just on (μ_t) , (ν_t) .

Remark 3.12 (Locality of the angle in the Wasserstein space). The angle $\angle^W \mu \eta \nu$ depends just on the germs of the curves μ and ν at t=0; i.e., given $T_1,T_2\in(0,1)$, called $\tilde{\mu}_t:=\mu_{T_1t},\tilde{\nu}_t:=\nu_{T_2t}$ for all $t\in[0,1]$ the restrictions of μ , ν to $[0, T_1]$ and $[0, T_2]$ respectively, it holds $\angle \mu \eta \nu = \angle \tilde{\mu} \eta \tilde{\nu}$. Indeed let Π , lift of the curve $(\mu_t)_{t\in[0,1]}$, be a test plan representing the gradient of $f\in S^2(X)$; fix $T\in(0,1)$ and let $\tilde{\mu}_t:=\mu_{Tt}$ for every $t \in [0,1]$ be the restriction of the curve μ to [0,T]; called \tilde{H} the lift of $(\tilde{\mu}_t)_{t \in [0,1]}$, it is easily seen that \tilde{H} represents the gradient of $\tilde{f} := Tf$. The claim follows.

Thanks to the locality expressed in Remark 3.12, given two curves $(\mu_t)_{t \in [0,1]}$, $(\nu_t)_{t \in [0,1]}$ such that they are of bounded compression once restricted to [0, T] for some $T \in (0, 1)$, we can define the angle between them as the angle between their restrictions $\tilde{\mu}$, $\tilde{\nu}$ to [0, T]. This will be always tacitly assumed throughout the paper.

Let us briefly discuss the particular but important case when (μ_t) and (ν_t) are W_2 -geodesics in a general m.m.s. (X, d, m). If (μ_t) is a \mathcal{W}_2 -geodesic with bounded compression then any lift Π of (μ_t) is a test plan and moreover is an optimal dynamical plan, i.e. $\Pi \in \mathsf{OptGeo}(\mu_0, \mu_1)$. Moreover, as a consequence of the Metric Brenier Theorem proved in [2] (see also [18, Theorem 5.2] for the present formulation), if (μ_t) has bounded compression and $\varphi \in S^2(X)$ is a Kantorovich potential from μ_0 to μ_1 , then any lift Π of (μ_t) represents the gradient of $-\varphi$. Therefore, specializing Definition 3.11 to this case we get the following notion.

Definition 3.13 (Angle between geodesics in W_2). Let (X, d, m) be an infinitesimally Hilbertian metric measure space, let (μ_t) , (ν_t) be \mathcal{W}_2 -geodesics with bounded compression, and such that $\mu_0 = \nu_0 =: \eta$. Assume there exist φ , $\psi \in S^2(X)$ Kantorovich potentials from μ_0 to μ_1 and from ν_0 to ν_1 respectively. Then the angle between $\mu = (\mu_t)$ and $\nu = (\nu_t)$ at t = 0 is defined by

$$[0, \pi] \ni \angle^{W} \mu \eta \nu := \arccos \left(\frac{\int_{X} \langle \nabla \varphi, \nabla \psi \rangle \, \mathrm{d} \eta}{\| \| \mathrm{D} \varphi \|_{L^{2}(X, n)} \| \| \mathrm{D} \psi \|_{L^{2}(X, n)}} \right).$$

The same definition makes sense if φ , $\psi \in S^2_{loc}(X)$ provided (μ_t) , (ν_t) have uniformly bounded supports.

Note that, thanks to Otto calculus and (3.10), Definition 3.13 is the analog for W_2 geometry of the angle between two geodesics in a general metric space in the sense of Definition 3.6.

3.4 The case of $RCD^*(K, N)$ spaces

In Theorem 3.9 we related the angle between three points with the angle between two geodesics, i.e. we related Definitions 3.1 and 3.6. Now, adding a curvature assumption on the space, we wish to relate Definition 3.13 with Definition 3.1 and Definition 3.6, i.e. the angle between two geodesics in W_2 with the angle between three points and the angle between two geodesics of *X*. To this aim the next lemma will be useful.

Lemma 3.14. Let (X, d, m) be an $RCD^*(K, N)$ space, and let φ_1, φ_2 be locally Lipschitz functions on X. Then the functions

$$F^{+}(x) := \lim_{\epsilon \downarrow 0} \frac{|\operatorname{lip}(\varphi_{1} + \epsilon \varphi_{2})|^{2}(x) - |\operatorname{lip}(\varphi_{1})|^{2}(x)}{2\epsilon},$$

$$F^{-}(x) := \lim_{\epsilon \uparrow 0} \frac{|\operatorname{lip}(\varphi_{1} + \epsilon \varphi_{2})|^{2}(x) - |\operatorname{lip}(\varphi_{1})|^{2}(x)}{2\epsilon}$$

are well defined at every $x \in X$ and it holds

$$F^+ = F^- = \langle \nabla \varphi_1, \nabla \varphi_2 \rangle$$
, \mathfrak{m} -a.e..

Proof. From the definition of local Lipschitz constant we know that the function $\epsilon \mapsto |\text{lip}(\varphi_1 + \epsilon \varphi_2)|^2(x)$ is convex for any x. Consider the function

$$F_{\epsilon}(x) := \epsilon \mapsto \frac{|\operatorname{lip}(\varphi_1 + \epsilon \varphi_2)|^2(x) - |\operatorname{lip}(\varphi_1)|^2(x)}{2\epsilon},$$

and observe that $\epsilon \mapsto F_{\epsilon}(x)$ is non-decreasing on $(-\infty, 0)$ and $(0, +\infty)$ for any fixed x. Hence F^+ and F^- are well-defined for any point $x \in X$ as

$$F^+(x) := \inf_{\epsilon > 0} F_\epsilon(x) = \lim_{\epsilon \downarrow 0} F_\epsilon(x), \quad F^-(x) := \sup_{\epsilon < 0} F_\epsilon(x) = \lim_{\epsilon \uparrow 0} F_\epsilon(x).$$

Since $(X, \mathbf{d}, \mathbf{m})$ is a $\mathrm{RCD}^*(K, N)$ metric measure space, it holds a local Poincaré inequality and it is locally doubling. Then, from [11, Theorem 6.1], we know $\mathrm{lip}(f)(x) = |\mathrm{D}f|(x)$ for \mathbf{m} -a.e. $x \in X$. The definition of infinitesimal Hilbertian space and of $\langle \nabla \varphi_1, \nabla \varphi_2 \rangle$, then gives

$$\langle \nabla \varphi_1, \nabla \varphi_2 \rangle = \underset{\varepsilon>0}{\mathrm{ess}} \inf_{\epsilon>0} \frac{|\mathrm{D}(\varphi_1 + \epsilon \varphi_2)|^2 - |\mathrm{D}\varphi_1|^2}{2\epsilon} = \underset{\varepsilon>0}{\mathrm{ess}} \inf_{\epsilon>0} F_{\epsilon}$$

and

$$\langle \nabla \varphi_1, \nabla \varphi_2 \rangle = \operatorname{ess\,sup}_{\epsilon \leq 0} \frac{|\mathrm{D}(\varphi_1 + \epsilon \varphi_2)|^2 - |\mathrm{D}\varphi_1|^2}{2\epsilon} = \operatorname{ess\,sup}_{\epsilon \leq 0} F_{\epsilon}.$$

Hence

$$\langle \nabla \varphi_1, \nabla \varphi_2 \rangle = \underset{\epsilon>0}{\operatorname{ess inf}} F_{\epsilon} = \underset{\epsilon<0}{\operatorname{ess sup}} F_{\epsilon}.$$

In particular, we infer $F^+ = F^- = \langle \nabla \varphi_1, \nabla \varphi_2 \rangle$ m-a.e..

In the next result we relate Definition 3.13 with Definition 3.1. Before stating it, let us recall [22, Theorem 1.1] that if (X, d, m) is an $\mathrm{RCD}^*(K, N)$ m.m.s., $\mu_0, \mu_1 \in \mathcal{P}_2(X)$ with $\mu_0 \ll m$, then there exists a unique \mathcal{W}_2 geodesic connecting μ_0 and μ_1 ; let us mention that the same result holds more generally for essentially non-branching m.m.s satisfying the weaker $\mathrm{MCP}(K, N)$ condition [10].

Proposition 3.15. Let (X, d, \mathfrak{m}) be an $\mathrm{RCD}^{\star}(K, N)$ m.m.s. and fix $p, q \in X$. For every $x \in X$ and R > 0 let $\mu_0^R = \nu_0^R = \eta^R := \frac{1}{\mathfrak{m}(B_R(x))}\mathfrak{m} \sqcup B_R(x)$ and let $\mu^R := (\mu_t^R)_{t \in [0,1]}$, $\nu^R := (\nu_t^R)_{t \in [0,1]}$ be the unique \mathcal{W}_2 -geodesics from μ_0^R to δ_p and from ν_0^R to δ_q respectively. Then

$$\angle pxq = \lim_{R \to 0} \angle^W \mu^R \eta^R \nu^R, \quad \text{for } \mathfrak{m}\text{-a.e. } x \in X.$$
 (3.11)

Proof. Calling as usual $r_p(\cdot) := d(p, \cdot), r_q(\cdot) := d(q, \cdot)$, Lemma 3.14 implies

$$\cos(\angle pxq) := \lim_{\varepsilon \to 0} \frac{|\mathrm{lip}(r_p + \varepsilon r_q)|^2(x) - |\mathrm{lip}(r_p)|^2(x)}{2\varepsilon} = \langle \nabla r_p, \nabla r_q \rangle(x), \text{ m-a.e. } x \in X.$$

On the other hand, it is easily seen that $\varphi(\cdot) := \frac{1}{2} r_p(\cdot)^2$, $\psi(\cdot) := \frac{1}{2} r_q(\cdot)^2$ are Kantorovich potentials from $\eta = \mu_0^R$ to δ_p and from $\eta = \nu_0^R$ to δ_q respectively. Moreover the geodesics (μ_t^R) , (ν_t^R) have uniformly bounded supports,

and bounded compression once restricted to $[0, 1 - \delta]$ for every $\delta \in (0, 1)$, see for instance [22, Corollary 1.7]. Then, Definition 3.13 yields

$$\lim_{R\downarrow 0} \cos(\angle^{W} \mu^{R} \eta^{R} v^{R}) := \lim_{R\downarrow 0} \frac{\int_{X} \langle \nabla \varphi, \nabla \psi \rangle \, \mathrm{d}\eta}{\||\mathbf{D}\varphi|\|_{L^{2}(X,\eta)} \||\mathbf{D}\psi|\|_{L^{2}(X,\eta^{R})}}$$

$$= \lim_{R\downarrow 0} \frac{\frac{1}{\mathfrak{m}(B_{R}(x))} \int_{B_{R}(x)} r_{p}(y) \, r_{q}(y) \, \langle \nabla r_{p}, \nabla r_{q} \rangle (y) \, \mathrm{d}\mathfrak{m}(y)}{\|r_{p}\|_{L^{2}(X,\eta^{R})} \|r_{q}\|_{L^{2}(X,\eta^{R})}}$$

$$= \langle \nabla r_{p}, \nabla r_{q} \rangle (x) \quad \text{for m-a.e. } x \in X.$$

The combination of the two formulas gives the claim.

Remark 3.16. For uniformity with the rest of the paper we decided to state Proposition 3.15 for $RCD^*(K, N)$ spaces, but using the results of [10] the same conclusion holds for essentially non-branching Lipschitz-infinitesimally Hilbertian spaces satisfying MCP(K, N).

In the next result we relate Definition 3.6 with the optimal transport picture.

Proposition 3.17. Assume that (X, d, m) is an $RCD^*(K, N)$ metric measure space. Let (μ_t^1) and (μ_t^2) be W_2 -geodesics with bounded compression and with $\mu_0^1 = \mu_0^2 =: \eta$; let $\Pi_1, \Pi_2 \in OptGeo(X)$ be corresponding lifts. Then, for i = 1, 2, we can find $\Gamma_i \subset Geo(X)$ with $\Pi_i(\Gamma_i) = 1$, such that for η -a.e. x there exist unique geodesics $\gamma^{x,i} \in \Gamma_i$ with $\gamma_0^{x,i} = x$, and the angle $\angle \gamma^{x,1} x \gamma^{x,2}$ exists according to the Definition 3.6. Moreover

$$\cos \angle \gamma^{x,1} x \gamma^{x,2} = \cos \angle \gamma^{x,2} x \gamma^{x,1} = \lim_{t \downarrow 0} \frac{\varphi_1(\gamma_t^{x,2}) - \varphi_1(\gamma_0^{x,2})}{t \operatorname{lip}(\varphi_1)(x) |\dot{\gamma}^{x,2}|} = \lim_{t \downarrow 0} \frac{\varphi_2(\gamma_t^{x,1}) - \varphi_2(\gamma_0^{x,1})}{t \operatorname{lip}(\varphi_2)(x) |\dot{\gamma}^{x,1}|} \\
= \frac{\langle \nabla \varphi_1, \nabla \varphi_2 \rangle(x)}{\operatorname{lip}(\varphi_1)(x) \operatorname{lip}(\varphi_2)(x)}, \quad \text{for } \eta\text{-a.e. } x, \tag{3.12}$$

where $-\varphi_i \in S^2(X)$ is any locally Lipschitz Kantorovich potential from $\eta = \mu_0^i$ to μ_1^i .

Proof. From [3, 18] we know that any lift Π_i of (μ_i^i) represents the gradient of φ_i , for i=1,2, i.e.

$$\lim_{t\downarrow 0} \int \frac{\varphi_i(\gamma_t) - \varphi_i(\gamma_0)}{t} d\Pi_i(\gamma) \ge \frac{1}{2} \int |D\varphi_i|^2(\gamma_0) d\Pi_i(\gamma) + \frac{1}{2} \overline{\lim_{t\downarrow 0}} \frac{1}{t} \iint_0^t |\dot{\gamma}_s|^2 ds d\Pi_i(\gamma).$$

From [17, Proposition 3.11] we then get for i = 1, 2:

$$\lim_{t \downarrow 0} \frac{\varphi_i(\gamma_t) - \varphi_i(\gamma_0)}{t} = \frac{1}{2} |\mathrm{D}\varphi_i|^2(\gamma_0) + \frac{1}{2} |\dot{\gamma}|^2 = \frac{1}{2} |\mathrm{lip}(\varphi_i)|^2(\gamma_0) + \frac{1}{2} |\dot{\gamma}|^2, \ \Pi_i\text{-a.e.}\ \gamma.$$
 (3.13)

In other words, for Π_i -a.e. γ , we have that γ represents $\nabla \varphi_i$ at γ_0 , i = 1, 2.

For any $\epsilon > 0$, consider the function $\varphi_1 + \epsilon \varphi_2$ and observe that

$$\overline{\lim_{t\downarrow 0}} \frac{(\varphi_1 + \epsilon \varphi_2)(\gamma_t) - (\varphi_1 + \epsilon \varphi_2)(\gamma_0)}{t} \le \frac{1}{2} |\operatorname{lip}(\varphi_1 + \epsilon \varphi_2)|^2(\gamma_0) + \frac{1}{2} |\dot{\gamma}|^2, \quad \forall \gamma \in \operatorname{Geo}(X).$$
(3.14)

The difference between (3.14) and (3.13), for i = 1, gives

$$\epsilon \overline{\lim_{t\downarrow 0}} \frac{\varphi_2(\gamma_t) - \varphi_2(\gamma_0)}{t} \le \frac{|\operatorname{lip}(\varphi_1 + \epsilon \varphi_2)|^2(\gamma_0) - |\operatorname{lip}(\varphi_1)|^2(\gamma_0)}{2}, \quad \Pi_1\text{-a.e. }\gamma.$$
(3.15)

Multiplying by $\epsilon^{-1} > 0$ both sides of (3.15) yields

$$\overline{\lim_{t\downarrow 0}} \frac{\varphi_2(\gamma_t) - \varphi_2(\gamma_0)}{t} \le \frac{|\operatorname{lip}(\varphi_1 + \varepsilon \varphi_2)|^2(\gamma_0) - |\operatorname{lip}(\varphi_1)|^2(\gamma_0)}{2\varepsilon}, \quad \Pi_1 \text{-a.e. } \gamma.$$

Letting $\epsilon \downarrow 0$ and using Lemma 3.14, we infer

$$\overline{\lim_{t \downarrow 0}} \frac{\varphi_2(\gamma_t) - \varphi_2(\gamma_0)}{t} \le \langle \nabla \varphi_2, \nabla \varphi_1 \rangle (\gamma_0), \quad \Pi_1\text{-a.e. } \gamma.$$

Following verbatim the same arguments after (3.13), but now for ϵ < 0, gives

$$\underline{\lim_{t\downarrow 0}} \frac{\varphi_2(\gamma_t) - \varphi_2(\gamma_0)}{t} \ge \langle \nabla \varphi_2, \nabla \varphi_1 \rangle (\gamma_0), \quad \Pi_1\text{-a.e. } \gamma.$$

Since from [22, Theorem 3.4] we can find $\Gamma_i \subset \text{Geo}(X)$ with $\Pi_i(\Gamma_i) = 1$, such that for η -a.e. x there exists unique geodesics $\gamma^{x,i} \in \Gamma_i$ with $\gamma_0^{x,i} = x$, it follows that

$$\lim_{t\downarrow 0} \frac{\varphi_2(\gamma_t^{\chi,1}) - \varphi_2(\gamma_0^{\chi,1})}{t} = \langle \nabla \varphi_2, \nabla \varphi_1 \rangle(x), \quad \eta\text{-a.e. } x.$$
 (3.16)

Recalling from (3.13) that Π_2 -a.e. γ represents the gradient of φ_2 at γ_0 , we get that for η -a.e. x the geodesic $\gamma^{x,2}$ represents the gradient of φ_2 at x. Therefore (3.16) proves that for η -a.e. x the angle $\angle \gamma^{x,1} x \gamma^{x,2}$ exists according to Definition 3.6 and coincides with arccos $(\langle \nabla \varphi_2, \nabla \varphi_1 \rangle(x))$. With the same arguments, just exchanging i=1with i=2, we get that also $\angle \gamma^{x,2} x \gamma^{x,1}$ exists for η -a.e. x, and that the identities (3.12) hold.

4 The cosine formula for angles in $RCD^*(K, N)$ spaces

The goal of this section is to prove Theorem 4.4, stating that the cosine formula holds for the angle between two geodesics in an $RCD^*(K, N)$ space. The first lemma states the almost everywhere uniqueness and extendability of geodesics in $RCD^*(K, N)$ spaces; this fact is already present in the literature under slightly different formulations so we just briefly sketch the proof.

Lemma 4.1. Let (X, d, m) be an RCD^{*}(K, N) space for some $K \in \mathbb{R}$, $N \in (1, \infty)$, and fix $p, q \in X$. Then for m-a.e. x there exist unique geodesics γ^{xp} , $\gamma^{xq} \in \text{Geo}(X)$ such that

- $\begin{array}{l} \bullet \ \ \gamma_0^{xp} = \gamma_0^{xq} = x, \gamma_1^{xp} = p, \gamma_1^{xq} = q, \\ \bullet \ \ both \ \gamma^{xp} \ \ and \ \gamma^{xq} \ \ are \ extendable \ to \ geodesics \ \tilde{\gamma}^{xp} \ \ and \ \tilde{\gamma}^{xq} \ \ having \ x \ \ as \ interior \ point; \ in \ other \ words \ there \ \ exist \ \tilde{\gamma}^{xp}, \ \tilde{\gamma}^{xq} \in \operatorname{Geo}(X) \ \ and \ \bar{t} \in (0, 1) \ \ such \ that \ \tilde{\gamma}^{xp}_{\bar{t}}, \ \tilde{\gamma}^{xq}_{\bar{t}} = x \ \ and \ \tilde{\gamma}^{xp}_{[\bar{t},1]} = \gamma^{xp}_{[0,1]}, \ \tilde{\gamma}^{xq}_{[\bar{t},1]} = \gamma^{xq}_{[0,1]}. \end{array}$

Proof. **Step 1**. $\forall p \in X$, m-a.e. $x \in X$ is an interior point of a geodesic with end point at p. Fix $p \in X$ and R > 0. Consider

$$\mu_0 := \frac{1}{\mathfrak{m}(B_R(p))} \mathfrak{m} \llcorner B_R(p)$$
 and $\mu_1 := \delta_p$.

Analyzing the optimal transport from μ_0 to μ_1 by following verbatim the proof of [21, Lemma 3.1] (i.e. use Jensen's inequality and the convexity property of the entropy granted by the curvature condition), we get that for m-a.e. $x \in B_R(0)$ there exists a geodesic $\gamma \in \text{Geo}(X)$ such that $\gamma_1 = p$ and $\gamma_t = x$, for some $t \in (0, 1)$. The claim then follows by the arbitrariness of R > 0.

Step 2. $\forall p \in X$, m-a.e. $x \in X$ there exists a unique geodesic from x to p.

The uniqueness of geodesics connecting a fixed $p \in X$ and \mathfrak{m} -a.e. $x \in X$ is a consequence of [22, Theorem 3.5] applied to the optimal transportation from the measures μ_0 , μ_1 above.

Step 3. Applying steps 1 and 2 to p and q, since the union of two negligible sets is still negligible, the thesis follows.

The next lemma will be useful to get good estimates on harmonic approximations of distance functions.

Lemma 4.2. Let B be a unit ball in an RCD $^*(K, N)$ metric measure space $(X, d, m), K \in \mathbb{R}, N \in (1, \infty)$. Then there exists a function $G: B \to \mathbb{R}$ with $G \in D(\Delta, B)$ such that

$$\Delta G \sqcup B = (\Delta G \sqcup B) \mathfrak{m} \sqcup B$$
, $\Delta G \sqcup B = 1$, $0 \le G \le C$ on B ,

where C = C(K, N) > 0 is a constant which depends only on K and N.

Proof. Since $(X, \mathbf{d}, \mathfrak{m})$ is a RCD^{*}(K, N) metric measure space, it satisfies a local (1-2)-Poincaré inequality and it is locally doubling. It is also known [3, Remark 6.9 and Theorem 6.10] that the metric d is induced by the Dirichlet form $f \mapsto \int |\mathrm{D} f|^2 \, \mathrm{d} \mathfrak{m}$. Therefore the standing assumptions of [8] are fulfilled and from [8, Corollary 1.2] we know that for any $f \in L^p(B, \mathfrak{m})$, p > 2, there exists a function $u_f \in W_0^{1,2}(B)$ such that

$$\int_{B} \langle \nabla u_f, \nabla v \rangle \, \mathrm{d}\mathfrak{m} = \int_{B} f v \, \mathrm{d}\mathfrak{m}$$

for any $v \in W_0^{1,2}(B)$. In other words, we know $u_f \in D(\Delta, B)$ and

$$\Delta u_f = f$$
 m-a.e..

Furthermore, from [8, Theorem 4.1] we know

$$\sup_{B} |u_f| \leq c \,\mathfrak{m}(B)^{-\frac{1}{p}} ||f||_{L^p(B,\mathfrak{m})}$$

where *c* only depends on the constants in the Poincaré inequality and in the doubling condition. In our case, *c* only depends on *N* and *K*.

Now, choosing
$$f = 1$$
 on B , we get that $G := u_f + c$ satisfies the thesis with $C = 2c$.

Using Lemma 4.2, in the next proposition we prove a key estimate in order to establish the cosine formula for angles.

Proposition 4.3. Let (X, d, m) be an $RCD^*(K, N)$ metric measure space, for some $K \in \mathbb{R}$, $N \in (1, \infty)$, and $fix \ x_0 \in X$. Let $R \ge 2$, $p, \hat{p} \in X$ such that $d(x_0, p) + d(x_0, \hat{p}) = d(p, \hat{p})$, and $d(x_0, p), d(x_0, \hat{p}) \ge R$. We denote $b_p(\cdot) := d(p, \cdot) - d(p, x_0)$ and $b_{\hat{p}}(\cdot) := d(\hat{p}, \cdot) - d(\hat{p}, x_0)$. Assume that there exists a function $\Phi(R|K, N)$ satisfying $\lim_{R \to +\infty} \Phi(R|K, N) = 0$ for fixed K, N, such that $0 \le b_{\hat{p}}(x) + b_p(x) \le \Phi(R|K, N)$ for any $x \in B_1(x_0)$.

Then there exists a harmonic approximation \mathbf{b}_p of b_p with the following properties:

- 1. $\mathbf{b}_p b_p \in W_0^{1,2}(B_1(x_0)), \mathbf{b}_p \in D(\Delta, B_1(x_0))$ with $\Delta \mathbf{b}_{p} \cup B_1(x_0) = 0$,
- 2. it holds

$$\|\mathbf{b}_{p}-b_{p}\|_{L^{\infty}(B_{1}(x_{0}))}+\frac{1}{\mathfrak{m}(B_{1}(x_{0}))}\int_{B_{1}(x_{0})}|\mathrm{D}(\mathbf{b}_{p}-b_{p})|^{2}\,\mathrm{d}\mathfrak{m}\leq\Psi(R|K,N),\tag{4.1}$$

where $\Psi : \mathbb{R}^3 \to \mathbb{R}_{>0}$ satisfies $\lim_{R \to +\infty} \Psi(R|K,N) = 0$ for fixed K,N.

Proof. From Proposition 2.8 we know there exists \mathbf{b}_p satisfying (1) of the thesis. Similarly, we can find a harmonic approximation $\mathbf{b}_{\hat{n}}$ of $b_{\hat{n}}$.

We are then left to show the validity of the estimate (4.1). To this aim, let $G: B_1(x_0) \to \mathbb{R}_{\geq 0}$ be given by Lemma 4.2, so that

$$\Delta G \sqcup B_1(x_0) = (\Delta G \sqcup B_1(x_0)) \mathfrak{m} \sqcup B_1(x_0), \quad \Delta G \sqcup B_1(x_0) = 1, \quad 0 \le G \le C \quad \text{on } B_1(x_0),$$

where C(K, N) depends only on K, N and in particular is independent of R.

From Laplacian Comparison Theorem 2.9 we know that b_p , $b_{\hat{p}} \in D(\Delta, B_1(x_0))$ and

$$\Delta b_{p} \sqcup B_1(x_0) \le \Psi(R|K,N) \,\mathfrak{m} \tag{4.2}$$

and

$$\Delta b_{\hat{n}} \sqcup B_1(x_0) \le \Psi(R|K,N) \,\mathfrak{m} \tag{4.3}$$

for some suitable $\Psi: \mathbb{R}^3 \to \mathbb{R}_{>0}$ satisfying $\lim_{R \to +\infty} \Psi(R|K,N) = 0$ for fixed K,N. Then we have

$$\Delta(b_p - \mathbf{b}_p - \Psi G) \sqcup B_1(x_0) \le 0, \quad \Delta(-b_{\hat{n}} + \mathbf{b}_{\hat{n}} + \Psi G) \sqcup B_1(x_0) \ge 0.$$

Applying the comparison statement of Proposition 2.8 to $(-b_p + \mathbf{b}_p) + \Psi G$ we get that

$$-b_p + \mathbf{b}_p \le 2\Psi \sup_{B_1(x_0)} |G| \le 2C(K, N)\Psi, \quad \text{m-a.e. on } B_1(x_0).$$
 (4.4)

Analogously, applying the comparison statement of Proposition 2.8 to $-b_{\hat{p}} + \mathbf{b}_{\hat{p}} + \Psi G$ we get

$$-b_{\hat{p}} + \mathbf{b}_{\hat{p}} \le 2\Psi \sup_{B_1(x_0)} |G| \le 2C(K, N)\Psi, \quad \text{m-a.e. on } B_1(x_0).$$

By assumption, we know there exists a function $\Phi(R|K,N)$ satisfying $\lim_{R\to+\infty}\Phi(R|K,N)=0$ for fixed K,N, such that $0 \le b_{\hat{p}}(x) + b_p(x) \le \Phi(R|K,N)$ for any $x \in B_1(x_0)$. Using maximum principle of Proposition 2.8, we know

$$0 \le \mathbf{b}_{\hat{n}}(x) + \mathbf{b}_{p}(x) \le \Phi(R|K, N)$$

for any $x \in B_1(x_0)$. The combination of the last three estimates gives

$$b_{p} - \mathbf{b}_{p} = (b_{p} + b_{\hat{p}}) - (\mathbf{b}_{p} + \mathbf{b}_{\hat{p}}) + (-b_{\hat{p}} + \mathbf{b}_{\hat{p}}) \le \Phi(R|K, N) + 2C(K, N)\Psi(R|K, N), \quad \text{on } B_{1}(x_{0}). \tag{4.5}$$

Putting together (4.4) and (4.5), we get

$$||b_p - \mathbf{b}_p||_{L^{\infty}(B_1(X_0))} \le 2 C(K, N) \Psi(R|K, N) + \Phi(R|K, N).$$
(4.6)

Next, write $B = B_1(x_0)$ for short. Recalling that $\Delta \mathbf{b}_{p} \perp B = 0$, combining (4.2) with (4.6) and using that $(\mathbf{b}_p - b_p) \in W_0^{1,2}(B)$ in order to integrate by parts, we obtain

$$\int_{B} |\mathrm{D}(\mathbf{b}_{p} - b_{p})|^{2} d\mathfrak{m} = -\int_{B} (\mathbf{b}_{p} - b_{p}) d\left(\Delta(\mathbf{b}_{p} - b_{p})\right)
= \int_{B} (\|b_{p} - \mathbf{b}_{p}\|_{L^{\infty}(B)} + (\mathbf{b}_{p} - b_{p})) d\left(\Delta b_{p}\right) + \|b_{p} - \mathbf{b}_{p}\|_{L^{\infty}(B)} \int_{B} d\left(\Delta(\mathbf{b}_{p} - b_{p})\right)
\leq \int_{B} (\|b_{p} - \mathbf{b}_{p}\|_{L^{\infty}(B)} + (\mathbf{b}_{p} - b_{p})) \Psi(R|K, N) d\mathfrak{m}
\leq 2\|b_{p} - \mathbf{b}_{p}\|_{L^{\infty}(B)} \mathfrak{m}(B) \Psi(R|K, N)
\leq 2(2C(K, N) \Psi(R|K, N) + \Phi(R|K, N)) \mathfrak{m}(B) \Psi(R|K, N), \tag{4.7}$$

where we used that, since $(\mathbf{b}_p - b_p) \in W_0^{1,2}(B)$, it holds

$$\int_{B} d \left(\Delta(\mathbf{b}_{p} - b_{p}) \right) = \int_{B} 1 d \left(\Delta(\mathbf{b}_{p} - b_{p}) \right) = - \int_{B} \langle \nabla(\mathbf{b}_{p} - b_{p}), \nabla 1 \rangle d\mathfrak{m} = 0.$$

Summing up (4.6) and (4.7) we get (4.1) by renaming with $\Psi(R|K, N)$ the quantity $2(2C(K, N) \Psi(R|K, N) + \Phi(R|K, N)) \Psi(R|K, N) + 2C(K, N) \Psi(R|K, N) + \Phi(R|K, N)$.

Theorem 4.4. Let (X, d, m) be an $\mathrm{RCD}^*(K, N)$ space for some $K \in \mathbb{R}$, $N \in (1, \infty)$, and fix $p, q \in X$. Then for m-a.e. $x \in X$ let γ^{xp} , $\gamma^{xq} \in \mathrm{Geo}(X)$ be the unique geodesics from x to p and from x to q given by Lemma 4.1. We may also assume that the tangent cone at x is unique and isomorphic as m.m. space to $(\mathbb{R}^k, d_E, \mathcal{L}_k)$, for some $k = k(x) \in \mathbb{N} \cap [1, N]$. Let $r_i \downarrow 0$ be any sequence, $\bar{p}, \bar{q} \in \mathbb{R}^k$ be the limit points of $\gamma^{xp}(r_i)$, $\gamma^{xq}(r_i)$ under the rescalings $(X, r_i^{-1}d, \mathfrak{m}_{r_i}^x, x)$ which converge to $(\mathbb{R}^k, d_E, \mathcal{L}_k, 0)$ in p-mGH sense. Then

$$\angle \gamma^{xp} x \gamma^{xq} = \angle pxq = \angle \bar{p}O\bar{q} = \lim_{t\downarrow 0} \arccos \frac{2t^2 - d^2(\gamma_t^{xp}, \gamma_t^{xq})}{2t^2}, \quad \textit{for } m\text{-}\textit{a.e. } x.$$
 (4.8)

Proof. **Step 1.** Fix p, $q \in X$. Combining Theorem 2.16, Remark 3.3, Theorem 3.9 and Lemma 4.1 we get that for \mathfrak{m} -a.e. $x \in X$

- we can find unique geodesics γ^{xp} , $\gamma^{xq} \in \text{Geo}(X)$ such that $\gamma_0^{xp} = \gamma_0^{xq} = x$, $\gamma_1^{xp} = p$, $\gamma_1^{xq} = q$, and both γ^{xp} , γ^{xq} are extendable beyond x in the sense of Lemma 4.1, so we can assume that γ^{xp} , γ^{xq} could be extended to $\hat{p} := \gamma_{-\varepsilon}^{xp}$, $\hat{q} := \gamma_{-\varepsilon}^{xq}$ respectively, for some $\varepsilon > 0$,
- both the angles $\angle pxq$ and $\angle \gamma^{xp}x\gamma^{xq}$ exist in the sense of Definitions 3.1, 3.6 respectively, and $\angle pxq = \angle \gamma^{xp}x\gamma^{xq}$,
- $x \in X$ is a Lebesgue point for $\langle \nabla r_p, \nabla r_q \rangle$ so that

$$\cos \angle pxq := \langle \nabla r_p, \nabla r_q \rangle(x) = \lim_{r \downarrow 0} \frac{1}{\mathfrak{m}(B_r(x))} \int_{B_r(x)} \langle \nabla r_p, \nabla r_q \rangle \, d\mathfrak{m}, \tag{4.9}$$

• the tangent to *X* at *x* is unique and euclidean.

From the locality of the angle (see Remark 3.7) we know that

$$\angle pxq = \angle \gamma^{xp} x \gamma^{xq} = \angle (\gamma^{xp}|_0^s) x (\gamma^{xq}|_0^t), \quad \forall s, t \in (0, 1), \tag{4.10}$$

where $(\gamma|_0^s)_t := \gamma_{st}$ for all $t \in [0, 1]$.

Let $r_i \downarrow 0$ be any sequence and let $(X, r_i^{-1}d, \mathfrak{m}_{r_i}^{\chi}, x)$ be the corresponding sequence of rescaled spaces. Since by assumption x is regular, we know that $(X, r_i^{-1}d, \mathfrak{m}_{r_i}^X, x)$ p-mGH converge to $(\mathbb{R}^k, d_E, \mathcal{L}_k, 0)$ for some $k = k(x) \in \mathbb{N} \cap [1, N]$. Since by assumption both γ^{xp} and γ^{xq} are extendable beyond x, they converge in p-GH sense to half lines ℓ_p , ℓ_q in \mathbb{R}^k such that $O \in \ell_p \cap \ell_q$ and both ℓ_p , ℓ_q are extendable to full lines of \mathbb{R}^k . We parametrize such half lines on $[0, +\infty)$ such that for every t > 0 one has that $\ell_p(t)$, $\ell_q(t)$ are the limit points of $\gamma^{xp}(r_i t), \gamma^{xq}(r_i t)$ respectively. Denote by $\bar{p} = \ell_p(1), \bar{q} = \ell_q(1) \in \mathbb{R}^k$ be the limit points of $\gamma^{xp}(r_i), \gamma^{xq}(r_i)$. By the uniqueness of the tangent space, the parametrized half lines ℓ_p , ℓ_q and the points \bar{p} , $\bar{q} \in \mathbb{R}^k$ do not depend on the choice of the rescaling sequence (r_i) .

Let $\hat{\ell}_p$, $\hat{\ell}_q$ be the half lines in \mathbb{R}^k antipodal to ℓ_p , ℓ_q respectively; in other words $\ell_p \cup \hat{\ell}_p$ and $\ell_q \cup \hat{\ell}_q$ are straight lines in \mathbb{R}^k intersecting at O. We parametrize $\hat{\ell}_p$, $\hat{\ell}_q$ on $(-\infty, 0]$ such that $\frac{d}{dt}|_{t=t_1}\hat{\ell}_p(t) = \frac{d}{dt}|_{t=t_2}\ell_p(t)$, $\frac{d}{dt}|_{t=t_1}\hat{\ell}_q(t) = \frac{d}{dt}|_{t=t_1}\ell_p(t)$ $\frac{d}{dt}|_{t=t_2}\ell_q(t)$, for all $t_1 < 0 < t_2$.

Step 2. We claim that

$$\angle \gamma^{xp} x \gamma^{xq} = \angle pxq = \angle \bar{p} O \bar{q} = \angle \ell_p O \ell_q. \tag{4.11}$$

Since the first identity is true by construction, and the last is trivially true because the ambient space is \mathbb{R}^k , it is enough to show that $\angle pxq = \angle \bar{p}O\bar{q}$. Given any sequence of rescalings $r_i \downarrow 0$, let $d^i(\cdot, \cdot) := \frac{1}{r_i}d(\cdot, \cdot)$ and define

$$b_n^i(\cdot) := d^i(p, \cdot) - d^i(p, x), \qquad b_a^i(\cdot) := d^i(q, \cdot) - d^i(q, x).$$
 (4.12)

Set also b_p^{∞} , $b_q^{\infty}: \mathbb{R}^k \to \mathbb{R}$ to be the Busemann functions associated to ℓ_p , ℓ_q , i.e.

$$b_p^{\infty}(\cdot) := \lim_{t \to +\infty} t - \mathrm{d}_E(\ell_p(t), \cdot), \qquad b_q^{\infty}(\cdot) := \lim_{t \to +\infty} t - \mathrm{d}_E(\ell_q(t), \cdot). \tag{4.13}$$

Since by construction we know that $(X, d^i, x) \to (\mathbb{R}^k, d_E, 0)$ in p-GH sense, $\gamma_{r_i t}^{xp} \to \ell_p(t)$ and $\gamma_{r_i t}^{xq} \to \ell_q(t)$ for every t > 0, it follows that

$$b_p^i o b_p^\infty$$
, $b_q^i o b_q^\infty$ pointwise in the sense of Definition 2.11. (4.14)

More strongly, since b_p^i are all Lipschitz with unit Lipschitz constant, by an Arzelá-Ascoli procedure (see for instance in [27, Proposition 2.12]) we get that the convergences are uniform on bounded subsets, in the sense of Definition 2.11. In particular, since the measures \mathfrak{m}_{i}^{χ} are converging weakly to \mathcal{L}_{k} we get that

$$b_p^i o b_p^\infty$$
, $b_q^i o b_q^\infty$ strongly in L^2 in the sense of Definition 2.12. (4.15)

Define $b_{\hat{p}}^i, b_{\hat{p}}^{\infty}, b_{\hat{q}}^i, b_{\hat{q}}^{\infty}$ analogously to (4.12)-(4.13):

$$\begin{split} b^i_{\hat{p}}(\cdot) &:= \operatorname{d}^i(\hat{p}, \cdot) - \operatorname{d}^i(\hat{p}, x), \qquad b^i_{\hat{q}}(\cdot) := \operatorname{d}^i(\hat{q}, \cdot) - \operatorname{d}^i(\hat{q}, x), \\ b^{\infty}_{\hat{p}}(\cdot) &:= \lim_{t \to -\infty} -t - \operatorname{d}_E(\hat{\ell}_p(t), \cdot), \qquad b^{\infty}_{\hat{q}}(\cdot) := \lim_{t \to -\infty} -t - \operatorname{d}_E(\hat{\ell}_q(t), \cdot). \end{split}$$

With analogous arguments as above we get

$$b^i_{\hat p} o b^\infty_{\hat p}$$
, $b^i_{\hat q} o b^\infty_{\hat q}$ uniformly on bounded subsets in the sense of Definition 2.11.

Since we are working in the euclidean space \mathbb{R}^k , it is not difficult to see that $b_p^\infty + b_{\bar{p}}^\infty = 0$ and $b_q^\infty + b_{\bar{q}}^\infty = 0$. Note that such equalities holds more generally in manifolds with non-negative Ricci curvature: the argument, used in the proof of the Cheeger-Gromoll Splitting Theorem, goes via maximum principle; here in any case one can argue more directly by using the geometry of the euclidean space. It follows that $b^i_p+b^i_{\hat p} \to 0$ and $b^i_p+b^i_{\hat p} \to 0$

uniformly on bounded sets in the sense of Definition 2.11. Hence there exists a function $\Phi(R|K,N)$ satisfying $\lim_{R\to +\infty} \Phi(R|K,N)=0$ for fixed K,N, such that $0\leq b_p^i(y)+b_{\hat{p}}^i(y)\leq \Phi(\frac{1}{r_i}|K,N)$ and $0\leq b_q^i(y)+b_{\hat{q}}^i(y)\leq \Phi(\frac{1}{r_i}|K,N)$ for any $y\in B_1^{d^i}(x)$.

Using Proposition 4.3, for every $i \in \mathbb{N}$ we can construct harmonic approximations \mathbf{b}_p^i , \mathbf{b}_q^i of b_p^i , b_q^i , respectively, in the unit ball $B_1^{\mathbf{d}^i}(x)$ of the space $(X, \mathbf{d}^i, \mathfrak{m}_{r_i}^x)$. Since $\mathbf{d}^i(p, x) = \frac{1}{r_i}\mathbf{d}(p, x) \to \infty$, $\mathbf{d}^i(q, x) = \frac{1}{r_i}\mathbf{d}(q, x) \to \infty$ and the spaces $(X, \mathbf{d}^i, \mathfrak{m}_{r_i}^x)$ are $\mathrm{RCD}^*(r_i^2K, N)$, so in particular $\mathrm{RCD}^*(-1, N)$ for i large enough, we infer that

$$\|\mathbf{b}_{p}^{i}-b_{p}^{i}\|_{L^{\infty}(B_{1}^{d_{i}}(x))}+\frac{1}{\mathfrak{m}_{r_{i}}^{\chi}(B_{1}^{d_{i}}(x))}\int_{B_{1}^{d_{i}}(x)}|\mathrm{D}(\mathbf{b}_{p}^{i}-b_{p}^{i})|^{2}\,\mathrm{d}\mathfrak{m}_{r_{i}}^{\chi}\to0\quad\text{as }i\to\infty,\tag{4.16}$$

$$\|\mathbf{b}_{q}^{i} - b_{q}^{i}\|_{L^{\infty}(B_{1}^{d_{i}}(x))} + \frac{1}{\mathfrak{m}_{r_{i}}^{x}(B_{1}^{d_{i}}(x))} \int_{B_{1}^{d_{i}}(x)} |\mathrm{D}(\mathbf{b}_{q}^{i} - b_{q}^{i})|^{2} \, \mathrm{d}\mathfrak{m}_{r_{i}}^{x} \to 0 \quad \text{as } i \to \infty. \tag{4.17}$$

The combination of (4.15), (4.16) and (4.17) yields

$$\mathbf{b}^i_p \llcorner B^{\mathrm{d}_i}_1(x) \to b^\infty_p \llcorner B^{\mathrm{d}_E}_1(O), \quad \mathbf{b}^i_q \llcorner B^{\mathrm{d}_i}_1(x) \to b^\infty_q \llcorner B^{\mathrm{d}_E}_1(O) \quad \text{strongly in L^2}.$$

Since by construction $\Delta \mathbf{b}_{p}^{i} \sqcup B_{1}^{d_{i}}(x) = 0$, by Proposition 2.14 we get that

$$\mathbf{b}_p^i \llcorner B_1^{\mathrm{d}_i}(x) o b_p^\infty \llcorner B_1^{\mathrm{d}_E}(O), \quad \mathbf{b}_q^i \llcorner B_1^{\mathrm{d}_i}(x) o b_q^\infty \llcorner B_1^{\mathrm{d}_E}(O) \quad \mathrm{strongly\ in}\ W^{1,2}.$$

But then the gradient estimates in (4.16)-(4.17) give that

$$b_p^i \sqcup B_1^{d_i}(x) \to b_p^\infty \sqcup B_1^{d_E}(0), \quad b_q^i \sqcup B_1^{d_i}(x) \to b_q^\infty \sqcup B_1^{d_E}(0) \quad \text{strongly in } W^{1,2}.$$
 (4.18)

In particular, for every $\rho \in (0, 1)$ we have

$$\lim_{i \to \infty} \frac{1}{\mathfrak{m}_{r_i}^{\chi}(B_{\rho}^{d_i}(x))} \int_{B_{\rho}^{d_i}(x)} \langle \nabla b_p^i, \nabla b_q^i \rangle \, d\mathfrak{m}_{r_i}^{\chi} = \frac{1}{\mathcal{L}_k(B_{\rho}^{d_E}(O))} \int_{B_{\rho}^{d_E}(O)} \langle \nabla b_p^{\infty}, \nabla b_q^{\infty} \rangle \, d\mathcal{L}_k. \tag{4.19}$$

We now analyze the two sides of (4.19). Recalling (4.9), from the very definitions of $\mathfrak{m}_{r_i}^x$ and of d^i it follows that

$$\lim_{i \to \infty} \frac{1}{\mathfrak{m}_{r_i}^{\chi}(B_{\rho}^{d_i}(x))} \int_{B_{\rho}^{d_i}(x)} \langle \nabla b_p^i, \nabla b_q^i \rangle \, \mathrm{d}\mathfrak{m}_{r_i}^{\chi} = \lim_{i \to \infty} \frac{1}{\mathfrak{m}(B_{r_i\rho}^{\mathrm{d}}(x))} \int_{B_{r_i\rho}^{\mathrm{d}}(x)} \langle \nabla b_p, \nabla b_q \rangle \, \mathrm{d}\mathfrak{m}$$

$$= \angle pxq. \tag{4.20}$$

On the other hand, since b_p^∞ , b_q^∞ are the Busemann functions of the lines ℓ_p , ℓ_q in \mathbb{R}^k it is readily seen that

$$\lim_{\rho \downarrow 0} \frac{1}{\mathcal{L}_{k}(B_{\alpha}^{d_{E}}(O))} \int_{B_{q}^{d_{E}}(O)} \langle \nabla b_{p}^{\infty}, \nabla b_{q}^{\infty} \rangle \, \mathrm{d}\mathcal{L}_{k} = \angle \bar{p} O \bar{q}. \tag{4.21}$$

Putting together (4.19), (4.20) and (4.21) finally yields

$$\begin{split} \angle \bar{p}O\bar{q} &= \lim_{\rho \downarrow 0} \frac{1}{\mathcal{L}_k(B^{\mathrm{d}_E}_{\rho}(O))} \int_{B^{\mathrm{d}_E}_{\rho}(O)} \left\langle \nabla b^{\infty}_p, \nabla b^{\infty}_q \right\rangle \mathrm{d}\mathcal{L}_k = \lim_{\rho \downarrow 0} \lim_{i \to \infty} \frac{1}{\mathfrak{m}^{\mathrm{x}}_{r_i}(B^{\mathrm{d}_i}_{\rho}(x))} \int_{B^{\mathrm{d}_i}_{\rho}(x)} \left\langle \nabla b^i_p, \nabla b^i_q \right\rangle \mathrm{d}\mathfrak{m}^{\mathrm{x}}_{r_i} \\ &= \angle p \mathrm{x} q, \end{split}$$

as desired.

Step 3. We claim that

$$\angle pxq = \lim_{t \downarrow 0} \arccos \frac{2t^2 - \mathrm{d}(\gamma_t^{xp}, \gamma_t^{xq})^2}{2t^2}.$$
 (4.22)

To this aim, first of all observe that the cosine formula in \mathbb{R}^k ensures that

$$\angle \bar{p}O\bar{q} = \arccos\frac{2 - d_E(\bar{p}, \bar{q})}{2}.$$
 (4.23)

Let now $t_i \downarrow 0$ be any sequence and set $r_i := t_i$. Define the rescaled spaces (X, d^i, x) as above, with $d^i(\cdot, \cdot) := r_i^{-1}d(\cdot, \cdot)$. Notice first of all that the p-GH convergence of (X, d^i, x) to $(\mathbb{R}^k, d_E, 0)$ ensures that

$$\lim_{i \to \infty} d^i(\gamma_{t_i}^{xp}, \gamma_{t_i}^{xq}) = d_E(\bar{p}, \bar{q}). \tag{4.24}$$

It follows that

$$\lim_{i \to \infty} \frac{2t_i^2 - d(\gamma_{t_i}^{xp}, \gamma_{t_i}^{xq})^2}{2t_i^2} = \lim_{i \to \infty} \frac{2 - d^i(\gamma_{t_i}^{xp}, \gamma_{t_i}^{xq})^2}{2} \stackrel{(4.24)}{=} \frac{2 - d_E(\bar{p}, \bar{q})}{2} \stackrel{(4.23)}{=} \cos(\angle \bar{p}O\bar{q})$$

$$\stackrel{(4.11)}{=} \cos(\angle pxq). \tag{4.25}$$

Since the sequence $t_i \downarrow 0$ was arbitrary, (4.25) implies (4.22).

The thesis then follows by combining (4.11) and (4.22).

Remark 4.5. The cosine formula in \mathbb{R}^k ensures that

$$\angle \bar{p}O\bar{q} = \arccos \frac{s^2 + t^2 - d_E(\ell_p(s), \ell_q(t))^2}{2st} \quad \forall s, t > 0.$$
 (4.26)

It is natural to ask if the same formula holds in the non-smooth case. This remains an open problem even for Ricci limit spaces, so a fortiori in $RCD^*(K, N)$ spaces.

Here let us briefly mention that with analogous arguments as above one can show the weaker statement

$$\angle pxq = \lim_{s,t\downarrow 0, \frac{1}{r} \le \frac{s}{s} \le C} \arccos \frac{s^2 + t^2 - d(\gamma_s^{xp}, \gamma_t^{xq})^2}{2st}, \quad \text{for every } C \ge 1.$$
 (4.27)

To this aim let $s_i \downarrow 0$, $t_i \downarrow 0$ be any two sequences. Up to subsequences, we may assume that for all $i \in \mathbb{N}$ it holds either $0 \le s_i \le t_i$ or $0 \le t_i \le s_i$. Without loss of generality we may assume the first case. Up to further subsequences we may also assume that s_i/t_i has a limit $\bar{s} \in (0,1]$ as $i \to \infty$. Let $r_i := t_i \downarrow 0$ and define the rescaled spaces (X, \mathbf{d}^i, x) as above, with $\mathbf{d}^i(\cdot, \cdot) := r_i^{-1}\mathbf{d}(\cdot, \cdot)$. Calling $s_i' := r_i^{-1}s_i \to \bar{s}$, $t_i' := r_i^{-1}t_i = 1$, the p-GH convergence of (X, \mathbf{d}^i, x) to $(\mathbb{R}^k, \mathbf{d}_E, 0)$ ensures

$$\lim_{i\to\infty} d^i(\gamma_{s_i}^{xp}, \gamma_{t_i}^{xq}) = d_E(\ell_p(\bar{s}), \bar{q}).$$

Then we have

$$\lim_{i \to \infty} \frac{s_i^2 + t_i^2 - d(\gamma_{s_i}^{xp}, \gamma_{t_i}^{xq})^2}{2s_i t_i} = \lim_{i \to \infty} \frac{(s_i')^2 + (t_i')^2 - d^i(\gamma_{s_i}^{xp}, \gamma_{t_i}^{xq})^2}{2s_i' t_i'}$$

$$= \lim_{i \to \infty} \frac{(s_i')^2 + (t_i')^2 - d_E(\ell_p(s_i'), \ell_q(t_i'))^2}{2s_i' t_i'}$$

$$\stackrel{(4,26)}{=} \cos(\angle \bar{p} O \bar{q})$$

$$\stackrel{(4,11)}{=} \cos(\angle p x q). \tag{4.28}$$

Since the sequences s_i , $t_i \downarrow 0$ were arbitrary, (4.28) implies (4.27).

Acknowledgement: The first author would like to thank Nicola Gigli, Shouhei Honda and Karl Theodor Sturm for discussions on the topic.

References

[1] L. Ambrosio, N. Gigli, A. Mondino, and T. Rajala, *Riemannian Ricci curvature lower bounds in metric measure spaces with o-finite measure*, Trans. Amer. Math. Soc., 367 (2015), pp. 4661–4701.

- [2] L. Ambrosio, N. Gigli, and G. Savaré, Calculus and heat flow in metric measure spaces and applications to spaces with Ricci bounds from below, Invent. Math., 195 (2014), pp. 289-391.
- –, Metric measure spaces with Riemannian Ricci curvature bounded from below, Duke Math. J., 163 (2014), pp. 1405– [3] 1490.
- [4] L. Ambrosio and S. Honda, New stability results for sequences of metric measure spaces with uniform Ricci bounds from below. Preprint, arXiv:1605.07908, (2016).
- [5] L. Ambrosio, A. Mondino, and G. Savaré, Nonlinear diffusion equations and curvature conditions in metric measure spaces, preprint arXiv:1509.07273, to appear in Mem. Amer. Math. Soc.
- [6] A. Björn and J. Björn, Nonlinear Potential Theory on Metric Spaces, EMS Tracts Math., vol.17 of European Mathematical Society (EMS), Zürich, 2011.
- [7] D. Burago, Y. Burago, and S. Ivanov, A course in metric geometry, vol. 33 of Graduate Studies in Mathematics, American Mathematical Society, 2001.
- [8] M. Biroli and U. Mosco, A Saint-Venant type principle for Dirichlet forms on discontinuous media, Ann. Mat. Pura Appl. (4) 169 (1995), pp 125-181.
- [9] K. Bacher and K.T. Sturm, Localization and tensorization properties of the curvature-dimension condition for metric measure spaces, J. Funct. Anal., 259 (2010), pp. 28-56.
- [10] F. Cavalletti and A. Mondino, Optimal maps in essentially non-branching spaces, preprint arXiv:1609.00782, to appear in Comm. Cont. Math. DOI: 10.1142/S0219199717500079.
- [11] J. Cheeger, Differentiability of Lipschitz functions on metric measure spaces, Geom. Funct. Anal., 9 (1999), pp. 428-517.
- [12] J. Cheeger and T. Colding, On the structure of spaces with Ricci curvature bounded below I, J. Diff. Geom., 45, (1997), pp. 406
- —, On the structure of spaces with Ricci curvature bounded below II, J. Diff. Geom., 54, (2000), pp. 13–35. [13] —
- , On the structure of spaces with Ricci curvature bounded below III, J. Diff. Geom., 54, (2000), pp. 37 74.
- [15] E. De Giorgi, New problems on minimizing movements, Boundary Value Problems for PDE and Applications, C. Baiocchi and J. L. Lions, eds., Masson, (1993), pp. 81-98.
- [16] M. Erbar, K. Kuwada, and K.-T. Sturm, On the equivalence of the entropic curvature-dimension condition and Bochner inequality on metric measure spaces, Invent. Math., 201 (2015), pp. 993-1071.
- [17] N. Gigli, On the differential structure of metric measure spaces and applications, Mem. Amer. Math. Soc., 236, (1113), (2015).
- [18] N. Gigli and B.-X. Han, The continuity equation on metric measure spaces, Calc. Var. Partial Differential Equations, 53 (2015), pp. 149-177.
- [19] N. Gigli and A. Mondino, A PDE approach to nonlinear potential theory in metric measure spaces, J. Math. Pures Appl. 100 (2013), pp. 505-534.
- [20] N. Gigli, A. Mondino, and G. Savaré, Convergence of pointed non-compact metric measure spaces and stability of Ricci curvature bounds and heat flows, Proc. Lond. Math. Soc., 111, (2015), no. 5, pp. 1071-1129.
- [21] N. Gigli, A. Mondino, and T. Rajala, Euclidean spaces as weak tangents of infinitesimally Hilbertian metric measure spaces with Ricci curvature bounded below, J. Reine Angew. Math., 705, (2015), pp. 233-244.
- [22] N. Gigli, T. Rajala, and K.T. Sturm, Optimal maps and exponentiation on finite dimensional spaces with Ricci curvature bounded from below, J. Geom. Anal., 26, (2016), no. 4, pp. 2914-2929.
- [23] S. Honda, A weakly second-order differential structure on rectifiable metric measure spaces, Geom. Topol. 18 (2014), no. 2, pp 633-668.
- [24] S. Lisini, Characterization of absolutely continuous curves in Wasserstein spaces, Calc. Var. Partial Differential Equations, 28 (2007), pp. 85-120.
- [25] J. Lott and C. Villani, Ricci curvature for metric-measure spaces via optimal transport, Ann. of Math. (2), 169 (2009), pp. 903-
- [26] A. Mondino, A new notion of angle between three points in a metric space, J. Reine Angew. Math, 706 (2015), pp. 103-121.
- [27] A. Mondino and A. Naber, Structure Theory of Metric-Measure Spaces with Lower Ricci Curvature Bounds, preprint arXiv:1405.2222, to appear in Journ. Europ. Math. Soc.
- [28] F. Otto, The geometry of dissipative evolution equations: the porous medium equation, Comm. Partial Differential Equations, 26 (2001), pp. 101-174.
- [29] T. Rajala, Local Poincaré inequalities from stable curvature conditions on metric spaces, Calc. Var. Partial Differential Equations, Vol. 44, Num. 3, (2012), pp. 477-494.
- [30] K.T. Sturm, On the geometry of metric measure spaces. I, Acta Math., Vol. 196, (2006), 65–131.
- -, On the geometry of metric measure spaces. II, Acta Math., Vol. 196, (2006), 133–177.
- [32] C. Villani, Optimal transport. Old and new, Grundlehren der Mathematischen Wissenschaften, 338, Springer-Verlag, Berlin, (2009).